

(Q1) (Alternative way to the one suggested by the hints)

In (Q2) we find that for a grav. field φ ,

$$\Gamma_{i00} = \partial_i \varphi \quad \forall i \in \{1, 2, 3\}$$

Work in 2D spacetime. $\Rightarrow \Gamma_{100} = \varphi'$

Assume metric doesn't dep. on time:

$$\begin{aligned} \Gamma_{100} &= \frac{1}{2} (g^{-1})^{1\alpha} (g_{\alpha e, 0} + g_{0e, \alpha} - g_{00, e}) \\ &= \frac{1}{2} (g^{-1})^{1\alpha} g_{00, \alpha} \\ &= -\frac{1}{2} (g^{-1})^{11} g_{00, 1} \end{aligned}$$

Assume g is almost the Minkowski metric: $g = \begin{bmatrix} +1 & 0 \\ 0 & -1 \end{bmatrix} + h$

where h has "small" components.

$$\Rightarrow g^{-1} \approx \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} + O(h), \quad g_{00, 1} = h_{00, 1} = h_{00}'$$

$$\Rightarrow \Gamma_{100} = +\frac{1}{2} h_{00, 1} \Rightarrow \frac{1}{2} h_{00}' = \varphi'$$

$$\Rightarrow g_{00} =$$

If we pick a homogeneous grav. field, $\varphi(x) = \gamma x_1$ where γ is the gravitational const. (9.8) with γ instead of g to avoid conflict of notation with the metric g .

We know (eqn (4.4) in the script) that

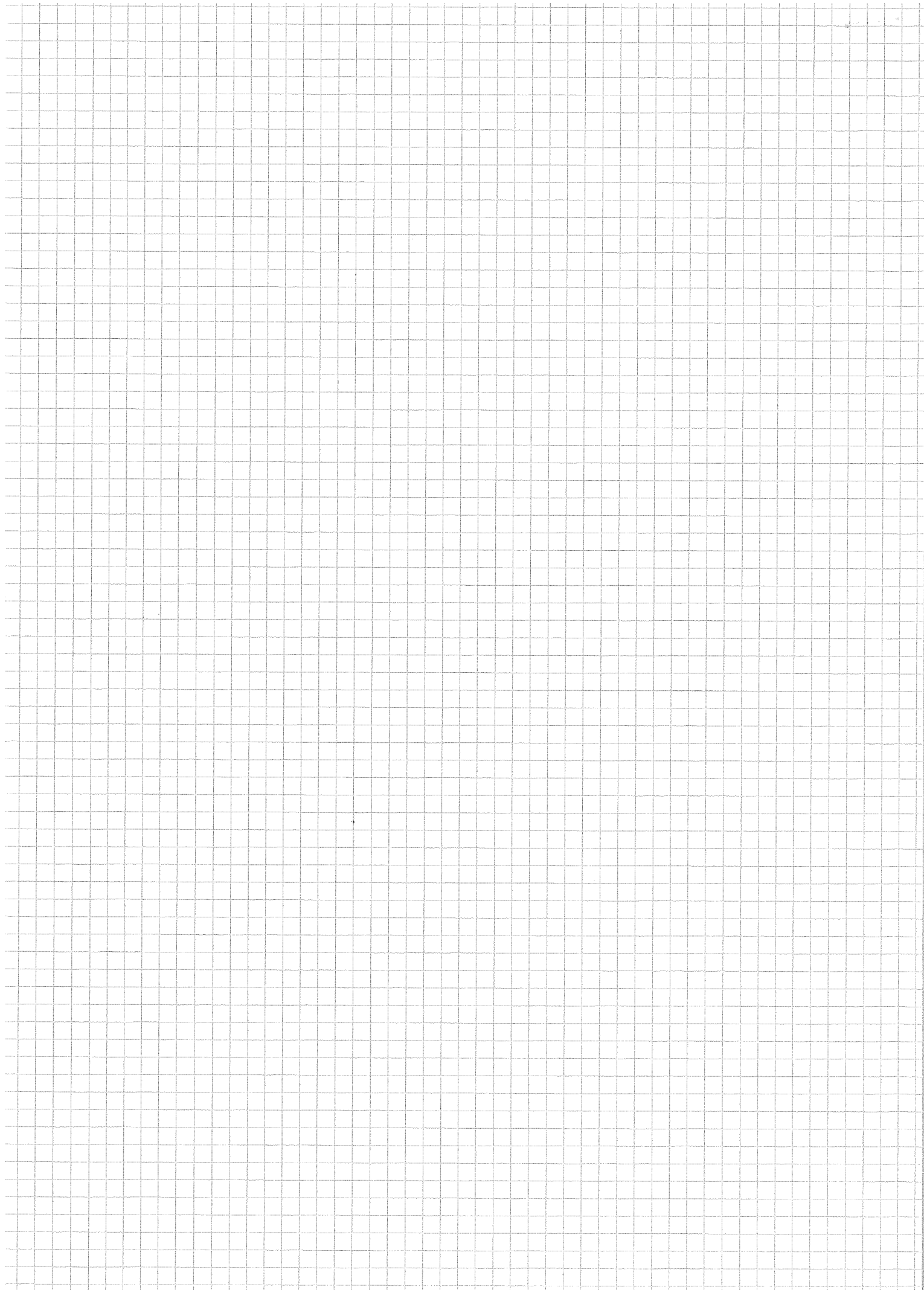
$$(\Delta \tau)^2 = g_{00}(x) (\Delta t)^2 \Rightarrow \Delta \tau = \sqrt{g_{00}(x)} \Delta t$$

