# Note on Lagrangian-Eulerian Methods for Uniqueness in Hydrodynamic Systems 

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#### Abstract

We discuss the Lagrangian-Eulerian framework for hydrodynamic models and provide a proof of Lipschitz dependence of solutions on initial data in path space. The paper presents a corrected version of the result in [1]. MSC Classification: 35Q30, 35Q35.


## 1 Introduction

Many hydrodynamical systems consist of evolution equations for fluid velocities forced by external stresses, coupled to evolution equations for the external stresses. In the simplest cases, the Eulerian velocity $u$ can be recovered from the stresses $\sigma$ via a linear operator

$$
\begin{equation*}
u=\mathbb{U}(\sigma) \tag{1}
\end{equation*}
$$

and the stress matrix $\sigma$ obeys a transport and stretching equation of the form

$$
\partial_{t} \sigma+u \cdot \nabla \sigma=F(\nabla u, \sigma),
$$

where $F$ is a nonlinear coupling depending on the model. The Eulerian velocity gradient is obtained in terms of the operator

$$
\begin{equation*}
\nabla_{x} u=\mathbb{G}(\sigma), \tag{2}
\end{equation*}
$$

and, in many cases, $\mathbb{G}$ is bounded in Hölder spaces of low regularity. Then, passing to Lagrangian variables,

$$
\tau=\sigma \circ X
$$

where $X$ is the particle path transformation $X(\cdot, t): \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$, a volume preserving diffeomorphism, the system becomes

$$
\left\{\begin{align*}
\partial_{t} X & =\mathcal{U}(X, \tau),  \tag{3}\\
\partial_{t} \tau & =\mathcal{T}(X, \tau) .
\end{align*}\right.
$$

with

$$
\begin{gather*}
\mathcal{U}(X, \tau)=\mathbb{U}\left(\tau \circ X^{-1}\right) \circ X, \\
\mathcal{T}(X, \tau)=F\left(\mathbb{G}\left(\tau \circ X^{-1}\right) \circ X, \tau\right) . \tag{4}
\end{gather*}
$$

In particular, $\tau$ solves an ODE

$$
\begin{equation*}
\frac{d}{d t} \tau=F(g, \tau) \tag{5}
\end{equation*}
$$

where $g=\nabla_{x} u \circ X$ is of the same order of magnitude as $\tau$ in appropriate spaces, and so the size of $\tau$ is readily estimated from the information provided by the ODE model, analysis of $\mathbb{G}$ and of the

[^0]operation of composition with $X$. The main additional observation that leads to Lipschitz dependence in path space is that derivatives with respect to parameters of expressions of the type encountered in the Lagrangian evolution (4),
$$
\mathbb{U}\left(\tau \circ X^{-1}\right) \circ X, \quad \mathbb{G}\left(\tau \circ X^{-1}\right) \circ X
$$
introduce commutators, and these are well behaved in spaces of relatively low regularity. The LagrangianEulerian method of [2] formalized these considerations leading to uniqueness and Lipschitz dependence on initial data in path space, with application to several examples including incompressible 2D and 3D Euler equations, the surface quasi-geostrophic equation (SQG), the incompressible porous medium equation, the incompressible Boussinesq system, and the Oldroyd-B system coupled with the steady Stokes system. In all these examples the operators $\mathbb{U}$ and $\mathbb{G}$ are time-independent.

The paper [1] considered time-dependent cases. When the operators $\mathbb{U}$ and $\mathbb{G}$ are time-dependent, in contrast to the time-independent cases studied in $[2], \mathbb{G}$ is not necessarily bounded in $L^{\infty}\left(0, T ; C^{\alpha}\right)$. This was addressed in [1] by using a Hölder continuity $\sigma \in C^{\beta}\left(0, T ; C^{\alpha}\right)$. While this treated the Eulerian issue, it was tacitly used but never explicitly stated in [1] that this kind of Hölder continuity is transferred to $\sigma$ from $\tau$ by composition with a smooth time-depending diffeomorphism close to the identity. This is false. In fact, we can easily give examples of $C^{\alpha}$ functions $\tau$ which are time-independent (hence analytic in time with values in $C^{\alpha}$ ) and diffeomorphisms $X(t)(a)=a+v t$ with constant $v$, such that $\sigma=\tau \circ X^{-1}$ is not continuous in $C^{\alpha}$ as a function of time. In this paper we present a correct version of the results in [1]. Instead of relying on the time regularity of $\tau$ alone, we also use the fact that $\mathbb{G}$ is composed from a time-independent bounded operator and an operator whose kernel is smooth and rapidly decaying in space. Then the time singularity is resolved by using the Lipschitz dependence in $L^{1}$ of Schwartz functions composed with smoothly varying diffeomorphisms near the identity.

A typical example of the systems we can treat is the Oldroyd-B system coupled with Navier-Stokes equations:

$$
\left\{\begin{array}{c}
\partial_{t} u-\nu \Delta u=\mathbb{H}(\operatorname{div}(\sigma-u \otimes u))  \tag{6}\\
\nabla \cdot u=0 \\
\partial_{t} \sigma+u \cdot \nabla \sigma=(\nabla u) \sigma+\sigma(\nabla u)^{T}-2 k \sigma+2 \rho K\left((\nabla u)+(\nabla u)^{T}\right) \\
u(x, 0)=u_{0}(x), \sigma(x, 0)=\sigma_{0}(x)
\end{array}\right.
$$

Here $(x, t) \in \mathbb{R}^{d} \times[0, T)$. The Leray-Hodge projector $\mathbb{H}=\mathbb{I}+R \otimes R$ is given in terms of the Riesz transforms $R=\left(R_{1}, \ldots, R_{d}\right)$, and $\nu, \rho K, k$ are fixed positive constants. This system is viscoelastic, and the behavior of the solution depends on the history of its deformation.
The non-resistive MHD system

$$
\left\{\begin{array}{c}
\partial_{t} u-\nu \Delta u=\mathbb{H}(\operatorname{div}(b \otimes b-u \otimes u)),  \tag{7}\\
\nabla \cdot u=0 \\
\nabla \cdot b=0 \\
\partial_{t} b+u \cdot \nabla b=(\nabla u) b, \\
u(x, 0)=u_{0}(x), b(x, 0)=b_{0}(x)
\end{array}\right.
$$

can also be treated by this method. The systems (6) and (7) have been studied extensively, and a review of the literature is beyond the scope of this paper.

## 2 The Lagrangian-Eulerian formulation

We show calculations for (6) in order to be explicit, and because the calculations for (7) are entirely similar. The solution map for $u(x, t)$ of (6) is

$$
\begin{equation*}
u(x, t)=\mathbb{L}_{\nu}\left(u_{0}\right)(x, t)+\int_{0}^{t} g_{\nu(t-s)} *(\mathbb{H}(\operatorname{div}(\sigma-u \otimes u)))(x, s) d s \tag{8}
\end{equation*}
$$

where

$$
\begin{equation*}
\mathbb{L}_{\nu}\left(u_{0}\right)(x, t)=g_{\nu t} * u_{0}(x)=\int_{\mathbb{R}^{d}} \frac{1}{(4 \pi \nu t)^{\frac{d}{2}}} e^{-\frac{|x-y|^{2}}{4 \nu t}} u_{0}(y) d y \tag{9}
\end{equation*}
$$

Thoroughout the paper we use

$$
g_{\nu t}(x)=\frac{1}{(4 \pi \nu t)^{\frac{d}{2}}} e^{-\frac{|x|^{2}}{4 \nu t}}
$$

