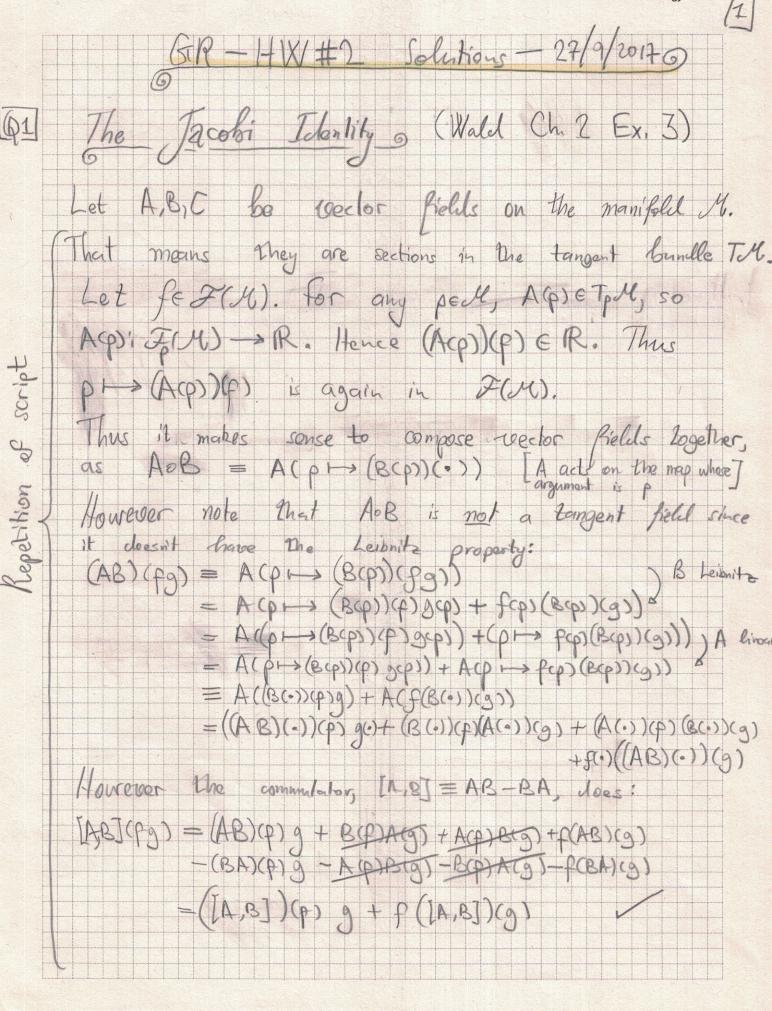


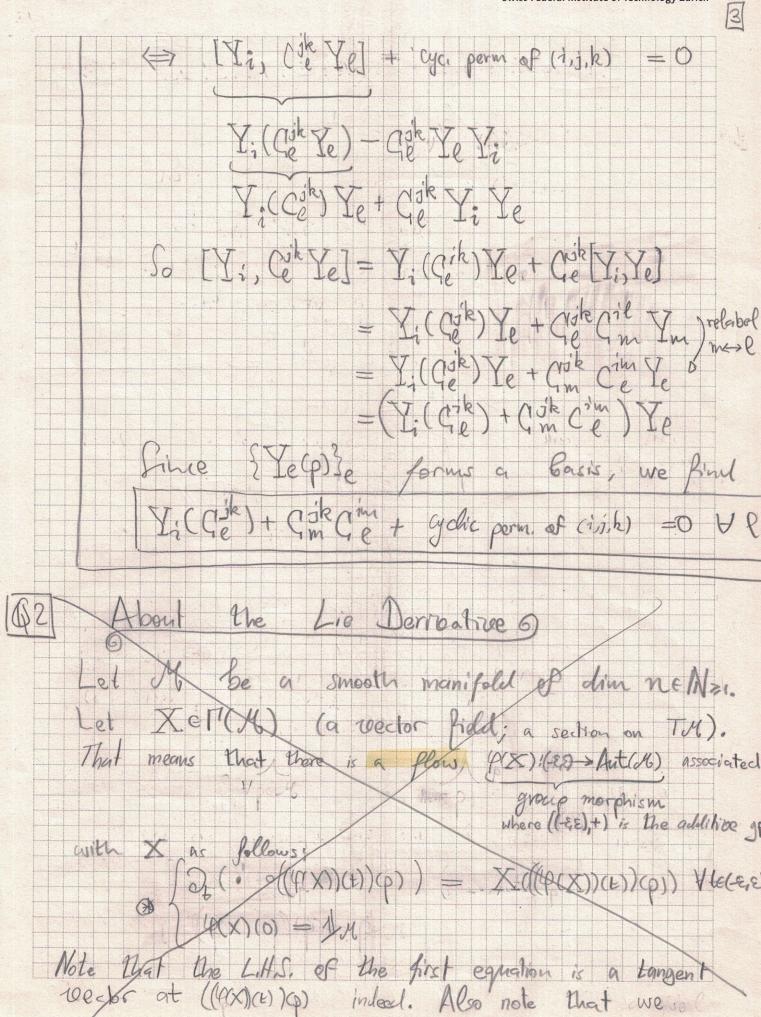
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$$\begin{split} \hline [B(G)] & (\underline{U}_{n}^{o} [A,]B, C]] + (Cyclic permutations) = 0 \\ \hline P_{1,0}^{o} [A, [B, C]] = [A, BC-CE] = ABC-ACE-BCA+CEA \\ & [A, [B, C]] + (Cyc. perm.) = [A,]BC] + [B,]C, A] + [C, [A, B]] \\ &= ABC-ACS-BCA+CEA \\ &+ BEA-BAC-CAE+BAC = 0 \\ \hline (AC) Let \{Y_{1,0}^{a} = \Gamma(A) & be n-vector fields c.t. Vpc.d. \\ \{Y_{1,0}^{a} \}_{i=1}^{i} is a basis of Tp.d. \\ \hline Note [Y_{1,1}, Y_{1,j}] \in \Pi(A_{0}), so the may expand it \\ at each point pc.d. using the basis $\{Y_{1,0}^{a} \}_{i=1}^{i} : \\ IX_{1,1} Y_{2,j}] =: CR^{i} Y_{k} (this eqn defines the \\ expansion coefficients CR^{i}). \\ \hline (C_{1,0}^{c} CR^{i} = -C_{k}^{i} \\ \hline Mole that since everything is a function of the point pc.d. \\ C_{k}^{c} are also p-dependent and so they are maps of the. \\ \hline (C_{k}^{c} CR^{i} C + C C + C = Y_{k}C_{k} + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (C_{k}^{c} CR^{i} C + C C + C = Y_{k}C_{k} + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C C + C = Y_{k}C_{k} + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{2,0}^{c} CR^{i} C + C C + C = Y_{k}C_{k} + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C C + C + C = Y_{k}C_{k} + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + C + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + C + Y_{k}C_{k} + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + C + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + C + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR^{i} C + C + C + C + C + C + Y_{k}C_{k} \\ \hline (M_{1,0}^{c} CR$$$



