# The Neumann Problem

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## 1 Formulation of the Problem

Let D be a bounded open subset in  $\mathbb{R}^d$  with  $\partial D$  its boundary such that D is sufficiently nice (to be stipulated later as Lipschitz). Let  $f \in L^2(D)$  and  $g: L^2(\partial D)$  be two given scalar fields and  $n: \partial D \to S^{d-1}$  be the normal unit vector to the boundary. Prove that

$$\begin{cases} \Delta \varphi &= f\\ (\nabla \varphi) \cdot n|_{\partial D} &= g \end{cases}$$
(1)

has a unique solution up to a constant for the unknown scalar field  $\varphi: D \to \mathbb{R}$ in  $H^1(D)$  if and only if

$$\int_{D} f = \int_{\partial D} g \tag{2}$$

(This last condition makes sense because  $L^2 \subseteq L^1$ )

#### 1.1 Sketch of Solution

- 1. Verify that if a solution of (1) exists, then (2) must be satisfied using the divergence theorem.
- 2. Formulate (1) as a variational problem:  $\varphi$  solves (1) iff

$$\int_{D} (\nabla \varphi) \cdot (\nabla \psi) = -\int_{D} f \psi + \int_{\partial D} g \psi \quad \forall \psi$$
(3)

Assuming (2) is satisfied.

3. Write (3) using the bilinear and linear respectively forms

$$\omega \left( \varphi, \psi \right) \;\; := \;\; \int_{D} \left( \nabla \varphi \right) \cdot \left( \nabla \psi \right)$$

and

$$\eta\left(\psi\right) := -\int_{D} f\psi + \int_{\partial D} g\psi$$

4. Use the Lax-Milgram theorem, which says that if  $\omega$  is continuous,  $\eta$  is continuous, and  $\omega$  is elliptic (meaning  $\omega(\psi, \psi) \ge \alpha \|\psi\|^2$  for all  $\psi$  for some  $\alpha > 0$ ) then there is a unique solution  $\varphi$  to the equation

$$\omega\left(\varphi,\,\cdot\right) - \eta \quad = \quad 0$$

In order to show that  $\omega$  and  $\eta$  are continuous, use the Cauchy-Schwarz inequality; in order to show that  $\omega$  is elliptic, use the Poincare inequality

$$\|\psi\| \leq C \|\nabla\psi\|$$

for some C > 0.

### 1.2 Solution

(We follow notes by Hervé Le Dret found on https://www.ljll.math.upmc.fr/~ledret/M1ApproxPDE.html) 1 Note. Regarding (2), we see that if  $\varphi$  solves (1), then using the divergence theorem we find

$$\int_{D} \Delta \varphi \equiv \int_{D} \nabla \cdot (\nabla \varphi)$$
(Div. thm.)
$$= \int_{\partial D} (\nabla \varphi) \cdot n$$

so that using  $\Delta \varphi = f$  and the boundary condition  $(\nabla \varphi) \cdot n|_{\partial D} = g$ , we arrive at the compatability condition (2). Conversely, if (2) does not hold then as seen, there cannot exist a solution.

First we introduce some notation:

**2 Definition.**  $\mathcal{D}(D)$  is the space of all infinitely differentiable scalar-fields on D such that which have compact support. Its dual,  $\mathcal{D}'(D)$ , is the space of all distributions on D. It is the space of all continuous linear forms  $\mathcal{D}(D) \to \mathbb{R}$ .

**3 Definition.** Cpt(D) is the set of all compact subsets of D.

#### 4 Definition. We define

$$L^{p}_{\text{loc}}(D) := \{ u: D \to \mathbb{R} \mid u|_{K} \in L^{p}(K) \,\forall K \in Cpt(D) \}$$

Note that  $L^{1}_{loc}(D)$  contains all  $L^{p}(D)$  spaces, and that  $L^{1}_{loc}(D)$  is continuously and injectively embedded in  $\mathcal{D}'(D)$ .

5 Claim. If  $u \in L^1_{loc}(D)$  is such that

$$\int_{D} u\varphi = 0 \quad \forall \varphi \in \mathcal{D}\left(D\right)$$

then u = 0 almost-everywhere.

