

# CONFORMAL INVARIANTS ASSOCIATED TO A MEASURE, I: POINTWISE INVARIANTS

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ABSTRACT. In this paper we study Riemannian manifolds  $(M^n, g)$  equipped with a smooth measure  $m$ . In particular, we show that Riemannian invariants of  $(M^n, g)$  give rise to conformal densities of the *Riemannian measure space*  $(M^n, g, m)$ . This leads to a natural definition of the Ricci and scalar curvatures of *RM-spaces*, both of which are conformally invariant. We also study some natural variational integrals.

## 1. INTRODUCTION

By a *Riemannian measure space* (or *RM-space*) we will mean a triple  $(M^n, g, m)$  consisting of a smooth oriented manifold  $M^n$ , a Riemannian metric  $g$ , and a smooth measure  $m$  defined on  $M^n$ . In this paper we show that every Riemannian invariant of  $(M^n, g)$  gives rise to a conformal density of the *RM-space*  $(M^n, g, m)$ . In particular, we define conformally invariant notions of the Riemannian, Ricci, and scalar curvature associated to  $(M^n, g, m)$ . When  $m$  is the Riemannian measure of  $g$  the conformally invariant curvatures agree with their classical counterparts (see Section 3).

In the Riemannian setting, the construction of conformal invariants is an important problem with connections to many fields in mathematics and physics; see the articles of Fefferman-Graham ([FG85], [FG02]) and Graham-Hirachi ([GH05]), for example. As we shall see, the additional structure of a measure gives rise to many conformal invariants which are not, in general, Riemannian invariants.

The dependence of our invariants on the measure  $m$  is mediated by the density function  $f$ , defined by

$$(1.1) \quad dm = e^{-f} dVol(g).$$

In particular the conformally invariant curvatures are local expressions in  $g$  and  $f$ . For example, the conformally invariant scalar curvature is defined by

$$\begin{aligned} R_n^m(g) &= e^{-\frac{2}{n}f} R(e^{-\frac{2}{n}f}g) \\ &= R(g) + \frac{2(n-1)}{n} \Delta_g f - \frac{(n-1)(n-2)}{n^2} |\nabla f|^2, \end{aligned}$$

where  $R(g)$  is the scalar curvature of  $g$ . It is invariant in the sense that  $R_n^m(e^{2w}g) = e^{-2w} R_n^m(g)$ .

Geometric quantities associated a metric and measure are of course not new, and go back at least to the work of Bakry-Émery [BE85], who introduced a notion of the

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Ricci curvature in this setting. In Perelman's recent work on the Ricci flow ([Per02]) he defined the scalar curvature associated to the  $BE$ -Ricci tensor. As we explain below, the Ricci curvature of Bakry-Émery and the associated scalar curvature defined by Perelman can be viewed as the infinite-dimensional versions of the conformally invariant Ricci and scalar curvature.

The diffeomorphism group  $\mathcal{D}$  of the manifold acts on an  $RM$ -space via pull-back:  $\varphi \in \mathcal{D} \Rightarrow \varphi^*(M^n, g, m) = (M^n, \varphi^*g, \varphi^*m)$ . This action gives rise to a notion of  $RM$ -invariance:  $T = T(g, m)$  is an  $RM$ -invariant if  $T(\varphi^*g, \varphi^*m) = \varphi^*T(g, m)$ . The conformally invariant curvatures are  $RM$ -invariants, as are the Ricci curvature of Bakry-Émery and the scalar curvature of Perelman. However, if we allow  $\mathcal{D}$  to act on the metric alone (keeping the measure  $m$  fixed), then in general  $T(\varphi^*g, m) \neq \varphi^*T(g, m)$ . This dependence on the diffeo-class of a metric leads to a fundamental result (Theorem 5.2): the behavior of a Riemannian invariant under conformal deformations of the metric is equivalent to the action of  $\mathcal{D}$  (via conjugation) on the associated density.

Using this equivalence we study a natural class of variational integrals defined by scalar conformal densities. An example of particular interest is the (weighted) integral of the conformally invariant scalar curvature:

$$\mathcal{S}^m[g] = \int R_n^m(g) e^{-\frac{2}{n}f} dm.$$

Viewed as a functional on the space of metrics (with  $m$  fixed), critical points of  $\mathcal{S}^m$  are conformal to an Einstein metric. We can also consider the constrained problem of restricting  $\mathcal{S}^m$  to the orbit of  $g$  under the action of  $\mathcal{D}$ . By Theorem 5.2, this is equivalent to the Yamabe problem.

Many of the results of this paper were announced in [CGY06]. In the sequel to this paper [CGY07] we will construct conformally covariant operators on  $RM$ -spaces, based on the methods of Fefferman-Graham ([FG85]). These are analogues of the  $GJMS$ -operators ([GJMS92]) defined on Riemannian manifolds. They also give rise to a family of "scalar curvatures" associated to an  $RM$ -space, and the conformally invariant curvatures of this paper play an important role in the construction.

## 2. $RM$ -INVARIANTS

Let  $(M^n, g)$  be a Riemannian manifold of dimension  $n \geq 2$ . We denote the Riemannian curvature tensor by  $Riem$ , the Ricci tensor by  $Ric$  (or  $Ric(g)$ ), and the scalar curvature by  $R$  (or  $R(g)$ ). Finally, we denote the volume form of  $g$  by  $dVol(g)$ .

Let  $\mathcal{G} \subset S^2T^*M^n$  denote the ray bundle consisting of metrics in the conformal class of  $g$ ; we also denote the conformal class by  $[g] = \{\hat{g} = e^{2w}g \mid w \in C^\infty(M^n)\}$ . If  $\delta_s : \mathcal{G} \rightarrow \mathcal{G}$  denotes the dilations  $\delta_s(g) = s^2g$ , then sections of  $\mathcal{G}$  which are homogeneous of degree  $\beta$  with respect to  $\delta_s$  are called *conformal densities of weight  $\beta$* .

Suppose  $m$  is a smooth measure defined on  $M^n$ . We define the *density function associated to  $m$*  by

$$(2.1) \quad dm = e^{-f} dVol(g).$$

That is,  $f$  is the logarithm of the Radon-Nikodym derivative of  $dVol(g)$  with respect to  $dm$ . Given a conformal metric  $\hat{g} = e^{2w}g$ , let  $\hat{f}$  denote the density function associated to  $\hat{g}$ ; i.e.,

$$(2.2) \quad dm = e^{-\hat{f}} dVol(\hat{g}).$$

Since  $dVol(\hat{g}) = e^{nw} dVol(g)$ , it follows that

$$(2.3) \quad \hat{f} = f + nw.$$

The measure  $m$  determines a canonical choice of conformal metric whose density function is zero; we call this the *canonical base metric*. The precise formulation of this fact is given by the following lemma:

**Lemma 2.1.** *Let  $(M^n, g, m)$  be an RM-space. Then there is a unique metric  $g_m \in [g]$  with the property*

$$(2.4) \quad dm = dVol(g_m).$$

We call  $g_m$  the canonical base metric in  $[g]$ .