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HW2 Q2

September 29, 2017

1 Notation

Let $\varphi: U_{\varphi} \to \mathbb{R}^n$ and $\psi: U_{\psi} \to \mathbb{R}^n$ be two charts near some $p \in \mathcal{M}$.

Then we define basis vectors of $T_p\mathcal{M}$ corresponding to these charts as $d_i^{\varphi} := [\partial_i (\cdot \circ \varphi^{-1})] \circ \varphi$. Note that this is really a vector field defined in a neighborhoud of p. In a point $q \in \mathcal{M}$ it is a tangent vector: d_i^{φ} at q is $\partial_i|_{\varphi(q)} (\cdot \circ \varphi^{-1})$. There are analogous definitions for ψ . We define the expansion coefficients of a vector field X in the basis corresponding to φ as X_i^{φ} :

$$X = X_i^{\varphi} d_i^{\varphi}$$

so that $X_i^{\varphi} \equiv X(\varphi_i)$ with $\varphi_i := \pi_i \circ \varphi$ and $\pi_i : \mathbb{R}^n \to \mathbb{R}$ is the natural projection. The transition rule (going from φ to ψ) for the expansion coefficients may be derived easily as

$$\begin{aligned} X_i^{\psi} &\equiv X\left(d_i^{\psi}\right) \\ &= X_j^{\varphi} d_j^{\varphi}\left(\psi_i\right) \end{aligned}$$

so that we define

$$M_{ij}^{\psi\varphi} := d_j^{\varphi}(\psi_i)$$

and get

$$X_i^{\psi} = M_{ij}^{\psi\varphi} X_j^{\varphi}$$

Similarly, we can move the basis vectors themselves:

$$\begin{array}{lcl} d_i^{\psi} & = & d_i^{\psi}\left(\varphi_j\right) d_j^{\varphi} \\ & = & M_{ji}^{\varphi\psi} d_j^{\varphi} \\ & =: & N_{ij}^{\psi\varphi} d_j^{\varphi} \end{array}$$