**6 Definition.** Let  $m \in \mathbb{N}$  and  $p \in [1, \infty]$ . The Sobolev space  $W^{m, p}(D)$  is defined as

$$W^{m, p}(D) := \left\{ u \in L^{p}(D) \mid \partial^{\alpha} u \in L^{p}(D) \; \forall \alpha \in \mathbb{N}^{d} : |\alpha| \le m \right\}$$

where we are using the multi-index notation for  $\alpha$ . Note also that as  $u \in L^p(D)$ , it is not necessarily differentiable in the usual sense, but it is in a distribution  $L^p(D) \subseteq L^1_{\text{loc}}(D) \subseteq \mathcal{D}'(D)$ ; as distributions are infinitely differentiable,  $\partial^{\alpha} u$ makes sense. Then the requirement is that  $\partial^{\alpha} u$  is a distribution that comes from a function in  $L^p(D)$ . We also define for convenience

$$H^m\left(D\right) := W^{m,2}\left(D\right)$$

and note  $W^{0, p}(D) \equiv L^{p}(D)$ . We also have a natural norm:

$$\|u\|_{W^{m,p}(D)} := \left(\sum_{a \in \mathbb{N}^d : |\alpha| \le m} \left( \|\partial^{\alpha} u\|_{L^p(\Omega)} \right)^p \right)^{\frac{1}{p}}$$

for all  $p \in [1, \infty)$  and

$$\|u\|_{W^{m,\infty}(D)} := \max_{\alpha \in \mathbb{N}^d : |\alpha| \le m} \|\partial^{\alpha} u\|_{L^{\infty}(D)}$$

Note that for p = 2 we can see that  $\|\cdot\|_{H^m(D)}$  comes from the inner product

$$\langle u, v \rangle_{H^m(D)} := \sum_{\alpha \in \mathbb{N}^d : |\alpha| \le m} \langle \partial^{\alpha} u, \partial^{\alpha} v \rangle_{L^2(\Omega)}$$

as  $L^{2}(\Omega)$  is a Hilbert space.

Note that  $W^{m, p}(D)$  are Banach spaces and so  $H^{m}(D)$  is a Hilbert space. For example, the step function H is in  $L^{1}$  but is its (distributional) derivative, the delta function  $\delta_{0}$  is not a function.

**7 Definition.** The closure of  $\mathcal{D}(D)$  in  $H^m(D)$  is denoted by  $H_0^m(D)$ . This is a sub-Hilbert-space in  $H^m(D)$ .

**8 Definition.** A bounded open subset  $C \subseteq \mathbb{R}^d$  is called Lipschitz if its boundary is "sufficiently regular" in the sense that it can be thought of as locally being the graph of a Lipschitz continuous function.

9 Claim. If D is Lipschitz then  $\exists !\gamma_0 : H^1(D) \to L^2(\partial D)$  continuous linear such that for all  $u \in C^1(\overline{D})$ ,  $\gamma_0(u|_D) = u|_{\partial D}$ . There is also a well defined continuous linear mapping  $\gamma_1 : H^2(D) \to L^2(\partial D)$  given by

$$\gamma_1(u) := \sum_{i=1}^a \gamma_0(\partial_i u) n_i$$

and such that for all  $u \in C^2(\overline{D})$ ,

$$\gamma_1\left(\left.u\right|_D\right) = (\nabla u) \cdot n$$

where n is the normal unit vector to  $\partial D$ .

**10 Definition.** Define  $\mathcal{H} := \{ \psi \in H^1(D) \mid \int_D \psi = 0 \}$  which makes sense as  $H^1(D) \subseteq L^2(D) \subseteq L^1(D)$ .

11 Claim.  $\mathcal{H}$  is a Hilbert space using the same scalar product as that of  $H^{1}(\Omega)$ .

*Proof.* We show that  $\mathcal{H} \in Closed(H^1(\Omega))$ . Let  $\{h_n\}_n$  be a sequence in  $\mathcal{H}$  such that  $h_n \to h$  in  $H^1(D)$  for some  $h \in H^1(D)$ . Want  $h \in \mathcal{H}$ . Then of course as  $H^1(D) \subseteq L^2(D)$ , we have  $h_n \to h$  in  $L^2(D)$ , and then by Cauchy-Schwarz in  $L^1(D)$  as well. Thus, since each  $h_n \in \mathcal{H}$ , its integral is zero and so

$$\begin{array}{rcl} 0 & = & \int_D h_r \\ & \rightarrow & \int_D h \end{array}$$

so that  $h \in \mathcal{H}$  as well.