The velocity gradient satisfies

$$
\begin{equation*}
(\nabla u)(x, t)=\mathbb{L}_{\nu}\left(\nabla u_{0}\right)(x, t)+\int_{0}^{t}\left(g_{\nu(t-s)} *(\mathbb{H} \nabla \operatorname{div}(\sigma-u \otimes u))\right)(x, s) d s \tag{10}
\end{equation*}
$$

We denote the Eulerian velocity and gradient operators

$$
\left\{\begin{align*}
\mathbb{U}(f)(x, t) & =\int_{0}^{t}\left(g_{\nu(t-s)} * \mathbb{H} \operatorname{div} f\right)(x, s) d s  \tag{11}\\
\mathbb{G}(f)(x, t) & =\int_{0}^{t}\left(g_{\nu(t-s)} * \mathbb{H} \nabla \operatorname{div} f\right)(x, s) d s
\end{align*}\right.
$$

Note that for a second order tensor $f, \mathbb{G}(f)=\nabla_{x} \mathbb{U}(f)=R \otimes R\left(\mathbb{U}\left(\nabla_{x} f\right)\right)$. Let $X$ be the Lagrangian path diffeomorphism, $v$ the Lagrangian velocity, and $\tau$ the Lagrangian added stress,

$$
\begin{gather*}
v=\frac{\partial X}{\partial t}=u \circ X,  \tag{12}\\
\tau=\sigma \circ X
\end{gather*}
$$

We also set

$$
\begin{gather*}
g(a, t)=(\nabla u)(X(a, t), t)=\mathbb{L}_{\nu}\left(\nabla u_{0}\right) \circ X(a, t) \\
+\mathbb{G}\left(\tau \circ X^{-1}\right) \circ X(a, t)-\mathbb{U}\left(\nabla_{x}\left((v \otimes v) \circ X^{-1}\right)\right) \circ X(a, t) . \tag{13}
\end{gather*}
$$

In Lagrangian variables the system is

$$
\left\{\begin{align*}
X(a, t)=a+\int_{0}^{t} \mathcal{V}(X, \tau, a, s) d s  \tag{14}\\
\tau(a, t)=\sigma_{0}(a)+\int_{0}^{t} \mathcal{T}(X, \tau, a, s) d s \\
v(a, t)=\mathcal{V}(X, \tau, t)
\end{align*}\right.
$$

where the Lagrangian nonlinearities $\mathcal{V}, \mathcal{T}$ are

$$
\left\{\begin{array}{c}
\mathcal{V}(X, \tau, a, s)=\mathbb{L}_{\nu}\left(u_{0}\right) \circ X(a, s)+\left(\mathbb{U}\left((\tau-v \otimes v) \circ X^{-1}\right)\right) \circ X(a, s)  \tag{15}\\
\mathcal{T}(X, \tau, a, s)=\left(g \tau+\tau g^{T}-2 k \tau+2 \rho K\left(g+g^{T}\right)\right)(a, s)
\end{array}\right.
$$

and $g$ is defined above in (13). The main result of the paper is
Theorem 1. Let $0<\alpha<1$ and $1<p<\infty$, be given. Let also $v_{1}(0)=u_{1}(0) \in C^{1+\alpha, p}$ and $v_{2}(0)=u_{2}(0) \in C^{1+\alpha, p}$ be given divergence-free initial velocities, and $\sigma_{1}(0), \sigma_{2}(0) \in C^{\alpha, p}$ be given initial stresses. Then there exists $T_{0}>0$ and $C>0$ depending on the norms of the initial data such that $\left(X_{1}, \tau_{1}, v_{1}\right),\left(X_{2}, \tau_{2}, v_{2}\right)$, with initial data $\left(I d, \sigma_{1}(0), u_{1}(0)\right),\left(I d, \sigma_{2}(0), u_{2}(0)\right)$, are bounded in $I d+\operatorname{Lip}\left(0, T_{0} ; C^{1+\alpha, p}\right) \times \operatorname{Lip}\left(0, T_{0} ; C^{\alpha, p}\right) \times L^{\infty}\left(0, T_{0} ; C^{1+\alpha, p}\right)$ and solve the Lagrangian form (14) of (6). Moreover,

$$
\begin{gather*}
\left\|X_{2}-X_{1}\right\|_{L i p\left(0, T_{0} ; C^{1+\alpha, p}\right)}+\left\|\tau_{2}-\tau_{1}\right\|_{L i p\left(0, T_{0} ; C^{\alpha, p}\right)}+\left\|v_{2}-v_{1}\right\|_{L^{\infty}\left(0, T_{0}, C^{1+\alpha, p}\right)} \\
\leq C\left(\left\|u_{2}(0)-u_{1}(0)\right\|_{1+\alpha, p}+\left\|\tau_{2}(0)-\tau_{1}(0)\right\|_{\alpha, p}\right) \tag{16}
\end{gather*}
$$

Remark 1. The solutions' Lagrangian stresses $\tau$ are Lipschitz in time with values in $C^{\alpha}$. Their Lagrangian counterparts $\sigma=\tau \circ X^{-1}$ are bounded in time with values in $C^{\alpha}$ and space-time Hölder continuous with exponent $\alpha$. The Eulerian version of the equations (6) is satisfied in the sense of distributions, and solutions are unique in this class.

The spaces $C^{\alpha, p}$ are defined in the next section. The proof of the theorem occupies the rest of the paper. We start by considering variations of Lagrangian variables. We take a family $\left(X_{\epsilon}, \tau_{\epsilon}\right)$ of flow maps depending smoothly on a parameter $\epsilon \in[1,2]$, with initial data $u_{\epsilon, 0}$ and $\sigma_{\epsilon, 0}$. Note that $v_{\epsilon}=\partial_{t} X_{\epsilon}$. We use the following notations

$$
\left\{\begin{array}{c}
u_{\epsilon}=\partial_{t} X_{\epsilon} \circ X_{\epsilon}^{-1}, g_{\epsilon}^{\prime}=\frac{d}{d \epsilon} g_{\epsilon}  \tag{17}\\
X_{\epsilon}^{\prime}=\frac{d}{d \epsilon} X_{\epsilon}, \eta_{\epsilon}=X_{\epsilon}^{\prime} \circ X_{\epsilon}^{-1} \\
v_{\epsilon}^{\prime}=\frac{d}{d \epsilon} v_{\epsilon} \\
\sigma_{\epsilon}=\tau_{\epsilon} \circ X_{\epsilon}^{-1} \\
\tau_{\epsilon}^{\prime}=\frac{d}{d \epsilon} \tau_{\epsilon}, \delta_{\epsilon}=\tau_{\epsilon}^{\prime} \circ X_{\epsilon}^{-1}
\end{array}\right.
$$

and

$$
\begin{equation*}
u_{\epsilon, 0}^{\prime}=\frac{d}{d \epsilon} u_{\epsilon}(0), \sigma_{\epsilon, 0}^{\prime}=\frac{d}{d \epsilon} \sigma_{\epsilon}(0) \tag{18}
\end{equation*}
$$

We represent

$$
\left\{\begin{align*}
X_{2}(a, t)-X_{1}(a, t) & =\int_{1}^{2} \mathcal{X}_{\epsilon}^{\prime} d \epsilon  \tag{19}\\
\tau_{2}(a, t)-\tau_{1}(a, t) & =\int_{1}^{2} \pi_{\epsilon} d \epsilon \\
v_{2}(a, t)-v_{1}(a, t)= & \int_{1}^{2} \frac{d}{d \epsilon} \mathcal{V}_{\epsilon} d \epsilon
\end{align*}\right.
$$

where

$$
\begin{gather*}
\mathcal{X}_{\epsilon}^{\prime}=\int_{0}^{t} \frac{d}{d \epsilon} \mathcal{V}_{\epsilon} d s, \pi_{\epsilon}=\int_{0}^{t} \frac{d}{d \epsilon} \mathcal{T}_{\epsilon} d s+\sigma_{\epsilon, 0}^{\prime}  \tag{20}\\
\mathcal{V}_{\epsilon}=\mathcal{V}\left(X_{\epsilon}, \tau_{\epsilon}\right), \mathcal{T}_{\epsilon}=\mathcal{T}\left(X_{\epsilon}, \tau_{\epsilon}\right)
\end{gather*}
$$