We also have a natural basis for $(T_p\mathcal{M})^*$, given by the dual of d_i^{φ} . Explicitly it is given by

$$e_i^{\varphi} := \cdot (\varphi_i)$$

That is, given any tangent vector X, $e_i^{\varphi}(X) \equiv X(\varphi_i) = X_i^{\varphi}$. The expansion coefficients of a 1-form ω are given by

$$\omega_i^{\varphi} = \omega \left(d_i^{\varphi} \right)$$

so that

$$\omega = \omega_i^{\varphi} e_i^{\varphi}$$

and the transformation rule for the expansion coefficients is

$$\begin{split} \omega_i^{\psi} &\equiv \omega \left(d_i^{\psi} \right) \\ &= \omega_j^{\varphi} e_j^{\varphi} \left(d_i^{\psi} \right) \end{split}$$

But $e_{j}^{\varphi}\left(d_{i}^{\psi}\right) \equiv d_{i}^{\psi}\left(\varphi_{j}\right) = N_{ij}^{\psi\varphi}$ so that we get

$$\omega_i^{\psi} = N_{ij}^{\psi\varphi}\omega_j^{\varphi}$$

and of course the dual basis vectors transform again in the opposite way compared to the expansion coefficients:

$$e_i^{\psi} = e_i^{\psi} \left(d_j^{\varphi} \right) e_j^{\varphi}$$
$$= d_j^{\varphi} \left(\psi_i \right) e_j^{\varphi}$$
$$= M_{ii}^{\psi\varphi} e_j^{\varphi}$$

We find that the expansion coefficients of a general (k, l) tensor T transform as

$$T^{\psi}_{i_1\cdots i_k j_1\cdots j_l} = M^{\psi\varphi}_{i_1i'_1}\cdots M^{\psi\varphi}_{i_ki'_k}N^{\psi\varphi}_{j_1j'_1}\cdots N^{\psi\varphi}_{j_lj'_l}T^{\varphi}_{i'_1\cdots i'_kj'_1\cdots j'_l}$$

2 Properties of the Transition Matrices

 $1 \ Claim. \ {\rm We \ have} \ N_{ij}^{\psi\varphi}M_{ik}^{\psi\varphi}=\delta_{jk} \ {\rm and} \ N_{ij}^{\psi\varphi}M_{kj}^{\psi\varphi}=\delta_{ik}.$

Proof. We start by plugging in the definitions

$$N_{ij}^{\psi\varphi}M_{ik}^{\psi\varphi} \equiv d_i^{\psi}(\varphi_j) d_k^{\varphi}(\psi_i)$$

we swap out φ_j and ψ_i for e_j^{φ} and e_i^{ψ} respectively, because it is more transparent then that these are dual vectors to the *d*'s. We get

$$\begin{split} N_{ij}^{\psi\varphi}M_{ik}^{\psi\varphi} &= d_i^{\psi}\left(e_j^{\varphi}\right)d_k^{\varphi}\left(e_i^{\psi}\right) \\ &= d_j^{\varphi*}\left(d_i^{\psi}\right)d_i^{\psi*}\left(d_k^{\varphi}\right) \\ &= \left\langle d_j^{\varphi}, d_i^{\psi}\right\rangle \left\langle d_i^{\psi}, d_k^{\varphi}\right\rangle \\ &= \left\langle d_j^{\varphi}, d_i^{\psi}\otimes d_i^{\psi*}d_k^{\varphi}\right\rangle \end{split}$$

Now we use the fact that $d_i^{\psi} \otimes d_i^{\psi*} = 1$ because $\left\{ d_i^{\psi} \right\}_{i=1}^n$ is an ONB of $T_p \mathcal{M}$ for each p in the domain of that basis. Thus

$$N_{ij}^{\psi\varphi}M_{ik}^{\psi\varphi} = \langle d_j^{\varphi}, d_k^{\varphi} \rangle$$

and again using the fact that $\{d_i^{\varphi}\}_{i=1}^n$ is a basis one obtains the proper result. The other result is obtained by repeating the argument with $\varphi \leftrightarrow \psi$.