*Proof.* Let  $f$  be the density function associated to  $g$ :

$$dm = e^{-f} dVol(g).$$

If we take

$$(2.5) \quad g_m = e^{-\frac{2}{n}f} g,$$

then from formula (2.3) it follows that  $dm = dVol(g_m)$ . The uniqueness of  $g_m$  is clear.  $\square$

As we noted in the introduction, given an RM-space the diffeomorphism group  $\mathcal{D}$  of the manifold acts on  $(M^n, g, m)$  by pull-back:

$$(2.6) \quad \varphi^*(M^n, g, m) = (M^n, \varphi^*g, \varphi^*m).$$

This leads to a natural notion of invariance:

**Definition 2.1.** *A covariant tensor field  $T = T(g, m)$  on  $M^n$  is a local RM-invariant if for every diffeomorphism  $\varphi : M^n \rightarrow M^n$ ,*

$$(2.7) \quad T(\varphi^*g, \varphi^*m) = \varphi^*T(g, m).$$

We will usually omit the adjective "local", unless the context requires it. We can also define RM-invariance for contravariant tensors:

$$T(\varphi^*g, \varphi^*m) = \varphi_*T(g, m),$$

and thus extend the notion to tensor fields of all types.

**Example 1.** The density function  $f$  associated to a measure, along with its covariant derivatives, is an example of a local RM-invariant:

**Lemma 2.2.** *Let  $(M^n, g, m)$  be a Riemannian measure space, and let  $f$  be given by*

$$(2.8) \quad dm = e^{-f} dVol(g).$$

*Then  $f$  is an  $RM$ -invariant. In addition,  $\nabla^k f$ , all tensor products of the form*

$$(2.9) \quad \sum_{I=(i_1, \dots, i_p)} c_I \nabla^{i_1} f \otimes \cdots \otimes \nabla^{i_p} f,$$

*and all contractions of (2.9) are  $RM$ -invariants.*

*Proof.* Given a diffeomorphism  $\varphi : M^n \rightarrow M^n$ , write

$$(2.10) \quad \varphi^* dm = e^{-f_\varphi} dVol(\varphi^* g).$$

Comparing (2.8) and (2.10) and using the invariance of the Riemannian volume form under pull-back we conclude

$$(2.11) \quad f_\varphi = \varphi^* f = f \circ \varphi.$$

Thus,  $f$  is an  $RM$ -invariant. From this fact it follows that all covariant derivatives of  $f$ , and all tensor powers of its covariant derivatives along with their contractions, are also  $RM$ -invariants. □

**Example 2.** The Bakry-Émery Ricci tensor and the scalar curvature introduced by Perelman are both examples of  $RM$ -invariants. This follows from Lemma 2.2, since the definitions are

$$(2.12) \quad Ric^m(g) = Ric(g) + \nabla^2 f, \quad (\text{Bakry-Émery})$$

$$R^m(g) = R(g) + 2\Delta f - |\nabla f|^2. \quad (\text{Perelman})$$

**Example 3.** Given a Riemannian invariant, one can construct various  $RM$ -invariants through multiplication by a weight. More precisely, the (multiplicative) group of smooth positive functions  $\mathcal{C}_+ = \{\gamma \in C^\infty(M^n) \mid \gamma > 0\}$  acts on a metric by conformal deformation:

$$(2.13) \quad \gamma : g \mapsto \gamma \cdot g.$$

Of course,  $\mathcal{C}_+$  also acts on sections of the tensor bundles  $\otimes^k T^* M^n \otimes^\ell T M^n$  by multiplication. Suppose  $I = I(g)$  is a Riemannian invariant. To simplify matters we assume  $I$  is a scalar function or more generally a covariant tensor field; thus

$$(2.14) \quad I(\varphi^* g) = \varphi^* I(g).$$

**Definition 2.2.** *Given a pair of real numbers  $(\alpha, \beta)$ , let*

$$(2.15) \quad \mathcal{I}_{\alpha, \beta}^m(g) = e^{\beta f} I(e^{\alpha f} g)$$

That is,  $\mathcal{I}_{\alpha, \beta}^m$  is obtained from  $I$  through 'conjugation' by the action of  $\mathcal{C}_+$ . It follows from Lemma 2.2 that  $\mathcal{I}_{\alpha, \beta}^m$  is an  $RM$ -invariant. As we shall see, an example of particular importance is when  $\alpha = -2/n$  and  $\beta = 0$ .

**2.1. Conformal densities.** An  $RM$ -conformal density is a conformal density (in the usual sense) which is also an  $RM$ -invariant. Using the construction of  $RM$ -invariants in Example 3 above, we now show that every Riemannian invariant induces such a conformal density, of any specified weight:

**Proposition 2.1.** *Let  $(M^n, g, m)$  be a Riemannian measure space with density function  $f$ :*

$$(2.16) \quad dm = e^{-f} dVol(g).$$

*Suppose  $I = I(g)$  is Riemannian invariant. Then for any real number  $\beta$ , the  $RM$ -invariant*

$$(2.17) \quad \mathcal{I}_{-2/n, \beta/n}^m(g) = e^{(\beta/n)f} I(e^{-\frac{2}{n}f} g)$$

*is a conformal density of weight  $\beta$ .*

*Proof.* Let  $\hat{g} = e^{2w}g$ ; by (2.3)

$$\begin{aligned} \mathcal{I}_{-2/n, \beta/n}^m(\hat{g}) &= e^{(\beta/n)\hat{f}} I(e^{-\frac{2}{n}\hat{f}} \hat{g}) \\ &= e^{(\beta/n)(f+nw)} I(e^{-\frac{2}{n}(f+nw)} e^{2w} g) \\ &= e^{\beta w} e^{(\beta/n)f} I(e^{-\frac{2}{n}f} g) \\ &= e^{\beta w} \mathcal{I}_{-2/n, \beta/n}^m(g). \end{aligned}$$

It follows that (2.17) defines a conformal density of weight  $\beta$ . □

It is clear from the definitions that

$$(2.18) \quad \mathcal{I}_{\alpha, \beta}^{dVol(g)}(g) = I(g).$$

In particular, for the canonical base metric we have

$$(2.19) \quad \mathcal{I}_{\alpha, \beta}^m(g_m) = I(g).$$

### 3. CONFORMALLY INVARIANT CURVATURES

If we take  $I$  to be the curvature of a Riemannian metric (scalar, Ricci, etc.), then by Proposition 2.1 we can construct conformally invariant 'curvatures' associated to an  $RM$ -space:

**Definition 3.1.** *Let  $(M^n, g, m)$  be an  $RM$ -space. The conformally invariant scalar curvature of  $(M^n, g, m)$  is given by*

$$(3.1) \quad \begin{aligned} R_n^m(g) &= e^{-\frac{2}{n}f} R(g_m) \\ &= R(g) + \frac{2(n-1)}{n} \Delta_g f - \frac{(n-1)(n-2)}{n^2} |\nabla f|^2, \end{aligned}$$

*where  $f$  is given by (2.16) and  $g_m$  is the canonical base metric.  $R_n^m$  is a conformal density of weight  $-2$ : If  $\hat{g} = e^{2w}g$ , then*

$$(3.2) \quad R_n^m(\hat{g}) = e^{-2w} R_n^m(g).$$

*In particular,  $R_n^m$  shares the same scaling properties as the scalar curvature.*

Notice when  $n = 2$ , the conformally invariant scalar curvature is given by

$$(3.3) \quad R_2^m(g) = R(g) + \Delta f.$$

It will be convenient to define a scale-invariant version of  $R_n^m$ .