12 Claim. (Variational formulation of the Neumann problem)  $\varphi \in H^2(D)$  solves the Neumann problem above iff for any  $\psi \in \mathcal{H}$ ,

$$\int_{D} \left( \nabla \varphi \right) \cdot \left( \nabla \psi \right) \quad = \quad - \int_{D} f \psi + \int_{\partial D} g \gamma_{0} \left( \psi \right)$$

Assuming f and g obey the compatability condition above.

*Proof.* Take an arbitrary  $\psi \in \mathcal{H}$  and multiply the equation  $\Delta \varphi = f$  with it to get

$$(\Delta \varphi) \psi = f \psi$$

Note that since  $\varphi \in H^2(D)$ ,  $\Delta \varphi \in L^2(D)$ ; also,  $\psi \in H^1(D)$  implies  $\psi \in L^2(D)$ . Then via Hoelder's inequality that  $(\Delta \varphi) \psi \in L^1(D)$  so that the left hand side is integrable. Since  $f \in L^2(D)$  and  $\psi \in L^2(D)$ , again by Hoelder  $f\psi \in L^1(D)$  so that we can integrate the equation and obtain:

$$\int_D \left( \Delta \varphi \right) \psi \quad = \quad \int_D f \psi$$

We now use Green's first identity on the left hand side to get

$$\begin{split} \int_{D} \left( \Delta \varphi \right) \psi &= -\int_{D} \left( \nabla \varphi \right) \cdot \left( \nabla \psi \right) + \int_{\partial D} \psi \left( \nabla \varphi \right) \cdot n \\ &= -\int_{D} \left( \nabla \varphi \right) \cdot \left( \nabla \psi \right) + \int_{\partial D} \psi g \end{split}$$

Of course this cannot really be written since  $\varphi$  is not a map on  $\partial D$  but only on D so that we must use 9 and then Green's first identity is written as

$$\int_{D} (\Delta \varphi) \psi = - \int_{D} (\nabla \varphi) \cdot (\nabla \psi) + \int_{\partial D} \gamma_{0} (\psi) \gamma_{1} (\varphi)$$

So that we find using the boundary condition that  $\varphi$  fulfills:

$$\int_{D} (\Delta \varphi) \psi = - \int_{D} (\nabla \varphi) \cdot (\nabla \psi) + \int_{\partial D} \gamma_{0} (\psi) g$$

We find

$$\int_{D} \left( \nabla \varphi \right) \cdot \left( \nabla \psi \right) \quad = \quad \int_{\partial D} g \gamma_0 \left( \psi \right) - \int_{D} f \psi$$

which is what we wanted to show.

Conversely, if we have some  $\varphi \in H^{2}(D)$  such that

$$\int_{D} (\nabla \varphi) \cdot (\nabla \psi) = \int_{\partial D} g \gamma_0(\psi) - \int_{D} f \psi \quad \forall \psi \in \mathcal{H}$$
(4)

Now because  $\mathcal{D}(D)$  is not actually contained within  $\mathcal{H}$ , we need a little song and dance about defining, for each  $\psi \in \mathcal{D}(D)$ ,

$$\tilde{\psi} := \psi - \frac{1}{\int_D} \int_D \psi$$

and now  $\tilde{\psi} \in \mathcal{H}$ . Note  $\psi$  and  $\tilde{\psi}$  differ by a constant, namely,  $\frac{1}{\int_D} \int_D \psi =: k$ ,

so that  $\nabla \psi = \nabla \tilde{\psi}$ . Hence if  $\psi \in \mathcal{D}(D)$ ,

$$\begin{split} \int_{D} (\nabla \varphi) \cdot (\nabla \psi) &= \int_{D} (\nabla \varphi) \cdot \left( \nabla \tilde{\psi} \right) \\ & \text{(By hypothesis)} \\ &= \int_{\partial D} g \gamma_0 \left( \tilde{\psi} \right) - \int_{D} f \tilde{\psi} \\ &= \int_{\partial D} g \gamma_0 \left( \psi - k \right) - \int_{D} f \left( \psi - k \right) \\ & \left( \psi \in \mathcal{D} \left( D \right) \Longrightarrow \gamma_0 \left( \psi \right) = 0 \land \gamma_0 \left( k \right) = k \right) \\ &= -\int_{D} f \psi - k \left( \int_{D} f - \int_{\partial D} g \right) \\ & \text{(Using the compatability condition)} \\ &= -\int_{D} f \psi \end{split}$$

Since this holds for all  $\psi \in \mathcal{D}(D)$ , we can use (5) to conclude  $\Delta \varphi = f$  in the distributional sense. But  $f \in L^2(D)$ , so this holds in  $L^2(D)$  as well.