where $\Delta \tau$ is the proper ^{elapsed} time of the clock, x_1 is its position, and Δt is the elapsed time in the chart.

We have one clock at height l (Q2) and another clock at height

$$\text{zero (Q1)} \Rightarrow \Delta \tau_1 = \quad , \quad \Delta \tau_2 =$$

$$\Rightarrow \Delta \tau_1 =$$



Q2. Newton's Eq-n as a Geodesic Eq-n

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In classical mechanics, when a particle is placed in a gravitational potential $\varphi: \mathbb{R}^3 \rightarrow \mathbb{R}$ its trajectory $\gamma: \mathbb{R} \rightarrow \mathbb{R}^3$ obeys the equation $\ddot{\gamma} = -(\nabla\varphi) \circ \gamma$. *

In 4D spacetime, a curve is given by a map $\tilde{\eta}: \mathbb{R} \rightarrow M$ where the domain \mathbb{R} is the parameter of the curve $\tilde{\eta}$ (not necessarily time) and M is the spacetime manifold, which by GR is a 4D pseudo-Riemannian manifold.

Hence we should have a chart $\psi: U \rightarrow \mathbb{R}^4 \quad \exists U \in \text{Open}(M)$.

Then $\psi \circ \tilde{\eta}|_{\tilde{\eta}^{-1}(U)}: \tilde{\eta}^{-1}(U) \rightarrow \mathbb{R}^4$ is a curve in \mathbb{R}^4 parametrized by an open subset of \mathbb{R} (since $\tilde{\eta}$ is cont.).

$$V := \tilde{\eta}^{-1}(U) \in \text{Open}(\mathbb{R})$$

$$\eta := \psi \circ \tilde{\eta}|_V$$

η is hence a curve in \mathbb{R}^4 .

If we want to "connect" η with γ then we could "declare" that $\eta_0 := \downarrow$ (viewing the 0th component of \mathbb{R}^4 as that which corresponds to time) and $(\eta_i)_{i=1,2,3} := \gamma$.