We have the following commutator expressions arising by differentiating in $\epsilon$ ([1], [2])):

$$
\begin{equation*}
\left(\frac{d}{d \epsilon}\left(\mathbb{U}\left(\tau_{\epsilon} \circ X_{\epsilon}^{-1}\right) \circ X_{\epsilon}\right)\right) \circ X_{\epsilon}^{-1}=\left[\eta_{\epsilon} \cdot \nabla_{x}, \mathbb{U}\right]\left(\sigma_{\epsilon}\right)+\mathbb{U}\left(\delta_{\epsilon}\right), \tag{21}
\end{equation*}
$$

where

$$
\begin{equation*}
\left[\eta_{\epsilon} \cdot \nabla_{x}, \mathbb{U}\right]\left(\sigma_{\epsilon}\right)=\eta_{\epsilon} \cdot \nabla_{x}\left(\mathbb{U}\left(\sigma_{\epsilon}\right)\right)-\mathbb{U}\left(\eta_{\epsilon} \cdot \nabla_{x} \sigma_{\epsilon}\right) \tag{22}
\end{equation*}
$$

and

$$
\begin{gather*}
\left(\frac{d}{d \epsilon} \mathbb{U}\left(v_{\epsilon} \otimes v_{\epsilon} \circ X_{\epsilon}^{-1}\right) \circ X_{\epsilon}\right) \circ X_{\epsilon}^{-1}  \tag{23}\\
=\left[\eta_{\epsilon} \cdot \nabla_{x}, \mathbb{U}\right]\left(u_{\epsilon} \otimes u_{\epsilon}\right)+\mathbb{U}\left(\left(v_{\epsilon}^{\prime} \otimes v_{\epsilon}+v_{\epsilon} \otimes v_{\epsilon}^{\prime}\right) \circ X_{\epsilon}^{-1}\right) .
\end{gather*}
$$

We note, by the chain rule,

$$
\begin{equation*}
\nabla_{a} \mathcal{V}=\left(\nabla_{a} X\right) g \tag{24}
\end{equation*}
$$

Consequently, differentiating $\mathcal{V}_{\epsilon}, g_{\epsilon}$ and the relation (24) we have

$$
\left\{\begin{array}{c}
\left(\frac{d}{d \epsilon} \mathcal{V}_{\epsilon}\right) \circ X_{\epsilon}^{-1}=\eta_{\epsilon} \cdot\left(\mathbb{L}_{\nu}\left(\nabla_{x} u_{\epsilon, 0}\right)\right)+\mathbb{L}_{\nu}\left(u_{\epsilon, 0}^{\prime}\right)  \tag{25}\\
+\left[\eta_{\epsilon} \cdot \nabla_{x}, \mathbb{U}\right]\left(\sigma_{\epsilon}-u_{\epsilon} \otimes u_{\epsilon}\right)+\mathbb{U}\left(\delta_{\epsilon}-\left(v_{\epsilon}^{\prime} \otimes v_{\epsilon}+v_{\epsilon} \otimes v_{\epsilon}^{\prime}\right) \circ X_{\epsilon}^{-1}\right) \\
g_{\epsilon}=\mathbb{L}\left(\nabla_{x} u_{\epsilon, 0}\right) \circ X_{\epsilon}+\mathbb{G}\left(\sigma_{\epsilon}\right) \circ X_{\epsilon}-\mathbb{U}\left(\nabla_{x}\left(u_{\epsilon} \otimes u_{\epsilon}\right)\right) \circ X_{\epsilon} \\
g_{\epsilon}^{\prime} \circ X_{\epsilon}^{-1}=\eta_{\epsilon} \cdot \mathbb{L}_{\nu}\left(\nabla_{x} \nabla_{x} u_{\epsilon, 0}\right)+\mathbb{L}_{\nu}\left(\nabla_{x} u_{\epsilon, 0}^{\prime}\right)+\left[\eta_{\epsilon} \cdot \nabla_{x}, \mathbb{G}\right]\left(\sigma_{\epsilon}\right)+\mathbb{G}\left(\delta_{\epsilon}\right) \\
-\left[\eta_{\epsilon} \cdot \nabla_{x}, \mathbb{U}\right]\left(\nabla_{x}\left(u_{\epsilon} \otimes u_{\epsilon}\right)\right)-\mathbb{U}\left(\nabla_{x}\left(\left(v_{\epsilon}^{\prime} \otimes v_{\epsilon}+v_{\epsilon} \otimes v_{\epsilon}^{\prime}\right) \circ X_{\epsilon}^{-1}\right)\right) \\
\frac{d}{d \epsilon}\left(\nabla_{a} \mathcal{V}_{\epsilon}\right)=\left(\nabla_{a} X_{\epsilon}^{\prime}\right) g_{\epsilon}+\left(\nabla_{a} X_{\epsilon}\right) g_{\epsilon}^{\prime} \\
\frac{d}{d \epsilon} \mathcal{T}_{\epsilon}=g_{\epsilon}^{\prime} \tau_{\epsilon}+g_{\epsilon} \tau_{\epsilon}^{\prime}+\tau_{\epsilon}^{\prime} g_{\epsilon}^{T}+\tau_{\epsilon}\left(g_{\epsilon}^{\prime}\right)^{T}-2 k \tau_{\epsilon}^{\prime}+2 \rho K\left(g_{\epsilon}^{\prime}+\left(g_{\epsilon}^{\prime}\right)^{T}\right)
\end{array}\right.
$$

## 3 Functions, operators, commutators

We consider function spaces

$$
\begin{equation*}
C^{\alpha, p}=C^{\alpha}\left(\mathbb{R}^{d}\right) \cap L^{p}\left(\mathbb{R}^{d}\right) \tag{26}
\end{equation*}
$$

with norm

$$
\begin{equation*}
\|f\|_{\alpha, p}=\|f\|_{C^{\alpha}\left(\mathbb{R}^{d}\right)}+\|f\|_{L^{p}\left(\mathbb{R}^{d}\right)} \tag{27}
\end{equation*}
$$

for $\alpha \in(0,1), p \in(1, \infty), C^{1+\alpha}\left(\mathbb{R}^{d}\right)$ with norm

$$
\begin{equation*}
\|f\|_{C^{1+\alpha}\left(\mathbb{R}^{d}\right)}=\|f\|_{L^{\infty}\left(\mathbb{R}^{d}\right)}+\|\nabla f\|_{C^{\alpha}\left(\mathbb{R}^{d}\right)} \tag{28}
\end{equation*}
$$

and

$$
\begin{equation*}
C^{1+\alpha, p}=C^{1+\alpha}\left(\mathbb{R}^{d}\right) \cap W^{1, p}\left(\mathbb{R}^{d}\right) \tag{29}
\end{equation*}
$$

with norm

$$
\begin{equation*}
\|f\|_{1+\alpha, p}=\|f\|_{C^{1+\alpha}\left(\mathbb{R}^{d}\right)}+\|f\|_{W^{1, p}\left(\mathbb{R}^{d}\right)} \tag{30}
\end{equation*}
$$