So now we need to establish that  $\varphi$  obeys the boundary conditions.

Using Green's formula now on the left-hand side of (4) again we find

$$-\int_{D} \left(\Delta\varphi\right)\psi + \int_{\partial D} \gamma_{0}\left(\psi\right)\gamma_{1}\left(\varphi\right) = \int_{\partial D} g\gamma_{0}\left(\psi\right) - \int_{D} f\psi \qquad \forall \psi \in \mathcal{H}$$

But now we may use  $\Delta \varphi = f$  to find

$$\int_{\partial D} \gamma_0 \left( \psi \right) \left( \gamma_1 \left( \varphi \right) - g \right) = 0 \qquad \forall \psi \in \mathcal{H}$$

If  $g \in H^{\frac{1}{2}}(D)$  then via  $\varphi \in H^{2}(D)$ ,  $\gamma_{1}(\varphi) \in H^{\frac{1}{2}}(\partial D)$ , so that there is some  $\psi \in \mathcal{H}$  so that  $\gamma_{0}(\psi) = \gamma_{1}(\varphi) - g$  and we find

$$\int_{\partial D} \left( \gamma_1 \left( \varphi \right) - g \right)^2 = 0$$

so it must be that  $\gamma_1(\varphi) - g = 0$  and  $\varphi$  obeys the Neumann boundary solution of (1) as needed.

13 Note. Defining the bilinear form  $\omega_{\Delta} : \mathcal{H} \times \mathcal{H} \to \mathbb{R}$  by

$$\omega_{\Delta}(\varphi, \psi) := \int_{D} (\nabla \varphi) \cdot (\nabla \psi) \qquad \forall (\varphi, \psi) \in \mathcal{H} \times \mathcal{H}$$

and a linear form  $\eta_{f,g}: \mathcal{H} \to \mathbb{R}$  via

$$\eta_{f,g}(\psi) := -\int_{D} f\psi + \int_{\partial D} g\gamma_{0}(\psi) \qquad \forall \psi \in \mathcal{H}$$

We see via that  $\varphi$  solves (1) iff

$$\omega_{\Delta}(\varphi, \cdot) = \eta_{f,g}$$

14 Claim. (Lax-Milgram) Let  $\mathcal H$  be a Hilbert space,  $\omega$  be a bilinear form and  $\eta$  a linear form, such that:

1.  $\omega$  is continuous:  $\exists M > 0$  such that

$$|\omega(\varphi, \psi)| \leq M \|\varphi\| \|\psi\| \quad \forall (\varphi, \psi) \in \mathcal{H}^2$$

2.  $\omega$  is  $\mathcal{H}$ -elliptic:  $\exists \alpha > 0$  such that

$$\omega\left(\psi,\,\psi\right) \geq \alpha \|\psi\|^2 \qquad \forall v \in \mathcal{H}$$

3.  $\eta$  is continuous:  $\exists C > 0$  such that

$$\left|\eta\left(\psi\right)\right| \leq C \left\|\psi\right\| \quad \forall v \in \mathcal{H}$$

Then  $\exists ! \varphi \in \mathcal{H}$  such that

$$\omega\left(\varphi,\,\cdot\right) \quad = \quad \eta \tag{5}$$

*Proof.* We start with uniqueness: Let  $\varphi_1$  and  $\varphi_2$  both satisfy (5). Using linearity of  $\omega$  in its first argument we have

$$\omega\left(\varphi_1 - \varphi_2, \cdot\right) = 0$$

In particular,

$$\omega \left( \varphi_1 - \varphi_2, \, \varphi_1 - \varphi_2 \right) = 0$$

Now using the fact that  $\omega$  is  $\mathcal{H}$ -elliptic we have actually that

$$\alpha \|\varphi_1 - \varphi_2\|^2 \le 0$$

and so  $\|\varphi_1 - \varphi_2\| = 0$  as  $\alpha > 0$ . But a norm is zero iff its argument is zero, so that  $\varphi_1 = \varphi_2$ .