**Definition 3.2.** *Let  $(M^n, g, m)$  be an RM-space. The  $\tau^m$ -curvature of  $(M^n, g, m)$  is the quantity*

$$(3.4) \quad \begin{aligned} \tau^m(g) &= e^{\frac{2}{n}f} R_n^m(g) \\ &= R(g_m). \end{aligned}$$

$\tau^m(g)$  is a conformal density of weight 0: If  $\hat{g} = e^{2w}g$ , then

$$(3.5) \quad \tau^m(\hat{g}) = \tau^m(g).$$

**Definition 3.3.** *Let  $(M^n, g, m)$  be an RM-space. The conformally invariant Ricci curvature of  $(M^n, g, m)$  is given by*

$$(3.6) \quad \begin{aligned} Ric_n^m(g) &= Ric(g_m) \\ &= Ric(g) + \left(\frac{n-2}{n}\right)\nabla_g^2 f + \frac{1}{n}(\Delta_g f)g + \left(\frac{n-2}{n^2}\right)df \otimes df - \frac{(n-2)}{n^2}|\nabla f|^2 g. \end{aligned}$$

$Ric_n^m$  is a conformal density of weight 0: If  $\hat{g} = e^{2w}g$ , then

$$(3.7) \quad Ric_n^m(\hat{g}) = Ric_n^m(g).$$

In particular,  $Ric_n^m$  shares the same scaling properties as the Ricci curvature.

Note that  $R_n^m(g)$  is the contraction of  $Ric_n^m(g)$ . Moreover, by (2.18)

$$(3.8) \quad \begin{aligned} Ric_n^{dVol(g)}(g) &= Ric(g), \\ R_n^{dVol(g)}(g) &= R(g), \\ \tau^{dVol(g)}(g) &= R(g). \end{aligned}$$

for any metric  $g$ . In particular, by (2.19), for the canonical base metric  $g_m$  we have

$$(3.9) \quad \begin{aligned} Ric_n^m(g_m) &= Ric(g_m), \\ R_n^m(g_m) &= R(g_m), \\ \tau^m(g_m) &= R(g_m). \end{aligned}$$

From these identities we see that certain properties of the usual Ricci and scalar curvatures are retained by their conformally invariant counterparts:

**Proposition 3.1.** *Let  $(M^n, g, m)$  be an RM-space.*

(i) *If  $Ric_n^m(g) > 0$  (respectively,  $\geq 0$ ,  $= 0$ ,  $\leq 0$ ,  $< 0$ ), then  $[g]$  contains a metric whose Ricci curvature is positive (resp., non-negative, zero, non-positive, negative).*

(ii) *If  $R_n^m(g) > 0$  (resp.,  $\geq 0$ ,  $= 0$ ,  $\leq 0$ ,  $< 0$ ), then  $[g]$  contains a metric whose scalar curvature is positive (resp., non-negative, zero, non-positive, negative).*

*Proof.* Both statements are immediate consequences of (3.9).  $\square$

The converse of Proposition 3.1 is false. For example, suppose we are given a metric  $g$  of positive Ricci curvature; then it is easy to construct a measure  $m$  so that  $Ric_n^m(g)$  has negative eigenvalues on an open set.

**3.1. The Bianchi identity.** The contracted second Bianchi identity implies that

$$(3.10) \quad (\delta Ric)_j = \nabla^i R_{ij} = \frac{1}{2} \nabla_j R,$$

where  $\delta$  is the divergence operator. If we introduce the *Bianchi operator*  $\mathcal{B} = \mathcal{B}_g$  on symmetric  $(0, 2)$ -tensors by

$$(3.11) \quad \mathcal{B}h = \delta h - \frac{1}{2} d(\text{tr}_g h),$$

then (3.10) can be written

$$(3.12) \quad \mathcal{B}Ric = 0.$$

Similar identities hold for the conformally invariant Ricci and scalar curvatures. First, a simple calculation gives

$$(3.13) \quad \nabla^i Ric_n^m(g)_{ij} = \frac{1}{2} \nabla_j R_n^m(g) + \left(\frac{n-2}{n}\right) Ric_n^m(g)_{ij} \nabla_i f.$$

To make this more compact, we define the operator

$$(3.14) \quad \delta_n^m = \delta - \left(\frac{n-2}{n}\right) \nabla f \lrcorner$$

where  $\lrcorner$  denotes the interior product. Note that  $\delta_n^m$  is the adjoint of  $\nabla$  relative to a weight. Let

$$(3.15) \quad d\nu_n = e^{-(\frac{n-2}{n})f} dVol(g),$$

and for functions  $u, v \in L^2(M^n)$  define

$$(3.16) \quad \langle u, v \rangle_{L^2(g, d\nu_n)} = \int uv \, d\nu_n.$$

If  $X, Y$  are smooth vector fields, we define

$$\langle X, Y \rangle_{L^2(g, d\nu_n)} = \int \langle X, Y \rangle_g \, d\nu_n,$$

and use this definition to extend the  $L^2$ -inner product to tensors of all types. In particular, for a symmetric  $(0, 2)$ -tensor  $h$  and 1-form  $\eta$  we have

$$(3.17) \quad \int \langle h, \nabla \eta \rangle d\nu_n = \int -\langle \delta_n^m h, \eta \rangle d\nu_n.$$

$\delta_n^m$  can be extended to act on sections of other bundles; for example, if  $\omega \in \Omega^p(TM^n)$  and  $\theta \in \Omega^{p-1}(TM^n)$  are forms, then

$$(3.18) \quad \int \langle \omega, d\theta \rangle d\nu_n = - \int \langle \delta_n^m \omega, \theta \rangle d\nu_n$$

defines  $\delta_n^m$  on the exterior algebra.

Using  $\delta_n^m$ , (3.13) becomes

$$(3.19) \quad \delta_n^m Ric_n^m(g) = \frac{1}{2} dR_n^m(g).$$

We will refer to the operator  $\delta_n^m$  as the  $m$ -adjoint. We can also define the Bianchi operator associated to a measure:

$$(3.20) \quad \mathcal{B}_n^m h = \delta_n^m h - \frac{1}{2} d(tr_g h).$$

Then (3.19) is equivalent to

$$(3.21) \quad \mathcal{B}_n^m Ric_n^m(g) = 0.$$

It is interesting to note that  $\mathcal{B}^m$  is conformally invariant:

$$(3.22) \quad \hat{g} = e^{2w} g \implies \mathcal{B}_{n,\hat{g}}^m = e^{-2w} \mathcal{B}_{n,g}^m.$$

Acting on forms,  $\delta_n^m : \Omega^p(TM^n) \rightarrow \Omega^{p-1}(TM^n)$  is conformally "covariant":

$$(3.23) \quad \hat{g} = e^{2w} g \implies \delta_{n,\hat{g}}^m(e^{2(1-p)w}\eta) = e^{2(p-2)w} \delta_{n,g}^m \eta.$$

In particular,  $\delta_n^m$  acting on 1-forms is invariant:  $\delta_{n,\hat{g}}^m = e^{-2w} \delta_{n,g}^m$ .