We also use spaces of paths, $L^{\infty}(0, T ; Y)$ with the usual norm,

$$
\begin{equation*}
\|f\|_{L^{\infty}(0, T ; Y)}=\sup _{t \in[0, T]}\|f(t)\|_{Y} \tag{31}
\end{equation*}
$$

spaces $\operatorname{Lip}(0, T ; Y)$ with norm

$$
\begin{equation*}
\|f\|_{L i p(0, T ; Y)}=\sup _{t \neq s, t, s \in[0, T]} \frac{\|f(t)-f(s)\|_{Y}}{|t-s|}+\|f\|_{L^{\infty}(0, T ; Y)} \tag{32}
\end{equation*}
$$

where $Y$ is $C^{\alpha, p}$ or $C^{1+\alpha, p}$ in the following. We use the following lemmas.
Lemma 1 ([2]). Let $0<\alpha<1,1<p<\infty$. Let $\eta \in C^{1+\alpha}\left(\mathbb{R}^{d}\right)$ and let

$$
\begin{equation*}
(\mathbb{K} \sigma)(x)=P . V . \int_{\mathbb{R}^{d}} k(x-y) \sigma(y) d y \tag{33}
\end{equation*}
$$

be a classical Calderon-Zygmund operator with kernel $k$ which is smooth away from the origin, homogeneous of degree $-d$ and with mean zero on spheres about the origin. Then the commutator $[\eta \cdot \nabla, \mathbb{K}]$ can be defined as a bounded linear operator in $C^{\alpha, p}$ and

$$
\begin{equation*}
\|[\eta \cdot \nabla, \mathbb{K}] \sigma\|_{C^{\alpha, p}} \leq C\|\eta\|_{C^{1+\alpha}\left(\mathbb{R}^{d}\right)}\|\sigma\|_{C^{\alpha, p}} \tag{34}
\end{equation*}
$$

Lemma 2 (Generalized Young's inequality). Let $1 \leq q \leq \infty$ and $C>0$. Suppose $K$ is a measurable function on $\mathbb{R}^{d} \times \mathbb{R}^{d}$ such that

$$
\begin{equation*}
\sup _{x \in \mathbb{R}^{d}} \int_{\mathbb{R}^{d}}|K(x, y)| d y \leq C, \sup _{y \in \mathbb{R}^{d}} \int_{\mathbb{R}^{d}}|K(x, y)| d x \leq C \tag{35}
\end{equation*}
$$

If $f \in L^{q}\left(\mathbb{R}^{d}\right)$, the function $T f$ defined by

$$
\begin{equation*}
T f(x)=\int_{\mathbb{R}^{d}} K(x, y) f(y) d y \tag{36}
\end{equation*}
$$

is well defined almost everywhere and is in $L^{q}$, and $\|T f\|_{L^{q}} \leq C\|f\|_{L^{q}}$.
The proof of this lemma for $1<q<\infty$ is done using duality, a straightforward application of Young's inequality and changing order of integration. The extreme cases $q=1$ and $q=\infty$ are proved directly by inspection.

For simplicity of notation, let us denote

$$
\begin{equation*}
M_{X}=1+\|X-\mathrm{Id}\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} \tag{37}
\end{equation*}
$$

Theorem 2. Let $0<\alpha<1,1<p<\infty$ and let $T>0$. Also let $X$ be a volume preserving diffeomorphism such that $X-\operatorname{Id} \in \operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)$. Then

$$
\begin{equation*}
\left\|\tau \circ X^{-1}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} M_{X}^{\alpha} \tag{38}
\end{equation*}
$$

If $X^{\prime} \in \operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)$, then

$$
\begin{equation*}
\left\|X^{\prime} \circ X^{-1}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} \leq\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} M_{X}^{1+2 \alpha} \tag{39}
\end{equation*}
$$

If $v \in \operatorname{Lip}\left(0, T ; W^{1, p}\right)$, then

$$
\begin{equation*}
\left\|v \circ X^{-1}\right\|_{L^{\infty}\left(0, T ; W^{1, p}\right)} \leq\|v\|_{L^{\infty}\left(0, T ; W^{1, p}\right)} M_{X} \tag{40}
\end{equation*}
$$

If in addition $\partial_{t} X^{\prime}, \partial_{t} X$ exist in $L^{\infty}\left(0, T ; C^{1+\alpha}\right)$, then

$$
\begin{equation*}
\left\|X^{\prime} \circ X^{-1}\right\|_{\operatorname{Lip(0,T;C^{\alpha })}} \leq\left\|X^{\prime}\right\|_{\operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)}\|X-\mathrm{Id}\|_{\operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)} M_{X}^{1+3 \alpha} \tag{41}
\end{equation*}
$$

Proof.

$$
\begin{equation*}
\left\|\tau \circ X^{-1}\right\|_{L^{p} \cap L^{\infty}}=\|\tau\|_{L^{p} \cap L^{\infty}} \tag{42}
\end{equation*}
$$

and, denoting the seminorm

$$
[\tau]_{\alpha}=\sup _{a \neq b, a, b \in \mathbb{R}^{2}} \frac{|\tau(a)-\tau(b)|}{|a-b|^{\alpha}}
$$

we have

$$
\begin{equation*}
\left[\tau \circ X^{-1}(t)\right]_{\alpha} \leq[\tau(t)]_{\alpha}\left\|\nabla_{x} X^{-1}(t)\right\|_{L^{\infty}}^{\alpha} \leq[\tau(t)]_{\alpha}\left(1+\|X-\mathrm{Id}\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\right)^{\alpha} \tag{43}
\end{equation*}
$$

Note that this shows that the same bound holds when we replace $X^{-1}$ by $X$. For the second and third part, it suffices to remark that

$$
\begin{equation*}
\nabla_{x}\left(X^{\prime} \circ X^{-1}\right)=\left(\left(\nabla_{a} X\right) \circ X^{-1}\right)^{-1}\left(\left(\nabla_{a} X^{\prime}\right) \circ X^{-1}\right) \tag{44}
\end{equation*}
$$

and the previous part gives the bound in terms of Lagrangian variables. For the last part, we note that

$$
\begin{gather*}
\frac{1}{t-s}\left(X^{\prime}\left(X^{-1}(x, t), t\right)-X^{\prime}\left(X^{-1}(x, s), s\right)\right) \\
=\int_{0}^{1}\left(\left(\partial_{t} X^{\prime}\right)\left(X^{-1}\left(x, \beta_{\tau}\right), \beta_{\tau}\right)+\left(\partial_{t} X^{-1}\right)\left(x, \beta_{\tau}\right)\left(\nabla_{a} X^{\prime}\right)\left(X^{-1}\left(x, \beta_{\tau}\right), \beta_{\tau}\right)\right) d \tau \tag{45}
\end{gather*}
$$

where

$$
\begin{equation*}
\beta_{\tau}=\tau t+(1-\tau) s \tag{46}
\end{equation*}
$$

Now noting that

$$
\begin{equation*}
\partial_{t} X^{-1}=-\left(\left(\partial_{t} X\right) \circ X^{-1}\right)\left(\left(\nabla_{a} X\right)^{-1} \circ X^{-1}\right) \tag{47}
\end{equation*}
$$

we have

$$
\begin{gather*}
\frac{1}{t-s}\left\|X^{\prime} \circ X^{-1}(t)-X^{\prime} \circ X^{-1}(s)\right\|_{C^{\alpha}}  \tag{48}\\
\leq\left(\left\|\partial_{t} X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{\alpha}\right)}+\left\|\partial_{t} X\right\|_{L^{\infty}\left(0, T ; C^{\alpha}\right)}\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\right)\left(1+\|X-\mathrm{Id}\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\right)^{1+3 \alpha}
\end{gather*}
$$

so that

$$
\begin{equation*}
\left\|X^{\prime} \circ X^{-1}\right\|_{L i p\left(0, T ; C^{\alpha}\right)} \leq\left\|X^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\|X-\mathrm{Id}\|_{\operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)}\left(1+\|X-\mathrm{Id}\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\right)^{1+3 \alpha} \tag{49}
\end{equation*}
$$