We turn to existence:

Define  $\Omega : \mathcal{H} \to \mathcal{H}'$  by  $\psi \mapsto \omega(\psi, \cdot)$ . Then the variational problem is to find a  $\varphi$  such that

$$\Omega\left(\varphi\right) = \eta$$

Since  $\omega$  is continuous,  $\Omega(\varphi)$  is continuous (and also linear by bilinearity) so that  $\Omega$  is well defined.  $\eta$  is also continuous so that the variational problem is in fact an equation to be solve in  $\mathcal{H}'$ , the dual of  $\mathcal{H}$ .

Since  $\eta$  is given and  $\varphi$  is the unknown, the question is whether  $\Omega$  is an epimorphism.

Claim.  $im(\Omega) \in Closed(\mathcal{H}').$ 

*Proof.* Let  $\{\pi_n\}_n$  be a sequence in  $im(\Omega)$  such that  $\pi_n \to \pi$  for some  $\pi \in \mathcal{H}'$ . If we can show that  $\pi \in im(\Omega)$  then our result is implied.

Note that since  $\pi_n$  converges, it is Cauchy. Since it is in  $im(\Omega)$ , we have a sequence  $\{\psi_n\}_n$  in  $\mathcal{H}$  such that  $\Omega(\psi_n) = \pi_n$  for all n. By  $\mathcal{H}$ -ellipticity, we have

$$\begin{aligned} \left\|\psi_{n}-\psi_{m}\right\|^{2} &\leq \frac{1}{\alpha}\omega\left(\psi_{n}-\psi_{m},\,\psi_{n}-\psi_{m}\right) \\ &= \frac{1}{\alpha}\left\langle\Omega\left(\psi_{n}\right)-\Omega\left(\psi_{m}\right),\,\psi_{n}-\psi_{m}\right\rangle_{\mathcal{H}',\,\mathcal{H}} \\ &= \frac{1}{\alpha}\left\langle\pi_{n}-\pi_{m},\,\psi_{n}-\psi_{m}\right\rangle_{\mathcal{H}',\,\mathcal{H}} \\ &\quad \text{(Cauchy-Schwarz)} \\ &\leq \frac{1}{\alpha}\left\|\pi_{n}-\pi_{m}\right\|\left\|\psi_{n}-\psi_{m}\right\| \end{aligned}$$

Thus if  $\|\psi_n - \psi_m\| = 0$  we are finished, as then that means  $\{\psi_n\}_n$  converges. Otherwise, we have

$$\|\psi_n - \psi_m\| \leq \frac{1}{\alpha} \|\pi_n - \pi_m\|$$

so that  $\{\psi_n\}_n$  is Cauchy, and by completeness of  $\mathcal{H}$ , converges. So that there is some  $\psi \in \mathcal{H}$  such that  $\psi_n \to \psi$ . But  $\Omega$  is continuous, so that

$$\pi = \lim_{n} \pi_{n}$$
$$= \lim_{n} \Omega(\psi_{n})$$
$$= \Omega\left(\lim_{n} \psi_{n}\right)$$
$$= \Omega(\psi)$$

and we find  $\pi \in im(\Omega)$  as desired.

Claim.  $\overline{im(\Omega)} = \mathcal{H}'$  (density)

*Proof.* We show this by showing that  $(im(\Omega))^{\perp} = \{0\}$ . Let  $\pi \in (im(\Omega))^{\perp}$ . Then

$$\langle \Omega (\psi), \pi \rangle_{\mathcal{H}'} = 0$$

for all  $\psi \in \mathcal{H}$ . If  $\delta : \mathcal{H}' \to \mathcal{H}$  is the isomorphism furnished by the Riesz representation theorem, then

$$\langle \Omega (\psi) , \pi \rangle_{\mathcal{H}'} = \langle \omega (\psi, \cdot) , \pi \rangle_{\mathcal{H}'} = \omega (\psi, \delta (\pi))$$

so that

$$\omega\left(\psi,\,\delta\left(\pi\right)\right) = 0$$

for all  $\psi \in \mathcal{H}$ . So pick  $\psi = \delta(\pi)$  to get

$$0 = \omega \left( \delta \left( \pi \right), \, \delta \left( \pi \right) \right) \\ \geq \alpha \| \delta \left( \pi \right) \|^{2}$$

by  $\mathcal{H}$ -ellipticity. But  $\alpha > 0$  so that  $\delta(\pi) = 0$ , hence  $\pi = 0$ . But  $\pi$  was arbitrary, so that  $(im(\Omega))^{\perp} = \{0\}$ .