**3.2. The Einstein condition.** There is a natural notion of an Einstein metric in the  $RM$ -setting.

**Definition 3.4.** *Let  $(M^n, g, m)$  be an  $RM$ -space. The conformally invariant trace-free Ricci tensor of  $(M^n, g, m)$  is given by*

$$(3.24) \quad \begin{aligned} E_n^m(g) &= Ric_n^m(g) - \frac{1}{n} R_n^m(g)g \\ &= Ric(g) - \frac{1}{n} R(g)g + \left(\frac{n-2}{n}\right) \nabla_g^2 f - \left(\frac{n-2}{n^2}\right) (\Delta_g f)g \\ &\quad + \left(\frac{n-2}{n^2}\right) df \otimes df - \left(\frac{n-2}{n^3}\right) |\nabla f|^2 g, \end{aligned}$$

Obviously  $E_n^m$  enjoys the same invariance properties as  $Ric_n^m$ . In the same way that the sign of the conformally invariant Ricci curvature detects the existence of a conformal metric with the same sign, the tensor  $E_n^m$  detects the existence of a conformal Einstein metric:

**Proposition 3.2.** *Let  $(M^n, g)$  be a smooth Riemannian manifold of dimension  $n \geq 3$ .*

(i) *If  $m$  is a measure for which  $E_n^m(g) = 0$ , then  $[g]$  contains an Einstein metric, namely, the canonical base metric.*

(ii) *If  $g$  is an Einstein metric and  $m$  is a measure for which  $E_n^m(g) = 0$ , then either  $dm = \lambda dVol(g)$  for some constant  $\lambda > 0$ , or  $(M^n, g)$  is homothetically isometric to  $(S^n, g_c)$ , where  $g_c$  denotes the round metric, and  $dm = \lambda \varphi^* dVol(g_c)$  for some conformal map  $\varphi$  and constant  $\lambda > 0$ .*



*Proof.* The proof of (i) follows from the conformal invariance of  $E_n^m$  and the definition of the canonical base metric.

To prove (ii), suppose  $g$  is Einstein and  $E_n^m(g) = 0$ . By part (i), the canonical base metric  $g_m = e^{-\frac{2}{n}f}g$  is also Einstein. This leads to two possibilities: either  $g_m$  is homothetic to  $g$ , in which case  $f = \text{const.}$  and  $dm = \lambda dVol(g)$  as claimed, or the conformal class of  $g$  admits two distinct Einstein metrics. In this latter case it follows from a theorem of Obata ([Oba71], Proposition 6.2) that  $(M^n, g)$  is homothetically isometric to  $(S^n, g_c)$ , and (also up to homothety)  $g_m$  is obtained from  $g$  by a conformal transformation. In particular,

$$\begin{aligned} m &= dVol(g_m) = \lambda dVol(\varphi^*g_c) \\ &= \lambda \varphi^* dVol(g_c). \end{aligned}$$

□

One consequence of the classical second Bianchi identity is that the scalar curvature of an Einstein metric must be constant. An analogous result holds in the setting of  $RM$ -spaces. By the Bianchi identity (3.19),

$$(3.25) \quad \delta_n^m E_n^m(g) = \frac{(n-2)}{2n} \nabla R_n^m(g) + \frac{(n-2)}{n^2} R_n^m(g) \nabla f.$$

Also,

$$\begin{aligned} dR_n^m(g) &= d(e^{-\frac{2}{n}f} \tau^m(g)) \\ &= -\frac{2}{n} R_n^m(g) df + e^{-\frac{2}{n}f} d\tau^m(g). \end{aligned}$$

Substituting this into (3.25) gives

$$(3.26) \quad \delta_n^m E_n^m(g) = \frac{(n-2)}{2n} e^{-\frac{2}{n}f} d\tau^m(g).$$

**Corollary 3.1.** *If  $(M^n, g, m)$  is an  $RM$ -space with  $n \geq 3$  and  $E_n^m(g) = 0$ , then  $\tau^m(g)$  is constant.*

*Remark 3.1.* By (3.5) the condition  $\tau^m(g) = \text{const.}$  is conformally invariant, while by (3.2) the condition  $R_n^m(g) = \text{const.}$  is not.

#### 4. RELATION TO THE CURVATURES OF BAKRY-ÉMERY AND PERELMAN

In this section we explore the relationship of the conformally invariant Ricci and scalar curvatures to the curvatures defined by Bakry-Émery [BE85] and Perelman [Per02]. To recall the definition of these latter invariants, once again we let  $(M^n, g, m)$  denote an  $RM$ -space, and let  $f$  be defined by (2.1):

$$dm = e^{-f} dVol(g).$$

The Bakry-Émery Ricci curvature is defined by

$$(4.1) \quad Ric^m(g) = Ric(g) + \nabla_g^2 f.$$

This definition arose in the analysis of infinite-dimensional diffusion processes, though this interpretation will not be relevant in the present paper. Instead, we will emphasize the role it plays as the curvature term in a "weighted" Wietzenböck formula, as we now explain.

In [Per02] (see also Lott, [Lot03]), Perelman defined the adjoint  $d^{*m}$  of the exterior derivative  $d$  relative to the measure  $m^1$ . More precisely, given a  $p$ -form  $\eta$  and smooth function  $\varphi$ ,  $d^{*m}$  is defined by

$$(4.2) \quad \int \langle d\varphi, \eta \rangle dm = - \int \varphi d^{*m} \eta dm.$$

By (2.1),

$$(4.3) \quad d^{*m} = d^* - \nabla f \lrcorner.$$

In case  $p = 1$ , in local coordinates we have

$$(4.4) \quad d^{*m} \eta = g^{kl} \nabla_k \eta_l - g^{kl} \nabla_k f \eta_l.$$

Similarly, one can define the adjoint of the covariant derivative, denoted  $\nabla^{*m}$ . Using these operators we can form the associated 'Hodge laplacian' and 'rough laplacian':

$$(4.5) \quad \square^m = dd^{*m} + d^{*m}d,$$

$$(4.6) \quad \Delta^m = \nabla^{*m} \nabla.$$

The following Weitzenböck formula for one-forms follows from the definitions of  $\square^m$  and  $\Delta^m$  (see [Per02], §1.3; [Lot03], §2):

$$(4.7) \quad \Delta^m = \square^m + Ric^m.$$

Perelman introduced a notion of the scalar curvature associated to the  $BE$ -Ricci tensor:

$$(4.8) \quad R^m = R + 2\Delta f - |\nabla f|^2.$$

In particular,  $R^m$  is *not* the trace of  $Ric^m$ . Perelman justified this definition in two ways: first, by noting that the classical contracted second Bianchi identity holds for  $Ric^m$  and  $R^m$ :

$$(4.9) \quad d^{*m} Ric^m = \frac{1}{2} dR^m.$$

Second, Perelman showed that the Dirac operator defined relative to  $dm$  in an analogous fashion leads to a Weitzenböck formula for the spin laplacian, where the usual scalar curvature term replaced by  $R^m$ .

As we now explain, the  $BE$ -Ricci curvature, and the scalar curvature defined by Perelman, can both be viewed as the infinite-dimensional limit of their conformally invariant counterparts. To see this we make use of a common construction, that of taking the product of  $M^n$  with a manifold of large dimension (see, e.g., [Lot03]).

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<sup>1</sup>So that our laplacians agree with the Euclidean laplace operator, we add the minus sign.