Then we can figure out which EoM η satisfies:

We let now the dot be the derivative w.r.t. the param. of the curve η , $\lambda \in V$:

$$\dot{\eta}_0(\lambda) \equiv 1 \Rightarrow \dot{\eta}_0(\lambda) = 1 \Rightarrow \ddot{\eta}_0(\lambda) = 0$$

From * we get:

$$\begin{aligned} (\partial_t^2 \gamma)_i(t) &= -(\nabla\varphi)_i(\gamma(t)) = -(\partial_i \varphi)(\gamma(t)) \\ \Rightarrow \ddot{\eta}_i &\equiv (\partial_\lambda^2 \eta)_i(\lambda) + (\partial_i \varphi)(\eta(\lambda)) = 0 \end{aligned}$$

But $\dot{\eta}_0(\lambda) = 1$, so we may rewrite this as:

$$\ddot{\eta}_i + [(\partial_i \varphi) \circ \eta] \dot{\eta}_0 \dot{\eta}_0 = 0$$

We then recall what the geodesic eq-n for a curve $\eta: V \rightarrow \mathbb{R}^4$ would look like given Christoffel symbols:

$$\ddot{\gamma}_i + \Gamma_{ijk} \dot{\gamma}_j \dot{\gamma}_k = 0$$

For the two equations to coincide, we may define

$$\Gamma_{0jk} := \quad \text{for } \ddot{\gamma}_0 = 0.$$

$$\Gamma_{i00} := \quad \text{for the other three eqns, } i \in \{1, 2, 3\}.$$

All other components should be zero.

The corresponding Riemann curvature tensor, given by eqn (2.15) in the script:

$$R_{ijkl} \stackrel{(2.15)}{=} \Gamma_{ljk} \Gamma_{iks} - \Gamma_{shj} \Gamma_{ils} =$$

$$=$$

Also $\Gamma_{ikhj}, \ell = 0$ if $k \in \{1, 2, 3\}$.

$$\Rightarrow R_{ijkl} =$$

$$\Rightarrow R_{i0k0} =$$

$$R_{ijk0} =$$

If Γ were to be associated with a Levi-Civita connection, we know it should obey the symmetries on pp. 29, namely:

$$\begin{cases} g_{ii} R_{ijkl} = -g_{jj} R_{jiikl} \\ g_{ii} R_{ijkl} = g_{kk} R_{klij} \end{cases}$$

$$g_{ii} R_{i0k0} = \stackrel{?}{=} -g_{0j} \underbrace{R_{j0k0}}_{=0}$$

$$\Rightarrow \text{Not true unless } \boxed{} \Leftrightarrow$$

Q3. On the Levi-Civita Connection

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Let M be a manifold.

Def.: A subset $N \subseteq M$ is called an embedded submanifold of M iff $\forall p \in N \exists$ chart $\psi: U \rightarrow \mathbb{R}^m$ ($U \in \text{Open}(M)$, $p \in U$) s.t. $\psi(U \cap N) = \psi(U) \cap W$ where $W \subseteq \mathbb{R}^m$ is a subspace of $\dim. n = \dim(N)$, $0 \leq n \leq m$. The atlas for N is the pairs $(U \cap N, \psi|_{U \cap N})$.

$\iota: N \hookrightarrow M$, the inclusion map, induces a tangent map:

$(d\iota)_p: T_p N \hookrightarrow T_p M$, which is injective as well. (We also use $\iota_* \equiv d\iota$)

$$(d\iota)_p(\iota_* \alpha) = (d\iota)_p(\iota_* \alpha)$$

$$\Downarrow$$

$$\iota_*(\iota_* \alpha) = \iota_*(\iota_* \alpha)$$

$$\iota_*(f \circ \alpha) = \iota_*(f \circ \alpha) \quad \forall f \in \mathcal{F}(M)$$

But since $\forall g \in \mathcal{F}(N)$, \exists extension $\tilde{g} \in \mathcal{F}(M)$: $\tilde{g} \circ \iota = g$, $\iota_* \equiv \iota_*$.

\Rightarrow We think of $T_p N$ as a linear subspace of $T_p M$.