We then have as an immediate result that  $\Omega$  is an epimorphism.

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15 Claim. (Poincare-Wirtinger inequality) Let D be a Lipschitz open subset of  $\mathbb{R}^{d}$ . Then there exists a constant C depending on D such that for all  $\psi \in H^{1}(D)$ ,

$$\left\|\psi - \frac{1}{\int_D} \int_D \psi\right\|_{L^2(D)} \leq C \|\nabla \psi\|_{L^2(D)}$$
(6)

Proof. Omitted.

16 Claim. (The Lax-Milgram theorem may be used) The conditions of 14 are fulfilled by  $\omega_{\Delta}$  and  $\eta_{f,g}$ .

*Proof.* We first show the elipticity of  $\omega_{\Delta}$ :

$$\omega_{\Delta}(\psi, \psi) \equiv \int_{D} (\nabla \psi) \cdot (\nabla \psi)$$
$$\equiv \|\nabla \psi\|_{L^{2}(D)}^{2}$$

Because  $\psi \in \mathcal{H}$ ,  $\int_D \psi = 0$  so that (6) implies

$$\|\psi\|_{L^2(D)}^2 \leq C^2 \|\nabla\psi\|_{L^2(D)}^2$$

Hence

$$\begin{aligned} \|\psi\|_{\mathcal{H}}^{2} &\equiv \|\psi\|_{H^{1}(D)}^{2} \\ &\equiv \|\psi\|_{L^{2}(D)}^{2} + \|\nabla\psi\|_{L^{2}(D)}^{2} \\ &\leq (1+C^{2}) \|\nabla\psi\|_{L^{2}(D)}^{2} \\ &= (1+C^{2}) \omega_{\Delta}(\psi,\psi) \end{aligned}$$

So that  $\omega_{\Delta}$  is  $\mathcal{H}$ -elliptic with constnat  $\alpha := (1 + C^2)^{-1}$ . We now show continuity of  $\omega_{\Delta}$ :

$$\begin{aligned} |\omega_{\Delta}(u, v)| &\equiv \left| \int_{D} (\nabla u) \cdot (\nabla v) \right| \\ &\leq \int_{D} |(\nabla u) \cdot (\nabla v)| \\ &\quad (\text{Cauchy-Schwarz}) \\ &\leq \|\nabla u\|_{L^{2}(D)} \|\nabla v\|_{L^{2}(D)} \\ &\leq \|u\|_{H^{1}(D)} \|v\|_{H^{1}(D)} \end{aligned}$$

For  $\eta_{f,g}$ , we have

$$\begin{aligned} |\eta_{f,g}(\psi)| &\equiv \left| -\int_{D} f\psi + \int_{\partial D} g\gamma_{0}(\psi) \right| \\ & (\text{Cauchy-Schwarz}) \\ &\leq \| f\|_{L^{2}(D)} \|\psi\|_{L^{2}(D)} + \|g\|_{L^{2}(\partial D)} \|\gamma_{0}(\psi)\|_{L^{2}(\partial D)} \\ & (\gamma_{0} \text{ is continuous, so for some constant } c > 0) \\ &\leq \| f\|_{L^{2}(D)} \|\psi\|_{H^{1}(D)} + \|g\|_{L^{2}(\partial D)} c\|\psi\|_{H^{1}(D)} \end{aligned}$$

17 Remark.  $\mathcal{H}$  is a subspace of  $H^{1}(D)$  which is  $L^{2}(D)$ -orthogonal to the constant maps: If c is a constant map, and  $\psi \in \mathcal{H}$ :

$$\langle \psi, c \rangle_{L^{2}(D)} \equiv \int_{D} \psi c$$

$$= c \int_{D} \psi$$

$$(\psi \in \mathcal{H})$$

$$= 0$$

and since the constants maps are also solutions of (1), we conclude that the general solution of (1) is taken from  $\mathcal{H} \oplus \mathbb{R}$ .