First, however, let us argue informally: If we let the dimension  $n \rightarrow \infty$  in (3.6), then in the limit we obtain the *BE*-Ricci curvature:

$$\begin{aligned} \lim_{n \rightarrow \infty} Ric_n^m(g) & \text{''} = \text{''} Ric + \nabla_g^2 f \\ & = Ric^m(g). \end{aligned}$$

Similarly, letting  $n \rightarrow \infty$  in the definition of  $R_n^m$  gives the scalar curvature defined by Perelman:

$$\begin{aligned} \lim_{n \rightarrow \infty} R_n^m(g) & \text{''} = \text{''} R + 2\Delta_g f - |\nabla f|^2 \\ & = R^m(g). \end{aligned}$$

In these formulas we used quotation marks to emphasize the fact that the process of letting the dimension go to infinity is not meant to be taken literally. However, we can make these formal observations more concrete by the aforementioned construction.

Let  $(\mathbb{T}^{d-n}, ds^2)$  denote the flat  $(d-n)$ -dimensional torus, and consider the product manifold  $N^d = M^n \times \mathbb{T}^{d-n}$  with the product metric  $h = g + ds^2$ . To define the conformally invariant Ricci and scalar curvatures on  $N^d$  we need to define a measure, so we take the product measure  $d\mu = dm \times dVol(ds^2)$ . Note that if  $dm = e^{-f} dVol(g)$  and  $d\mu = e^{-\tilde{f}} dVol(h)$ , it follows that  $\tilde{f} = f$ . Therefore, given tangent vectors  $X, Y \in TM^n$ , if we denote their lift to  $TN^d$  also by  $X, Y$ , then by the definition of  $Ric_n^m$

$$\begin{aligned} \lim_{d \rightarrow \infty} Ric_d^m(h)(X, Y) &= \lim_{d \rightarrow \infty} \left\{ Ric(h)(X, Y) + \left(\frac{d-2}{d}\right) \nabla_h^2 \tilde{f}(X, Y) + \frac{1}{d} (\Delta_h \tilde{f}) h(X, Y) \right. \\ &\quad \left. + \left(\frac{d-2}{d^2}\right) (d\tilde{f} \otimes d\tilde{f})(X, Y) - \frac{(d-2)}{d^2} |\nabla \tilde{f}|^2 h(X, Y) \right\} \\ &= \lim_{d \rightarrow \infty} \left\{ Ric(g)(X, Y) + \left(\frac{d-2}{d}\right) \nabla_g^2 f(X, Y) + \frac{1}{d} (\Delta_g f) g(X, Y) \right. \\ &\quad \left. + \left(\frac{d-2}{d^2}\right) df(X) df(Y) - \frac{(d-2)}{d^2} |\nabla f|^2 g(X, Y) \right\} \\ &= Ric(g)(X, Y) + \nabla^2 f(X, Y) \\ &= Ric_g^m(X, Y). \end{aligned}$$

Similarly,

$$\lim_{d \rightarrow \infty} R_d^m(h) = R^m(g).$$

A similar argument can be used to relate the  $m$ -divergence defined in Section 3.1 with the operator  $d^{*m}$ . Comparing (3.14) and (4.3) one sees that formally

$$\text{''} \lim_{n \rightarrow \infty} \delta_n^m = d^{*m}. \text{''}$$

This can also be made precise using the construction above.

5.  $\mathcal{D}$  AND  $\mathcal{C}_+$ -ACTIONS ON  $RM$ -SPACES

As we observed in Section 2, there are two natural group actions on  $RM$ -spaces: the natural  $\mathcal{D}$ -action defined by (2.6),

$$\varphi : (M^n, g, m) \mapsto (M^n, \varphi^*g, \varphi^*m),$$

and the action of  $\mathcal{C}_+$  via conformal rescaling:

$$(5.1) \quad \gamma : (M^n, g, m) \mapsto (M^n, \gamma g, m).$$

Note that  $\mathcal{C}_+$  acts on the metric alone, while  $\mathcal{D}$  acts by pull-back on *both* the metric and the measure. On the other hand, if  $T = T(g, m)$  is an  $RM$ -invariant and we pull back the metric  $g$  by a diffeomorphism  $\varphi \in \mathcal{D}$  (leaving  $m$  fixed), then by (2.6),

$$(5.2) \quad T(\varphi^*g, m) = \varphi^*T(g, (\varphi^{-1})^*m).$$

One way of interpreting (5.2) is to view pull-backs of the metric as determining a change in the measure,  $m \mapsto (\varphi^{-1})^*m$ . By a Theorem of Moser, this is actually an equivalence; that is, *all* measures can be realized in this manner:

**Theorem 5.1.** *Let  $(M^n, g, m)$  be an  $RM$ -space and  $T = T(g, m)$  a local  $RM$ -invariant. Given a measure  $\mu$  with*

$$(5.3) \quad \int d\mu = \int dm,$$

*there is a diffeomorphism  $\varphi \in \mathcal{D}$  with*

$$(5.4) \quad T(\varphi^*g, m) = \varphi^*T(g, \mu).$$

*Conversely, given a diffeomorphism  $\varphi \in \mathcal{D}$ , there is a measure  $\mu$  satisfying (5.3) and (5.4).*

*Proof.* By Moser's theorem [Mos65], equation (5.3) implies the existence of a diffeomorphism  $\varphi \in \mathcal{D}$  such that  $\mu = (\varphi^{-1})^*m$ . Therefore, (5.4) follows from (5.2).

The converse is obvious; just take  $\mu = (\varphi^{-1})^*m$ .  $\square$

Our next result shows that the behavior of a Riemannian invariant  $I = I(g)$  under conformal deformations of the metric is *equivalent* to the behavior of the associated conformal density  $\mathcal{I}_{-2/n,0}^m(g)$  under conjugation by a diffeomorphism. To simplify the statement we normalize the measure and the metric:

**Theorem 5.2.** *Let  $(M^n, g, m)$  be an  $RM$ -space, where for convenience we assume  $g$  and  $m$  have the same total mass:*

$$(5.5) \quad \int dVol(g) = \int dm.$$

*Suppose  $I = I(g)$  is a Riemannian invariant, and  $\mathcal{I}^m = \mathcal{I}_{-2/n,0}^m(g)$  is the conformal density of weight 0 defined in (2.17). Given  $\hat{g} \in [g]$  of the same volume, there is a diffeomorphism  $\varphi \in \mathcal{D}$  such that*

$$(5.6) \quad I(\hat{g}) = (\varphi^{-1})^*\mathcal{I}^m(\varphi^*g).$$

Conversely, given a diffeomorphism  $\varphi \in \mathcal{D}$ , there is a conformal metric  $\hat{g} \in [g]$  (of the same volume) such that (5.6) holds.

*Proof.* Given  $\hat{g} \in [g]$  with the same volume of  $g$  (and consequently, the same mass as  $m$ ), by Moser's Theorem there is a diffeomorphism  $\varphi \in \mathcal{D}$  such that

$$(5.7) \quad (\varphi^{-1})^* dm = dVol(\hat{g}).$$

Therefore, by (5.2),

$$\begin{aligned} \mathcal{I}^m(\varphi^*g) &= \varphi^*(\mathcal{I}^{(\varphi^{-1})^*m}(g)) \\ &= \varphi^*(\mathcal{I}^{dVol(\hat{g})}(g)) \\ &= \varphi^*(\mathcal{I}^{dVol(\hat{g})}(\hat{g})) \quad (\text{by conformal invariance of } \mathcal{I}^m) \\ &= \varphi^*(I(\hat{g})). \quad (\text{by (2.18)}) \end{aligned}$$

Since  $I$  is a Riemannian invariant, pulling back both sides by  $\varphi^{-1}$  gives (5.6).