2 Corollary. We have $d_l^{\varphi}\left(N_{ij}^{\psi\varphi}\right)M_{ik}^{\psi\varphi} = -N_{ij}^{\psi\varphi}d_l^{\varphi}\left(M_{ik}^{\psi\varphi}\right)$.

Proof. Apply d_l^{φ} on the foregoing equation. Since δ_{ik} is a constant scalar function, we get zero on the left hand side (as a tangent vector working on any scalar function is zero). On the right we use the Leibniz property of d_l^{φ} .

3 Claim. We have $d_k^{\varphi}\left(M_{ii'}^{\psi\varphi}\right) = d_{i'}^{\varphi}\left(M_{ik}^{\psi\varphi}\right)$.

Proof. If we expand out the definitions we will find that this boils down to the fact that $[d_i^{\varphi}, d_k^{\varphi}] = 0$, which is always true for basis tangent vectors which correspond to charts, which is what d_i^{φ} is. Indeed,

$$M_{ii', k} - M_{ik, i'} \equiv d_k^{\varphi} \left(M_{ii'} \right) - d_{i'}^{\varphi} \left(M_{ik} \right)$$

= $d_k^{\varphi} \left(d_{i'}^{\varphi} \left(\psi_i \right) \right) - d_{i'}^{\varphi} \left(d_k^{\varphi} \left(\psi_i \right) \right)$
= $\left[d_k^{\varphi}, d_{i'}^{\varphi} \right] \left(\psi_i \right)$

and $\left[d_i^{\varphi}, d_j^{\varphi}\right] = 0$ because

$$\begin{split} \left(\begin{bmatrix} d_i^{\varphi}, d_j^{\varphi} \end{bmatrix} \right) (f) &\equiv d_i^{\varphi} d_j^{\varphi} f - (i \leftrightarrow j) \\ &= \begin{bmatrix} \partial_i \left(d_j^{\varphi} f \circ \varphi^{-1} \right) \end{bmatrix} \circ \varphi - (i \leftrightarrow j) \\ &= \begin{bmatrix} \partial_i \left(\begin{bmatrix} \partial_j \left(f \circ \varphi^{-1} \right) \end{bmatrix} \circ \varphi \circ \varphi^{-1} \right) \end{bmatrix} \circ \varphi - (i \leftrightarrow j) \\ &= \begin{bmatrix} \partial_i \left(\partial_j \left(f \circ \varphi^{-1} \right) \right) \end{bmatrix} \circ \varphi - (i \leftrightarrow j) \\ &= 0 \end{split}$$

as $\partial_i \partial_j = \partial_j \partial_i$.

3 Some short hand notation to make the calculation lighter

From this point onwards, since the charts φ and ψ are fixed, we omit them from the notation. Thus φ is considered the "original" chart and ψ the "new" chart. Consequently, all expansion coefficients in the original chart φ will have φ simply dropped expansion coefficients in the new chart ψ will be denoted by a bar above. We also abbreviate $M_{ij}^{\psi\varphi}$ simply as M_{ij} and the same for N. Finally we also abbreviate $d_i^{\varphi}(O) \equiv O_{,i}$ for any object O (typically O is an expansion coefficient in φ or ψ carrying itself some indices, but the application of d_i^{φ} always will be noted with a comma after all other indices).