Theorem 3. Let $0<\alpha<1,1<p<\infty$ and let $T>0$. There exists a constant $C$ independent of $T$ and $\nu$ such that for any $0<t<T$,

$$
\begin{align*}
&\left\|\mathbb{L}_{\nu}\left(u_{0}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left\|u_{0}\right\|_{\alpha, p} \\
&\left\|\mathbb{L}_{\nu}\left(u_{0}\right)\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha, p}\right)} \leq C\left\|u_{0}\right\|_{1+\alpha, p} \\
&\left\|\mathbb{L}_{\nu}\left(\nabla u_{0}\right)(t)\right\|_{\alpha, p} \leq \frac{C}{(\nu t)^{\frac{1}{2}}}\left\|u_{0}\right\|_{\alpha, p}  \tag{50}\\
&\left\|\mathbb{L}_{\nu}\left(\nabla u_{0}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left\|u_{0}\right\|_{1+\alpha, p}
\end{align*}
$$

hold.
Proof.

$$
\begin{gather*}
\left\|\mathbb{L}_{\nu}\left(u_{0}\right)(t)\right\|_{\alpha, p} \leq\left\|g_{\nu t}\right\|_{L^{1}}\left\|u_{0}\right\|_{\alpha, p}=\left\|u_{0}\right\|_{\alpha, p} \\
\left\|\mathbb{L}_{\nu}\left(u_{0}\right)(t)\right\|_{1+\alpha, p} \leq\left\|g_{\nu t}\right\|_{L^{1}}\left\|u_{0}\right\|_{1+\alpha, p}=\left\|u_{0}\right\|_{1+\alpha, p} \\
\left\|\mathbb{L}_{\nu}\left(\nabla u_{0}\right)(t)\right\|_{\alpha, p} \leq\left\|\nabla g_{\nu t}\right\|_{L^{1}}\left\|u_{0}\right\|_{1+\alpha, p}=\frac{C}{(\nu t)^{\frac{1}{2}}}\left\|u_{0}\right\|_{\alpha, p},  \tag{51}\\
\left\|\mathbb{L}_{\nu}\left(\nabla u_{0}\right)(t)\right\|_{\alpha, p} \leq\left\|g_{\nu t}\right\|_{L^{1}}\left\|\nabla u_{0}\right\|_{\alpha, p} \leq\left\|u_{0}\right\|_{1+\alpha, p} .
\end{gather*}
$$

Theorem 4. Let $0<\alpha<1,1<p<\infty$ and let $T>0$. There exists a constant $C$ such that

$$
\begin{equation*}
\|\mathbb{U}(\sigma)\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left(\frac{T}{\nu}\right)^{\frac{1}{2}}\|\sigma\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{52}
\end{equation*}
$$

Proof.

$$
\begin{gather*}
\|\mathbb{U}(\sigma)(t)\|_{C^{\alpha, p}} \leq C \int_{0}^{t}\left\|\nabla g_{\nu(t-s)}\right\|_{L^{1}}\|\sigma(s)\|_{\alpha, p} d s \\
\leq \frac{C}{\nu^{\frac{1}{2}}} \int_{0}^{t} \frac{1}{(t-s)^{\frac{1}{2}}} d s\|\sigma\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq \frac{C}{\nu^{\frac{1}{2}}} \sqrt{T}\|\sigma\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{53}
\end{gather*}
$$

Theorem 5. Let $0<\alpha<1,1<p<\infty$ and let $T>0$. There exist constants $C_{1}, C_{2}$ depending only on $\alpha$ and $\nu$, and $C_{3}(T, X), C_{4}(T, X)$ such that

$$
\begin{gather*}
\left\|\mathbb{G}\left(\tau \circ X^{-1}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C_{1}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\tau(0)\|_{\alpha, p}\left(1+C_{3}(T, X)\right) \\
+C_{2}\|\tau\|_{\operatorname{Lip}\left(0, T ; C^{\alpha, p}\right)} C_{4}(T, X) \tag{54}
\end{gather*}
$$

where $C_{3}(T, X)$ and $C_{4}(T, X)$ are of the form $C T^{\frac{1}{2}}\left(\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}+\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{4}\right)$.
Proof. Since $\mathbb{G}=(R \otimes R) \mathbb{H} \Gamma$ where

$$
\begin{equation*}
\Gamma\left(\tau \circ X^{-1}\right)=\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\tau \circ X^{-1}(s)\right) d s \tag{55}
\end{equation*}
$$

we can replace $\mathbb{G}$ by $\Gamma$. Then $\Gamma\left(\tau \circ X^{-1}\right)$ can be written as

$$
\begin{align*}
\Gamma\left(\tau \circ X^{-1}\right)(t)= & \int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\left(\tau \circ X^{-1}\right)(s)-\left(\tau \circ X^{-1}\right)(t)\right) d s  \tag{56}\\
& +\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\tau \circ X^{-1}\right)(t) d s
\end{align*}
$$

But

$$
\begin{equation*}
\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\tau \circ X^{-1}\right)(t) d s=\tau \circ X^{-1}(t)-g_{\nu t} *\left(\tau \circ X^{-1}\right)(t) \tag{57}
\end{equation*}
$$

so the second term is bounded by $2\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} M_{X}^{\alpha}$ by Theorem 2 . Now we let

$$
\begin{equation*}
\tau \circ X^{-1}(x, s)-\tau \circ X^{-1}(x, t)=\Delta_{1} \tau(x, s, t)+\Delta_{2} \tau(x, s, t) \tag{58}
\end{equation*}
$$

where

$$
\begin{align*}
\Delta_{1} \tau(x, s, t) & =\tau\left(X^{-1}(x, s), s\right)-\tau\left(X^{-1}(x, s), t\right)  \tag{59}\\
\Delta_{2} \tau(x, s, t) & =\tau\left(X^{-1}(x, s), t\right)-\tau\left(X^{-1}(x, t), t\right)
\end{align*}
$$

But since

$$
\begin{equation*}
\left\|\Delta_{1} \tau(s, t)\right\|_{C^{\alpha, p}} \leq|t-s| M_{X}^{\alpha}\|\tau\|_{L i p\left(0, T ; C^{\alpha, p}\right)} \tag{60}
\end{equation*}
$$

by the proof of Theorem 2 we get

$$
\begin{equation*}
\left\|\int_{0}^{t} \Delta g_{\nu(t-s)} * \Delta_{1} \tau(s, t) d s\right\|_{\alpha, p} \leq \frac{C t}{\nu}\|\tau\|_{\operatorname{Lip(0,T;C^{\alpha ,p})}} M_{X}^{\alpha} \tag{61}
\end{equation*}
$$

On the other hand,

$$
\begin{equation*}
\int_{0}^{t} \Delta g_{\nu(t-s)} * \Delta_{2} \tau(s, t) d s=\int_{0}^{t} \int_{\mathbb{R}^{d}} K(x, z, t, s) \tau(z, t) d z d s \tag{62}
\end{equation*}
$$

where

$$
\begin{equation*}
K(x, z, t, s)=\Delta g_{\nu(t-s)}(x-X(z, s))-\Delta g_{\nu(t-s)}(x-X(z, t)) \tag{63}
\end{equation*}
$$

We use the following lemma.
Lemma 3. $K(x, z, t, s)$ is $L^{1}$ in both the $x$ variable and the $z$ variable, and

$$
\begin{equation*}
\sup _{z}\|K(\cdot, z, t, s)\|_{L^{1}}, \sup _{x}\|K(x, \cdot, t, s)\|_{L^{1}} \leq \frac{C\|X-\mathrm{Id}\|_{L i p\left(0, T ; L^{\infty}\right)}}{|t-s|^{\frac{1}{2}} \nu^{\frac{3}{2}}} \tag{64}
\end{equation*}
$$