Conversely, suppose  $\varphi \in \mathcal{D}$  is given. Then

$$(5.8) \quad (\varphi^{-1})^* m = e^{nw} dVol(g)$$

for some function  $w$ . Let  $\hat{g} = e^{2w}g$ ; then (5.8) can be rewritten

$$(5.9) \quad (\varphi^{-1})^* m = dVol(\hat{g}).$$

Therefore,

$$\begin{aligned} \mathcal{I}^m(\varphi^*g) &= \varphi^*(\mathcal{I}^{(\varphi^{-1})^*m}(g)) \\ &= \varphi^*(\mathcal{I}^{dVol(\hat{g})}(g)) \\ &= \varphi^*(\mathcal{I}^{dVol(\hat{g})}(\hat{g})) \quad (\text{by conformal invariance of } \mathcal{I}^m) \\ &= \varphi^*(I(\hat{g})), \quad (\text{by (2.18)}) \end{aligned}$$

and pulling back by  $\varphi$  we get (5.6). □

*Remark 5.3.* Theorem 5.2 can be regarded as the fundamental theorem of the geometry of  $RM$ -spaces, in that it gives the relationship between  $RM$ -conformal densities and Riemannian invariants.

**Corollary 5.1.** *Let  $(M^n, g, m)$  be an  $RM$ -space, where for convenience we assume  $g$  and  $m$  have the same total mass. Given  $\hat{g} \in [g]$  with the same volume as  $g$ , there is a diffeomorphism  $\varphi \in \mathcal{D}$  such that*

$$(5.10) \quad R(\hat{g}) = (\varphi^{-1})^* \tau^m(\varphi^*g).$$

Conversely, given a diffeomorphism  $\varphi \in \mathcal{D}$ , there is a conformal metric  $\hat{g} \in [g]$  such that (5.10) holds.

## 6. VARIATIONAL THEORY

In this section we consider variational problems associated to  $RM$ -spaces. We begin with some general features of variational integrals, and describe a natural constrained problem. In the succeeding section we explore in detail an interesting example associated to the scalar curvature.

Let  $\mathcal{M} = \mathcal{M}(M^n)$  denote the space of Riemannian metrics on  $M^n$ , and  $\mathcal{P}$  the space of all probability measures. Sometimes we may need to normalize the volume, so we denote by  $\mathcal{M}_1$  the subspace of unit-volume metrics.

**Definition 6.1.** *An  $RM$ -functional is a (differentiable) mapping  $\mathcal{F} : \mathcal{M} \times \mathcal{P} \rightarrow \mathbb{R}$  which is invariant under the action of  $\mathcal{D}$  in (2.7):*

$$(6.1) \quad \mathcal{F}(\varphi^*g, \varphi^*m) = \mathcal{F}(g, m).$$

The basic example of an  $RM$ -functional is the integral of a scalar local  $RM$ -invariant  $q = q(g, m)$ :

$$(6.2) \quad \mathcal{F}(g, m) = \int q(g, m) dVol(g).$$

Of particular interest will be the case where  $q$  is the conformal density  $\mathcal{I}^m(g) = \mathcal{I}_{-2/n,0}^m(g)$  associated to a Riemannian invariant  $I(g)$ , as defined in Proposition 2.1. Let

$$(6.3) \quad \mathcal{J}^m[g] = \int \mathcal{I}^m(g) dm,$$

and denote the corresponding Riemannian functional by

$$(6.4) \quad J[g] = \int I(g) dVol(g).$$

**Example.** Let  $I(g) = R(g)$ , the scalar curvature. Then

$$(6.5) \quad \mathcal{S}^m[g] = \int \tau^m(g) dm,$$

while the Riemannian functional is the total scalar curvature:

$$(6.6) \quad S[g] = \int R(g) dVol(g).$$

Let  $\nabla J$  denote the  $L^2$ -gradient of the functional  $J[\cdot]$  (assuming it exists):

$$(6.7) \quad \frac{d}{ds} J[g + sh] \Big|_{s=0} = \int \langle \nabla J(g), h \rangle_g dVol(g).$$

**Proposition 6.1.** *Suppose  $(M^n, g, m)$  is an  $RM$ -space, and  $I = I(g)$  is a scalar Riemannian invariant. Let  $\mathcal{I}^m = \mathcal{I}_{-2/n,0}^m$  denote the associated conformal density, and define the functionals  $J$  and  $\mathcal{J}^m$  by (6.4) and (6.3), respectively. Then*

(i) *The functional  $\mathcal{J}^m$  is conformally invariant:*

$$(6.8) \quad \mathcal{J}^m[e^{2w}g] = \mathcal{J}^m[g].$$

(ii) If the functional  $J$  is differentiable with  $L^2$ -gradient  $\nabla J$ , then  $\mathcal{J}^m$  is also differentiable, and

$$(6.9) \quad \nabla \mathcal{J}^m(g) = e^{\frac{2}{n}f} \left\{ \nabla J(g_m) - \frac{1}{n} \text{tr}_{g_m}(\nabla J(g_m))g_m \right\}.$$

*Proof.* The conformal invariance of  $\mathcal{J}^m$  is obvious from the definition. To prove (6.9), we first note that the gradient of  $J$  can be expressed in terms of the  $L^2$ -adjoint of the linearized operator:

$$\begin{aligned} \frac{d}{ds} J[g + sh] \Big|_{s=0} &= \int \frac{d}{ds} I(g + sh) \Big|_{s=0} dVol(g) + \int I(g) \frac{d}{ds} dVol(g + sh) \Big|_{s=0} \\ &= \int I'(g)[h] dVol(g) + \int \frac{1}{2} I(g)(\text{tr}_g h) dVol(g) \\ &= \left\langle I'(g)[h], 1 \right\rangle_{L^2} + \left\langle h, \frac{1}{2} I(g)g \right\rangle_{L^2} \\ &= \int \left\langle h, I'(g)^*(1) + \frac{1}{2} I(g)g \right\rangle_g dVol(g), \end{aligned}$$

which implies

$$(6.10) \quad \nabla J(g) = I'(g)^*(1) + \frac{1}{2} I(g)g.$$

Similarly, for  $\mathcal{J}^m$  we find

$$\begin{aligned} \frac{d}{ds} \mathcal{J}^m[g + sh] \Big|_{s=0} &= \int \frac{d}{ds} \mathcal{I}^m(g + sh) \Big|_{s=0} dm \\ &= \int (\mathcal{I}^m)'(g)[h] dm \\ &= \left\langle (\mathcal{I}^m)'(g)[h], 1 \right\rangle_{L^2(g, dm)} \\ &= \left\langle h, (\mathcal{I}^m)'(g)^*(1) \right\rangle_{L^2(g, dm)}, \end{aligned}$$

hence

$$(6.11) \quad \nabla \mathcal{J}^m(g) = (\mathcal{I}^m)'(g)^*(1).$$

**Claim 6.1.**

$$(6.12) \quad (\mathcal{I}^m)'(g)^*u = e^{\frac{2}{n}f} \left\{ I'(g_m)^*u - \frac{1}{n} \text{tr}_{g_m} \left( I'(g_m)^*u \right) g_m \right\},$$

where  $g_m = e^{-\frac{2}{n}f}g$  is the canonical base metric.