Hence the transformation law for a vector's expansion coefficients

$$\overline{X}_i = M_{ij}X_i$$

The transformation law for a dual vector's expansion coefficients

$$\overline{\mu}_i = N_{ij}\mu_i$$

Transformation law for a (1, 1) tensor's expansion coefficients

$$\overline{T}_{ij} = M_{ii'} N_{jj'} T_{i'j'}$$

Transformation law for a basis vector

 $\overline{d}_i = N_{ij}d_j$

In the exercise, we "define" the Lie derivative along a vector field X of the (1, 1) tensor T via its components as

$$(L_X T)_{ij} = T_{ij,k} X_k - T_{kj} X_{i,k} + T_{ik} X_{k,j}$$

To see how it transforms, we must see how its constituent parts transform:

$$\overline{X}_{i,j} \equiv \overline{d}_j \left(\overline{X}_i \right)$$

$$= N_{jj'} d_{j'} \left(M_{ii'} X_{i'} \right)$$

$$= N_{jj'} d_{j'} \left(M_{ii'} \right) X_{i'} + N_{jj'} M_{ii'} d_{j'} \left(X_{i'} \right)$$

$$= N_{jj'} M_{ii',j'} X_{i'} + N_{jj'} M_{ii'} X_{i',j'}$$

So that $\overline{X}_{i,j}$ does not transform like a (1, 1) tensor, due to the extra first term (the second term alone is how it should have transformed had it been a (1, 1) tensor).

We have also

$$\overline{T}_{ij,k} \equiv \overline{d}_k \left(\overline{T}_{ij}\right)$$

$$= N_{kk'} d_{k'} \left(M_{ii'} N_{jj'} T_{i'j'}\right)$$

$$= N_{kk'} M_{ii',k'} N_{jj'} T_{i'j'} + N_{kk'} M_{ii'} N_{jj',k'} T_{i'j'} + N_{kk'} M_{ii'} N_{jj'} T_{i'j',k'}$$

So that $\overline{T}_{ij,k}$ does not transform like a (1, 2) tensor, due to the extra first two terms (the third term alone is how a (1, 2) tensor should have transformed).

We check however the transformation law of $\overline{(L_X T)}_{ij}$:

$$\begin{split} (L_X T)_{ij} &= \overline{T}_{ij,k} \overline{X}_k - \overline{T}_{kj} \overline{X}_{i,k} + \overline{T}_{ik} \overline{X}_{k,j} \\ &= (N_{kk'} M_{ii',k'} N_{jj'} T_{i'j'} + N_{kk'} M_{ii'} N_{jj',k'} T_{i'j'} + N_{kk'} M_{ii'} N_{jj'} T_{i'j',k'}) M_{kk''} X_{k''} \\ &- M_{kk'} N_{jj'} T_{k'j'} (N_{kk''} M_{ii',k''} X_{i'} + N_{kk''} M_{ii'} X_{i',k''}) \\ &+ M_{ii'} N_{kk'} T_{i'k'} (N_{jj'} M_{kk'',j'} X_{k''} + N_{jj'} M_{kk''} X_{k'',j'}) \\ &(\text{Regroup to terms containing derivatives of N and M and those that don't}) \\ &= N_{kk'} M_{ii',k'} N_{jj'} M_{kk''} T_{i'j'} X_{k''} + N_{kk'} M_{ii'} N_{jj',k'} M_{kk''} T_{i'j'} X_{k''} \\ &- M_{kk'} N_{jj'} N_{kk''} M_{ii',k''} T_{k'j'} X_{i'} + M_{ii'} N_{kk'} N_{jj'} M_{kk'''} T_{i'k'} X_{k''} \\ &- M_{kk'} N_{jj'} N_{kk''} M_{ii'} T_{k'j'} X_{i',k''} \\ &+ N_{kk'} M_{ii'} N_{jj'} M_{kk''} T_{i'j',k'} X_{k''} + M_{ii'} N_{kk'} N_{jj'} M_{kk''} T_{i'k'} X_{k'',j'} \\ \end{aligned}$$

We know what the answer *should* be:

$$\overline{(L_X T)}_{ij} \stackrel{?}{=} M_{ii'} N_{jj'} (L_X T)_{i'j'} = M_{ii'} N_{jj'} T_{i'j', k} X_k - M_{ii'} N_{jj'} T_{kj'} X_{i', k} + M_{ii'} N_{jj'} T_{i'k} X_{k, j'}$$