Proof. We define

$$
\begin{equation*}
S(x)=4 \pi e^{-|x|^{2}}\left(|x|^{2}-\frac{d}{2}\right) \tag{65}
\end{equation*}
$$

so that

$$
\begin{equation*}
\left(\Delta g_{\nu(t-s)}\right)=(4 \pi \nu(t-s))^{-\left(\frac{d}{2}+1\right)} S\left(\frac{x}{(4(t-s))^{\frac{1}{2}}}\right) \tag{66}
\end{equation*}
$$

Then

$$
\begin{gather*}
\int|K(x, z, t, s)| d z=\int(4 \pi \nu(t-s))^{-\left(\frac{d}{2}+1\right)}\left|S\left(\frac{x-X(z, s)}{(4 \nu(t-s))^{\frac{1}{2}}}\right)-S\left(\frac{x-X(z, t)}{(4 \nu(t-s))^{\frac{1}{2}}}\right)\right| d z \\
=\int(4 \pi \nu(t-s))^{-\left(\frac{d}{2}+1\right)}\left|S\left(\frac{x-y}{(4 \nu(t-s))^{\frac{1}{2}}}\right)-S\left(\frac{x-X(y, t-s)}{(4 \nu(t-s))^{\frac{1}{2}}}\right)\right| d y  \tag{67}\\
= \\
(4 \pi \nu(t-s))^{-1} \pi^{-\left(\frac{d}{2}+1\right)} \int\left|S(u)-S\left(u-\frac{(X-\mathrm{Id})\left(x-(4(t-s))^{\frac{1}{2}} u, t-s\right)}{(4 \nu(t-s))^{\frac{1}{2}}}\right)\right| d u .
\end{gather*}
$$

However, for each $u$

$$
\begin{align*}
\mid S(u)-S(u & \left.-\frac{(X-\mathrm{Id})\left(x-(4 \nu(t-s))^{\frac{1}{2}} u, t-s\right)}{(4 \nu(t-s))^{\frac{1}{2}}}\right)\left|\leq\left|\frac{(X-\mathrm{Id})\left(x-(4 \nu(t-s))^{\frac{1}{2}} u, t-s\right)}{(4 \nu(t-s))^{\frac{1}{2}}}\right|\right.  \tag{68}\\
& \times \sup \left\{|\nabla S(u-z)|:|z| \leq\left|\frac{(X-\mathrm{Id})\left(x-(4 \nu(t-s))^{\frac{1}{2}} u, t-s\right)}{(4 \nu(t-s))^{\frac{1}{2}}}\right|\right\}
\end{align*}
$$

and we have

$$
\begin{equation*}
\left|\frac{(X-\mathrm{Id})\left(x-(4 \nu(t-s))^{\frac{1}{2}} u, t-s\right)}{(4 \nu(t-s))^{\frac{1}{2}}}\right| \leq\|(X-\mathrm{Id})\|_{L i p\left(0, T ; L^{\infty}\right)} \frac{|t-s|^{\frac{1}{2}}}{\nu^{\frac{1}{2}}} \leq C T^{\frac{1}{2}} \tag{69}
\end{equation*}
$$

and obviously

$$
\begin{equation*}
\tilde{S}(u)=\sup _{z \leq C T^{\frac{1}{2}}}|(\nabla S)(u-z)| \tag{70}
\end{equation*}
$$

is integrable in $\mathbb{R}^{d}$; because $\nabla S$ is Schwartz,

$$
\begin{equation*}
|(\nabla S)(x)| \leq \frac{C_{d}}{\left(1+2 C^{2} T+|x|^{2}\right)^{d}} \tag{71}
\end{equation*}
$$

for some constant $C_{d}$, but if $|z| \leq C T^{\frac{1}{2}}$, then $|u-z|^{2} \geq|u|^{2}-C^{2} T$ and

$$
\begin{equation*}
|(\nabla S)(u-z)| \leq \frac{C_{d}}{\left(1+C^{2} T+|u|^{2}\right)^{d}} \tag{72}
\end{equation*}
$$

and the right side of above is clearly integrable with bound depending only on $d$ and $T$. Therefore, we have

$$
\begin{equation*}
\int|K(x, z, t, s)| d z \leq|t-s|^{-\frac{1}{2}} \nu^{-\frac{3}{2}}\|(X-\mathrm{Id})\|_{L i p\left(0, T ; L^{\infty}\right)} C(d, T) \tag{73}
\end{equation*}
$$

Similarly,

$$
\begin{gather*}
\int|K(x, z, t, s)| d x=\int(4 \pi \nu(t-s))^{-\left(\frac{d}{2}+1\right)}\left|S\left(\frac{x-X(z, s)}{(4 \nu(t-s))^{\frac{1}{2}}}\right)-S\left(\frac{x-X(z, t)}{(4 \nu(t-s))^{\frac{1}{2}}}\right)\right| d x \\
=\int(4 \pi \nu(t-s))^{-1} \pi^{-\left(\frac{d}{2}+1\right)}\left|S(y)-S\left(y+\frac{X(z, s)-X(z, t)}{(4 \nu(t-s))^{\frac{1}{2}}}\right)\right| d y \tag{74}
\end{gather*}
$$

and again we have

$$
\begin{equation*}
\left|\frac{X(z, s)-X(z, t)}{(4 \nu(t-s))^{\frac{1}{2}}}\right| \leq\|(X-\mathrm{Id})\|_{L i p\left(0, T ; L^{\infty}\right)}|t-s|^{\frac{1}{2}} \nu^{-\frac{1}{2}} \leq C T^{\frac{1}{2}} \tag{75}
\end{equation*}
$$

Therefore, we have the bound

$$
\begin{equation*}
\int|K(x, z)| d x \leq|t-s|^{-\frac{1}{2}} \nu^{-\frac{3}{2}}\|(X-\mathrm{Id})\|_{L i p\left(0, T ; L^{\infty}\right)} C(d, T) \tag{76}
\end{equation*}
$$

From Lemma 3 and generalized Young's inequality, we have

$$
\begin{equation*}
\left\|\int_{0}^{t} \Delta g_{\nu(t-s)} * \Delta_{2} \tau(s, t) d s\right\|_{L^{p} \cap L^{\infty}} \leq \frac{C}{\nu}\left(\left(\frac{t}{\nu}\right)^{\frac{1}{2}}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\right)\|\tau\|_{L^{\infty}\left(0, T ; L^{p} \cap L^{\infty}\right)} \tag{77}
\end{equation*}
$$

For the Hölder seminorm, we measure the finite difference. Let us denote $\delta_{h} f(x, t)=f(x+h, t)-f(x, t)$. If $|h|<t$, then

$$
\begin{equation*}
\delta_{h}\left(\int_{0}^{t} \Delta g_{\nu(t-s)} * \Delta_{2} \tau(s, t) d s\right)=\int_{0}^{t} \delta_{h}\left(\Delta g_{\nu(t-s)}\right) * \Delta_{2} \tau(s, t) d s \tag{78}
\end{equation*}
$$

If $0<t-s<|h|$, then $\left\|\delta_{h} \Delta g_{\nu(t-s)}\right\|_{L^{1}} \leq 2\left\|\Delta g_{\nu(t-s)}\right\|_{L^{1}} \leq \frac{C}{\nu(t-s)}$ and since

$$
\begin{equation*}
\left\|\Delta_{2} \tau(s, t)\right\|_{L^{\infty}} \leq|t-s|^{\alpha}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{79}
\end{equation*}
$$

we have

$$
\begin{equation*}
\left\|\int_{t-|h|}^{t} \delta_{h}\left(\Delta g_{\nu(t-s)}\right) * \Delta_{2} \tau(s, t) d s\right\|_{L^{\infty}} \leq \frac{C}{\nu \alpha}|h|^{\alpha}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{80}
\end{equation*}
$$