*Proof.* The claim follows from the identity

$$(6.13) \quad (\mathcal{I}^m)'(g)[h] = I'(g_m)[e^{-\frac{2}{n}f}\mathring{h}],$$

where

$$\mathring{h} = h - \frac{1}{n}(\text{tr}_g h)g,$$

by simply integrating by parts. To verify (6.13), let  $f_s$  be defined by

$$(6.14) \quad dm = e^{-f_s} dVol(g + sh).$$

Then

$$(6.15) \quad \frac{d}{ds} f_s \Big|_{s=0} = \frac{1}{2} tr_g h.$$

Since  $\mathcal{I}^m(g) = I(e^{-\frac{2}{n}f}g)$ , we find

$$\begin{aligned} (\mathcal{I}^m)'(g)[h] &= \frac{d}{ds} \mathcal{I}^m(g + sh) \Big|_{s=0} \\ &= \frac{d}{ds} I(e^{-\frac{2}{n}f_s}(g + sh)) \Big|_{s=0} \\ &= I'(e^{-\frac{2}{n}f}g) \left[ e^{-\frac{2}{n}f} \left( h - \frac{1}{n} (tr_g h) g \right) \right]. \end{aligned}$$

□

Using (6.12) and (6.10) we conclude

$$\begin{aligned} \nabla \mathcal{J}^m(g) &= (\mathcal{I}^m)'(g)^*(1) \\ &= e^{\frac{2}{n}f} \left\{ I'(g_m)^*(1) - \frac{1}{n} tr_{g_m} \left( I'(g_m)^*(1) \right) g_m \right\} \\ &= e^{\frac{2}{n}f} \left\{ \nabla J(g_m) - \frac{1}{n} tr_{g_m} (\nabla J(g_m)) g_m \right\}. \end{aligned}$$

□

**Example.** The gradient of the total scalar curvature is the 'gravitational tensor':

$$\nabla S(g) = -Ric(g) + \frac{1}{2} R(g)g,$$

whose trace is

$$tr_g \nabla S(g) = \frac{(n-2)}{2} R(g).$$

Therefore, by Proposition 6.1 the gradient of  $\mathcal{S}^m$  is

$$\begin{aligned} \nabla \mathcal{S}^m(g) &= e^{\frac{2}{n}f} \left\{ -Ric(g_m) + \frac{1}{2} R(g_m)g_m - \frac{1}{n} \left[ \frac{(n-2)}{2} R(g_m) \right] g_m \right\} \\ (6.16) \quad &= e^{\frac{2}{n}f} \left\{ -Ric(g_m) + \frac{1}{n} R(g_m)g_m \right\} \\ &= -e^{\frac{2}{n}f} E_n^m(g), \end{aligned}$$

where  $E_n^m(g)$  is the conformally invariant trace-free Ricci tensor (see Section 3.2).

**Proposition 6.2.** *Let  $(M^n, g_m)$  be an RM-space of dimension  $n \geq 3$ . Then  $g$  is critical for  $\mathcal{S}^m$  if and only if it is conformal to an Einstein metric.*



**6.1. A constrained variational problem.** Since we are fixing the measure when studying  $\mathcal{J}^m$ , a natural constrained variational problem is to restrict the functional to the orbit of a metric  $g$  under the action of  $\mathcal{D}$ . By Theorem 5.2, this is equivalent to restricting the associated functional  $J$  to the conformal class of  $g$ :

**Theorem 6.1.** *Let  $(M^n, g, m)$  be an RM-space, where  $m$  is a probability measure and we normalize  $g$  to have unit volume. Let  $I, J, \mathcal{I}^m$ , and  $\mathcal{J}^m$  be as above. Then*

$$(6.17) \quad \mathcal{J}^m|_{\mathcal{D}(g)} = J|_{[g]_1},$$

where  $\mathcal{D}(g) = \{\varphi^*g \mid \varphi \in \mathcal{D}\}$  is the orbit of  $g$  under the action of  $\mathcal{D}$ , and  $[g]_1$  denotes metrics conformal to  $g$  of unit volume.

More precisely: Given  $\varphi \in \mathcal{D}$ , there is a conformal metric  $\hat{g} \in [g]_1$  with

$$(6.18) \quad \mathcal{J}^m[\varphi^*g] = J[\hat{g}].$$

Conversely, given  $\hat{g} \in [g]_1$ , there is a diffeomorphism  $\varphi \in \mathcal{D}$  such that (6.18) holds.

*Proof.* This equivalence is essentially a corollary of Theorem 5.2 and its proof. Given  $\varphi \in \mathcal{D}$ , there is a conformal metric  $\hat{g} \in [g]_1$  with

$$(\varphi^{-1})^*\mathcal{I}^m(\varphi^*g) = I(\hat{g}).$$

Pulling back by  $\varphi$  gives

$$(6.19) \quad \mathcal{I}^m(\varphi^*g) = \varphi^*I(\hat{g}).$$

Also, by (5.9),

$$(6.20) \quad dm = \varphi^*dVol(\hat{g}).$$

Therefore,

$$(6.21) \quad \begin{aligned} \mathcal{J}^m[\varphi^*g] &= \int \mathcal{I}^m(\varphi^*g) dm \\ &= \int \varphi^*I(\hat{g}) \varphi^*dVol(\hat{g}) \\ &= \int I(\hat{g}) dVol(\hat{g}) \\ &= J[\hat{g}]. \end{aligned}$$

Conversely, by Theorem 5.2, if  $\hat{g}$  is a conformal metric of unit volume then there is a diffeomorphism  $\varphi \in \mathcal{D}$  satisfying (6.19) and (6.20). Then (6.21) implies

$$J[\hat{g}] = \mathcal{J}^m[\varphi^*g],$$

and (6.18) follows. □

7. THE TOTAL  $\tau^m$ -CURVATURE

Now let  $I = I(g)$  be the scalar curvature, and consider the variational integral associated to its  $RM$ -counterpart; i.e., the  $\tau^m$ -curvature. As a corollary of Theorem 6.1, we can give a different characterization of the Yamabe invariant:

**Theorem 7.1.** *Let  $(M^n, g, m)$  be an  $RM$ -space of dimension  $n \geq 3$ , where  $m$  is a probability measure and we normalize  $g$  to have unit volume. Then*

$$(7.1) \quad \inf_{\varphi \in \mathcal{D}} \int \tau^m(\varphi^*g) \, dm = Y(M^n, [g]),$$

where  $Y(M^n, [g])$  is the Yamabe invariant of  $g$ . Moreover, the infimum on the left-hand side of (7.1) is attained by a diffeomorphism if and only if there is a conformal metric which attains the Yamabe invariant.