If $g \in \Gamma(T^*M \otimes T^*M)$ is a metric on M , a \perp metric $h \in \Gamma(T^*N \otimes T^*N)$ is induced via the eqn:

$$h(X, Y) := g((d\iota)(X), (d\iota)(Y))$$

$$\forall (X, Y) \in \Gamma(TN)^2.$$

$\forall p \in N$, define $P_p: T_p M \rightarrow T_p N$ as the orthogonal proj. assoc. to g : g is an inner-prod. on $T_p M$, so it defines what it means for two vectors to be orthogonal;

$$[u \perp_{g_p} v \iff g_p(u, v) = 0] \quad \forall (u, v) \in (T_p M)^2$$

Then the orthogonal projection P_p satisfies:

$$P_p u = u \quad \forall u \in T_p N \quad \text{bes. } P_p \text{ is a proj.}$$

$$g_p(v, (d\iota)u) = g_p(v, P_p(d\iota)u) = g_p(P_p v, (d\iota)u)$$

\uparrow
 bes. it is an orthogonal proj.

$$\forall u \in T_p N, v \in T_p M.$$

6] Let ∇ be the g -Levi-Civita connection.

Cl. $(S_X Y)_p = \tilde{P}_p(\nabla_{\tau_X} \tau_* Y) \quad \forall (X, Y) \in T(TN)^2, p \in N.$
 with \tilde{P} being the proj. E w/ its codomain TN .

Pf. Let $(X, Y) \in T(TN)^2, p \in N.$

Cl. $\tilde{P} \circ \nabla_{\tau_X} \tau_*$ is an affine connection on N .

Pf. (i) W.T.S. $\tilde{P} \circ \nabla_{\tau_X} \tau_* Y$ is $\mathcal{F}(N)$ -lin. in X :

$$\tilde{P} \circ \nabla_{\tau_{X_1+X_2}} \tau_* Y =$$

↑
linear

∇ is an affine conn. \Downarrow

\tilde{P} is linear \Downarrow

$$(\tilde{P} \circ \nabla_{\tau_{fX}} \tau_* Y)_p =$$

↑
 $\tau_* X = f \tau_* X$

✓
 $f \in \mathcal{F}(N)$
 $p \in N$

\tilde{P}_p is linear \Downarrow

✓

(ii) W.T.S. $\tilde{P} \circ \nabla_{\tau_X} \tau_* Y$ is \mathbb{R} -lin. in Y :

(iii) W.T.S. $\tilde{P} \circ \nabla_{\tau_X} \tau_* Y$ obeys Leibniz in Y :

$$(\tilde{P} \circ \nabla_{\tau_X} \tau_* fY)_p =$$

↑
 \tilde{f} extends f to N (trivially)

$$\tilde{P} \circ (\tau_x)_* = 1$$

$\Rightarrow \tilde{P} \circ \nabla_{\tau_x} \tau_x \circ$ is an affine conn. on M . ✓

Cl. The torsion of ∇ is zero.

Pr. The torsion is given by:

$$\tilde{P} \nabla_{\tau_x} \tau_x Y - \tilde{P} \nabla_{\tau_x Y} \tau_x X - [X, Y] =$$

$\tilde{P} d_x = 1$
 \downarrow
 $= i$

Cl. $\tau_x [X, Y] = [\tau_x X, \tau_x Y]$

Pr.

$= \tilde{P} \circ$ as ∇ has no torsion.

Cl. $(\tilde{P} \circ \nabla_{\tau_x})(h) = 0$

Pr. We use the inductive formula: $Z(T(X_1, \dots, X_k, \mu_1, \dots, \mu_r))$

$$(\nabla_Z T)(X_1, \dots, X_k, \mu_1, \dots, \mu_r) = \nabla_Z (T(X_1, \dots, X_k, \mu_1, \dots, \mu_r)) - T(\nabla_Z X_1, X_2, \dots, X_k, \mu_1, \dots, \mu_r) - \dots - T(X_1, \dots, X_k, \mu_1, \dots, \nabla_Z \mu_r)$$

to get:
 $(\tilde{P} \nabla_{\tau_x} h)(X, Y) =$

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Orthogonal
proj

$= 0$ in particular,

Hence by uniqueness of the Levi-Civita
connection we obtain the result.