If $|h|<t-s<t$, then following lines of Lemma $3 \delta_{h}\left(\Delta g_{\nu(t-s)}\right)$ is a $L^{1}$ function with

$$
\begin{equation*}
\left\|\delta_{h}\left(\Delta g_{\nu(t-s)}\right)\right\|_{L^{1}} \leq \frac{C|h|}{(\nu(t-s))^{\frac{3}{2}}} \tag{81}
\end{equation*}
$$

and we have

$$
\begin{gather*}
\left\|\int_{0}^{t-|h|} \delta_{h}\left(\Delta g_{\nu(t-s)}\right) * \Delta_{2} \tau(s, t) d s\right\|_{L^{\infty}} \\
\leq \begin{cases}\frac{C}{\nu^{\frac{3}{2}}}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}|h|^{\frac{1}{2} \frac{t^{\alpha}}{\alpha}} & \alpha \leq \frac{1}{2}, \\
\frac{C}{\nu^{\frac{3}{2}}}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}|h| \frac{t^{\alpha-\frac{1}{2}}}{\alpha-\frac{1}{2}} & \alpha>\frac{1}{2} .\end{cases} \tag{82}
\end{gather*}
$$

If $|h| \geq t$, then we only have the first term. Therefore, we have

$$
\begin{equation*}
\frac{1}{|h|^{\alpha}}\left\|\delta_{h}\left(\int_{0}^{t} \Delta g_{\nu(t-s)} * \Delta_{2} \tau(s, t) d s\right)\right\|_{L^{\infty}} \leq \frac{C(\alpha)}{\nu}\|X-\operatorname{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{83}
\end{equation*}
$$

We note that

$$
\begin{equation*}
\|\tau(t)\|_{\alpha, p} \leq\|\tau(0)\|_{\alpha, p}+t\|\tau\|_{\operatorname{Lip}\left(0, T ; C^{\alpha, p}\right)} \tag{84}
\end{equation*}
$$

To summarize, we have

$$
\begin{gather*}
\left\|\Gamma\left(\tau \circ X^{-1}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \\
\leq C(\alpha)\left(1+\frac{1}{\nu}\right)\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\tau(0)\|_{\alpha, p}+C(\alpha)\left(1+\frac{1}{\nu}\right)\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha} T\|\tau\|_{L i p\left(0, T ; C^{\alpha, p}\right)} \\
+\frac{C(\alpha)}{\nu}\left(\frac{T}{\nu}\right)^{\frac{1}{2}} \max \left\{\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha},\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{4}\right\}\left(\|\tau(0)\|_{\alpha, p}+T\|\tau\|_{L i p\left(0, T ; C^{\alpha, p}\right)}\right), \tag{85}
\end{gather*}
$$

and this completes the proof.
Theorem 6. Let $0<\alpha<1,1<p<\infty$ and let $T>0$. Let $X^{\prime} \in \operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)$ with $\partial_{t} X^{\prime} \in$ $L^{\infty}\left(0, T ; C^{1+\alpha}\right)$. There exists a constant $C$ such that

$$
\begin{gather*}
\left\|\left[X^{\prime} \circ X^{-1} \cdot \nabla, \mathbb{U}\right](\sigma)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \\
\leq C\left(\left(\frac{T}{\nu}\right)^{\frac{1}{2}}+\frac{T}{\nu}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\right) M_{X}^{1+3 \alpha}\left\|X^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\|\sigma\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{86}
\end{gather*}
$$

Proof. First, we denote

$$
\begin{equation*}
\eta=X^{\prime} \circ X^{-1} \tag{87}
\end{equation*}
$$

Then we have

$$
\begin{gather*}
{[\eta \cdot \nabla, \mathbb{U}](\sigma)(t)} \\
=\eta(t) \cdot \nabla \int_{0}^{t} g_{\nu(t-s)} * \mathbb{H} \operatorname{div} \sigma(s) d s-\int_{0}^{t} g_{\nu(t-s)} * \mathbb{H} \operatorname{div}(\eta(s) \cdot \nabla \sigma(s)) d s \\
=[\eta(t) \cdot \nabla, \mathbb{H}] \int_{0}^{t} g_{\nu(t-s)} * \operatorname{div} \sigma(s) d s+\mathbb{H} \int_{0}^{t}\left(\nabla g_{\nu(t-s)}\right) *(\nabla \cdot \eta(s) \sigma(s)) d s  \tag{88}\\
-\mathbb{H} \int_{0}^{t}\left(\nabla \nabla g_{\nu(t-s)}\right) *(\eta(s)-\eta(t)) \sigma(s) d s \\
+\mathbb{H} \int_{0}^{t}\left(\eta(t) \cdot\left(\nabla \nabla g_{\nu(t-s)}\right) * \sigma(s)-\left(\nabla \nabla g_{\nu(t-s)}\right) *(\eta(t) \sigma(s))\right) d s,
\end{gather*}
$$

where $\left(\nabla \nabla g_{\nu(t-s)}\right) *(\eta(s)-\eta(t)) \sigma(s), \eta(t) \cdot\left(\nabla \nabla g_{\nu(t-s)}\right) * \sigma(s)$, and $\left(\nabla \nabla g_{\nu(t-s)}\right) *(\eta(s) \sigma(s))$ represent

$$
\begin{gather*}
\sum_{i, j}\left(\partial_{i} \partial_{j} g_{\nu(t-s)} *\right)\left(\eta_{i}(s)-\eta_{i}(t)\right) \sigma_{j k}(s) \\
\sum_{i, j} \eta_{i}(t)\left(\partial_{i} \partial_{j} g_{\nu(t-s)}\right) * \sigma_{j k}(s), \text { and respectively } \sum_{i, j}\left(\partial_{i} \partial_{j} g_{\nu(t-s)}\right) *\left(\eta_{i}(s) \sigma_{j k}(s)\right) \tag{89}
\end{gather*}
$$

The first term is bounded by Lemma 1 and the second term is estimated directly

$$
\begin{align*}
\left\|[\eta(t) \cdot \nabla, \mathbb{H}] \int_{0}^{t} g_{\nu(t-s)} * \operatorname{div} \sigma(s) d s\right\|_{\alpha, p} \leq C\|\eta(t)\|_{C^{1+\alpha}}\left(\frac{t}{\nu}\right)^{\frac{1}{2}}\|\sigma\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}  \tag{90}\\
\left\|\mathbb{H} \int_{0}^{t}\left(\nabla g_{\nu(t-s)}\right) *(\nabla \cdot \eta(s) \sigma(s)) d s\right\|_{\alpha, p} \leq C\left(\frac{t}{\nu}\right)^{\frac{1}{2}}\|\eta\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\sigma\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}
\end{align*}
$$

The third term is bounded by

$$
\begin{equation*}
\frac{C t}{\nu}\|\eta\|_{L i p\left(0, T ; C^{\alpha}\right)}\|\sigma\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{91}
\end{equation*}
$$

by the virtue of Theorem 2. For the last term, note that

$$
\begin{gather*}
\quad\left(\eta(t) \cdot\left(\nabla \nabla g_{\nu(t-s)}\right) * \sigma(s)-\left(\nabla \nabla g_{\nu(t-s)}\right) *(\eta(t) \sigma(s))\right)(x) \\
=\int_{\mathbb{R}^{d}} \nabla \nabla g_{\nu(t-s)}(z) z \cdot\left(\int_{0}^{1} \nabla \eta(x-(1-\lambda) z, t) d \lambda\right) \sigma(x-z, s) d z \tag{92}
\end{gather*}
$$

and note that $\nabla \nabla g_{\nu(t-s)}(z) z$ is a $L^{1}$ function with

$$
\begin{equation*}
\left\|\nabla \nabla g_{\nu(t-s)}(z) z\right\|_{L^{1}} \leq \frac{C}{(\nu(t-s))^{\frac{1}{2}}} \tag{93}
\end{equation*}
$$