*Remark 7.2.* Recall the  $\sigma$ -constant of a manifold  $M^n$  is defined to be

$$(7.2) \quad \sigma(M^n) = \sup_{g \in \mathcal{M}} Y(M^n, [g]).$$

If we let

$$(7.3) \quad \sigma(M^n, m) \equiv \sup_{g \in \mathcal{M}_1} \inf_{\varphi \in \mathcal{D}} \mathcal{S}^m(\varphi^*g),$$

then it follows from Theorem 7.1 that

$$(7.4) \quad \sigma(M^n, m) = \sigma(M^n),$$

independent of the measure  $m$ .

More generally, a metric is a critical point of the total scalar curvature constrained to a fixed conformal class if and only if it has constant scalar curvature. The corresponding result for the total  $\tau^m$ -curvature is

**Proposition 7.1.** *When the dimension  $n \geq 3$ , a metric  $g$  is critical for  $\mathcal{S}^m|_{\mathcal{D}(g)}$  if and only if  $\tau^m(g) = \text{const}$ .*

*Proof.* Let  $\{\varphi_t\}$  be the 1-parameter family of diffeomorphisms generated by a vector field  $X$ . By (6.16),

$$(7.5) \quad \begin{aligned} 0 &= \frac{d}{dt} \mathcal{S}^m(\varphi_t^*g) \Big|_{t=0} \\ &= \mathcal{S}^m(g)'[L_X g] \\ &= - \int \langle e^{\frac{2}{n}f} E_n^m(g), L_X g \rangle \, dm \\ &= - \int \langle E_n^m(g), L_X g \rangle \, d\nu_n. \end{aligned}$$

The Proposition will follow from the next Lemma:

**Lemma 7.1.** *For any vector field  $X$ ,*

$$(7.6) \quad \int -\langle E_n^m(g), L_X g \rangle d\nu_n = \frac{(n-2)}{n} \int X\tau^m(g) dm.$$

*Proof.* Let  $\alpha = X^\flat$  be the one-form dual to the vector field  $X$ , then

$$\langle E_n^m(g), L_X g \rangle = 2\langle E_n^m(g), \nabla\alpha \rangle.$$

Therefore,

$$\begin{aligned} \int -\langle E_n^m(g), L_X g \rangle d\nu_n &= \int -2\langle E_n^m(g), \nabla\alpha \rangle d\nu_n \\ &= \int 2\langle \delta_n^m E_n^m(g), \alpha \rangle d\nu_n \quad (\text{by (3.17)}) \\ &= \int \frac{(n-2)}{n} \langle e^{-\frac{2}{n}f} \nabla\tau^m(g), \alpha \rangle d\nu_n \quad (\text{by (3.26)}) \\ &= \frac{(n-2)}{n} \int \langle \nabla\tau^m(g), \alpha \rangle dm \\ &= \frac{(n-2)}{n} \int X\tau^m(g) dm. \end{aligned}$$

□

To complete the proof of the Proposition, for any vector field  $X$  (7.6) implies

$$\begin{aligned} 0 &= \frac{d}{dt} \mathcal{S}^m(\varphi_t^* g) \Big|_{t=0} \\ &= - \int \langle E_n^m(g), L_X g \rangle dV ol(g) \\ &= \frac{(n-2)}{n} \int X\tau^m(g) dm, \end{aligned}$$

and it follows that  $\tau^m(g)$  must be constant. □

In view of Proposition 7.1, it is natural to introduce the following linear functional  $\mathcal{G}^m(g) : \mathfrak{X}(M^n) \rightarrow \mathbb{R}$ :

$$(7.7) \quad \mathcal{G}^m(g)(X) = \int X\tau^m(g) dm.$$

This is the obvious extension of the Futaki invariant from Kähler geometry [Fut83], or the Kazadan-Warner integral from conformal geometry [KW74].

**Proposition 7.2.** (i) *The functional  $\mathcal{G}^m(g)$  is conformally invariant: if  $\hat{g} = e^{2w}g$ , then*

$$\mathcal{G}^m(\hat{g}) = \mathcal{G}^m(g).$$

(ii) *If  $X$  is a conformal Killing vector field, then  $\mathcal{G}^m(g)(X) = 0$ .*

(iii) If  $E_n^m(g) = 0$ , then  $\mathcal{G}^m(g) = 0$ . In particular, unless  $\mathcal{G}^m(g) \equiv 0$  for some measure  $m$ , the conformal class of  $g$  does not contain an Einstein metric.

*Proof.* (i) The conformal invariance of  $\mathcal{G}^m(g)$  follows from the conformal invariance of  $\tau^m$ .

(ii) If  $X$  is a conformal vector field, then  $L_X g = \psi g$  for some function  $\psi$ . Therefore,  $\langle E_n^m(g), L_X g \rangle = 0$ , and from (7.6) it follows that  $\mathcal{G}^m(g)(X) = 0$ .

(iii) If  $E_n^m(g) = 0$ , then (7.6) obviously implies  $\mathcal{G}^m(g) = 0$ . □

## REFERENCES

- [BE85] E. Bakry and Michel Émery, *Diffusions hypercontractives*, Sminaire de probabilités, XIX, 1983/84, Lecture Notes in Math., vol. 1123, Springer, Berlin, 1985, pp. 177–206.
- [CGY06] Sun-Yung Alice Chang, Matthew J. Gursky, and Paul Yang, *Conformal invariants associated to a measure*, Proc. Nat. Acad. Sci. **103** (2006), 2535–2540.
- [CGY07] ———, *Conformal invariants associated to a measure, II: Conformally covariant operators*, preprint, 2007.
- [FG85] Charles Fefferman and C. Robin Graham, *Conformal invariants*, The mathematical heritage of Élie Cartan (Lyon, 1984), Astrisque, vol. Numero Hors Serie, 1985, pp. 95–116.
- [FG02] C. Fefferman and C.R. Graham, *Q-curvature and Poincaré metrics*, Math. Res. Lett. **9** (2002), 139–152.
- [Fut83] A. Futaki, *An obstruction to the existence of Einstein Kähler metrics*, Invent. Math. **73** (1983), no. 3, 437–443.
- [GH05] C. Robin Graham and Kengo Hirachi, *The ambient obstruction tensor and Q-curvature*, AdS/CFT correspondence: Einstein metrics and their conformal boundaries, Eur. Math. Soc., Zürich, IRMA Lect. Math. Theor. Phys., vol. 8, 2005, pp. 59–71.
- [GJMS92] C. Robin Graham, Ralph Jenne, Lionel J. Mason, and George A.J. Sparling, *Conformally invariant powers of the Laplacian, I. Existence.*, J. London Math. Soc. **46** (1992), 557–565.
- [KW74] Jerry L. Kazdan and F. W. Warner, *Curvature functions for compact 2-manifolds*, Ann. of Math. **99** (1974), no. 2, 14–47.
- [Lot03] John Lott, *Some geometric properties of the Bakry-Émery-Ricci tensor*, Comment. Math. Helv. **78** (2003), no. 4, 865–883.
- [Mos65] Jürgen Moser, *On the volume elements on a manifold*, Trans. Amer. Math. Soc. **120** (1965), 286–294.
- [Oba71] Morio Obata, *The conjectures on conformal transformations of Riemannian manifolds*, J. Differential Geometry **6** (1971), 247–258.
- [Per02] Grisha Perelman, *The entropy formula for the Ricci flow and its geometric applications*, preprint, 2002.

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