So we identify those terms in $\overline{(L_X T)}_{ij}$ as C (for "correct", the last two lines) and R (for "rest", the first two lines):

$$C := N_{kk'} M_{ii'} N_{jj'} M_{kk''} T_{i'j',k'} X_{k''} - M_{kk'} N_{jj'} N_{kk''} M_{ii'} T_{k'j'} X_{i',k''} + M_{ii'} N_{kk'} N_{jj'} M_{kk''} T_{i'k'} X_{k'',j'}$$

and

$$R := N_{kk'} M_{ii', k'} N_{jj'} M_{kk''} T_{i'j'} X_{k''} + N_{kk'} M_{ii'} N_{jj', k'} M_{kk''} T_{i'j'} X_{k''} - M_{kk'} N_{jj'} N_{kk''} M_{ii', k''} T_{k'j'} X_{i'} + M_{ii'} N_{kk'} N_{jj'} M_{kk'', j'} T_{i'k'} X_{k'}$$

We want to show that

$$C \stackrel{!}{=} M_{ii'} N_{jj'} T_{i'j',k} X_k - M_{ii'} N_{jj'} T_{kj'} X_{i',k} + M_{ii'} N_{jj'} T_{i'k} X_{k,j'}$$
(1)

and that R = 0.

We start with the first task. In order to do that we must we must "cancel out" factors of M and N. Take for instance the first term in C:

$$N_{kk'}M_{ii'}N_{jj'}M_{kk''}T_{i'j',k'}X_{k''} = (M_{ii'}N_{jj'}T_{i'j',k'})(N_{kk'}M_{kk''}X_{k''})$$

Using 1 we find for that term

$$\begin{split} N_{kk'}M_{ii'}N_{jj'}M_{kk''}T_{i'j',\,k'}X_{k''} &= (M_{ii'}N_{jj'}T_{i'j',\,k'}) \left(N_{kk'}M_{kk''}X_{k''}\right) \\ &= (M_{ii'}N_{jj'}T_{i'j',\,k'}) \left(\delta_{k'k''}X_{k''}\right) \\ &= (M_{ii'}N_{jj'}T_{i'j',\,k'}) X_{k'} \end{split}$$

so we get the first term on the RHS of (1) correctly. We proceed similarly using 1 twice more to find that (1) is correct. We go on to prove that R = 0: We use 1 three more times to find:

$$R \equiv N_{kk'}M_{ii', k'}N_{jj'}M_{kk''}T_{i'j'}X_{k''} + N_{kk'}M_{ii'}N_{jj', k'}M_{kk''}T_{i'j'}X_{k''} -M_{kk'}N_{jj'}N_{kk''}M_{ii', k''}T_{k'j'}X_{i'} + M_{ii'}N_{kk'}N_{jj'}M_{kk'', j'}T_{i'k'}X_{k''} = M_{ii', k}N_{jj'}T_{i'j'}X_k - N_{jj'}M_{ii', k}T_{kj'}X_{i'} +M_{ii'}N_{jj', k}T_{i'j'}X_k + M_{ii'}N_{kk'}N_{jj'}M_{kk'', j'}T_{i'k'}X_{k''}$$

for the first line, in the second term we relabel as $i' \leftrightarrow k$ to get

$$M_{ii',k}N_{jj'}T_{i'j'}X_k - N_{jj'}M_{ii',k}T_{kj'}X_{i'} = M_{ii',k}N_{jj'}T_{i'j'}X_k - N_{jj'}M_{ik,i'}T_{i'j'}X_k$$

But now use 3 so the first line of our most recent expression for R is zero.

We go on to the next line. We relabel in the second term $j' \leftrightarrow k'$ and $k \leftrightarrow k''$ to get

$$M_{ii'}N_{jj',k}T_{i'j'}X_k + M_{ii'}N_{kk'}N_{jj'}M_{kk'',j'}T_{i'k'}X_{k''} = M_{ii'}(N_{jj',k} + N_{k''j'}N_{jk'}M_{k''k,k'})T_{i'j'}X_k$$

Now we deal with the term $N_{k''j'}N_{jk'}M_{k''k,k'}$. In fact we can rewrite it as

$$\begin{split} N_{k''j'}N_{jk'}M_{k''k,\,k'} &= N_{k''j'}N_{jk'}M_{k''k',\,k} \\ &= -N_{k''j',\,k}N_{jk'}M_{k''k'} \\ &= -N_{k''j',\,k}\delta_{jk''} \\ &= -N_{jj',\,k} \end{split}$$

so that we really get zero. In the last expression, used again the fact that $M_{ij,k} = M_{ik,j}$ (in 3) as proven above already, as well as 2. The proof is finally complete.