Therefore,

$$
\begin{gather*}
\left\|\left(\eta(t) \cdot\left(\nabla \nabla g_{\nu(t-s)}\right) * \sigma(s)-\left(\nabla \nabla g_{\nu(t-s)}\right) *(\eta(t) \sigma(s))\right)\right\|_{\alpha, p} \\
\leq \frac{C}{(\nu(t-s))^{\frac{1}{2}}}\|\eta(t)\|_{C^{1+\alpha}}\|\sigma(s)\|_{\alpha, p} \tag{94}
\end{gather*}
$$

so that the last term is bounded by

$$
\begin{equation*}
C\left(\frac{t}{\nu}\right)^{\frac{1}{2}}\|\eta(t)\|_{C^{1+\alpha}}\|\sigma\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{95}
\end{equation*}
$$

We finish the proof by replacing $\eta$ by $X^{\prime}$ using Theorem 2 .
Theorem 7. Let $0<\alpha<1,1<p<\infty$ and let $T>0$. Let $X^{\prime} \in \operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)$ with $\partial_{t} X^{\prime} \in$ $L^{\infty}\left(0, T ; C^{1+\alpha}\right)$. There exists a constant $C(\alpha)$ depending only on $\alpha$ such that

$$
\begin{align*}
& \left\|\left[X^{\prime} \circ X^{-1} \cdot \nabla, \mathbb{G}\right]\left(\tau \circ X^{-1}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \\
\leq & \left(\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}+\left\|X^{\prime}\right\|_{\operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)} T^{\frac{1}{2}}\right) R \tag{96}
\end{align*}
$$

where $R$ is a polynomial function on $\|\tau\|_{\text {Lip }\left(0, T ; C^{\alpha, p}\right)},\|X-\mathrm{Id}\|_{\text {Lip }\left(0, T ; C^{1+\alpha}\right)}$, whose coefficients depend on $\alpha, \nu$, and $T$, and in particular it grows polynomially in $T$ and bounded below.

Proof. Again we denote $\eta=X^{\prime} \circ X^{-1}$. Also it suffices to bound

$$
\begin{equation*}
[\eta \cdot \nabla, \Gamma]\left(\tau \circ X^{-1}\right)=\eta(t) \cdot \nabla \Gamma\left(\tau \circ X^{-1}\right)-\Gamma\left(\eta \cdot \nabla\left(\tau \circ X^{-1}\right)\right) \tag{97}
\end{equation*}
$$

where $\Gamma$ is as defined in (55), since

$$
\begin{equation*}
[\eta \cdot \nabla, \mathbb{G}]=(R \otimes R) \mathbb{H}[\eta \cdot \nabla, \Gamma]+[\eta(t) \cdot \nabla,(R \otimes R) \mathbb{H}] \Gamma \tag{98}
\end{equation*}
$$

and the second term is bounded by Lemma 1. For the first term, we have

$$
\begin{equation*}
[\eta \cdot \nabla, \Gamma]\left(\tau \circ X^{-1}\right)(t)=I_{1}+I_{2}+I_{3}+I_{4}+I_{5}+I_{6} \tag{99}
\end{equation*}
$$

where

$$
\begin{gather*}
I_{1}=\int_{0}^{t} \eta(t) \cdot\left(\nabla \Delta g_{\nu(t-s)} *\left(\tau \circ X^{-1}(t)\right)\right)-\nabla \Delta g_{\nu(t-s)} *\left(\eta(t) \tau \circ X^{-1}(t)\right) d s \\
I_{2}=\int_{0}^{t} \eta(t) \cdot\left(\nabla \Delta g_{\nu(t-s)} *\left(\tau \circ X^{-1}(s)-\tau \circ X^{-1}(t)\right)\right) \\
\\
-\nabla \Delta g_{\nu(t-s)} *\left(\eta(t)\left(\tau \circ X^{-1}(s)-\tau \circ X^{-1}(t)\right)\right) d s  \tag{100}\\
I_{3}=-\int_{0}^{t} \nabla \Delta g_{\nu(t-s)} *\left((\eta(s)-\eta(t))\left(\tau \circ X^{-1}(s)\right)\right) d s \\
I_{4}=\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\nabla \cdot(\eta(s)-\eta(t)) \tau \circ X^{-1}(s)\right) d s \\
I_{5}=\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\nabla \cdot \eta(t)\left(\tau \circ X^{-1}(s)-\tau \circ X^{-1}(t)\right)\right) d s \\
I_{6}=-\frac{1}{\nu}\left(\nabla \cdot \eta(t) \tau \circ X^{-1}(t)-g_{\nu t} *\left(\nabla \cdot \eta(t) \tau \circ X^{-1}(t)\right)\right)
\end{gather*}
$$

First, $I_{1}+I_{6}$ can be bounded:

$$
\begin{align*}
I_{1}+I_{6}=\frac{1}{\nu}(\eta(t) \cdot \nabla & \left.\left(g_{\nu t} *\left(\tau \circ X^{-1}(t)\right)\right)-\nabla\left(g_{\nu t} *\left(\eta(t) \tau \circ X^{-1}(t)\right)\right)\right) \\
& -\frac{1}{\nu} g_{\nu t} *\left(\nabla \cdot \eta(t)\left(\tau \circ X^{-1}(t)\right)\right) \tag{101}
\end{align*}
$$

and the first term is treated in the same way as (92). Since the first term is

$$
\begin{equation*}
\frac{1}{\nu}\left(\int_{\mathbb{R}^{d}} \nabla g_{\nu t}(y) y \cdot \int_{0}^{1} \nabla \eta(x-(1-\lambda) y, t) d \lambda\left(\tau \circ X^{-1}\right)(x-y, t) d y\right) \tag{102}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|\nabla g_{\nu t}(y) y\right\|_{L^{1}} \leq C \tag{103}
\end{equation*}
$$

the $C^{\alpha, p}$-norm of the first term is bounded by

$$
\begin{equation*}
\frac{C}{\nu}\|\eta(t)\|_{C^{1+\alpha}}\left\|\tau \circ X^{-1}(t)\right\|_{\alpha, p} \tag{104}
\end{equation*}
$$

The $C^{\alpha, p}$-norm of the second term is also bounded by the same bound. Therefore,

$$
\begin{equation*}
\left\|I_{1}+I_{6}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq \frac{C}{\nu} M_{X}^{1+3 \alpha}\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{105}
\end{equation*}
$$

The term $I_{3}$ is bounded due to Theorem 2. Since $\eta \in \operatorname{Lip}\left(0, T ; C^{\alpha}\right)$ we have

$$
\begin{gather*}
\left\|I_{3}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq \frac{C}{\nu}\left(\frac{T}{\nu}\right)^{\frac{1}{2}} M_{X}^{1+4 \alpha}\|X-\mathrm{Id}\|_{\operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)}  \tag{106}\\
\left\|X^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}
\end{gather*}
$$

The terms $I_{4}$, and $I_{5}$ are treated in the spirit of Theorem 5. We treat $L^{p} \cap L^{\infty}$ norm and Hölder seminorm separately. For the term $I_{5}$, we have

$$
\begin{equation*}
I_{5}=\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\nabla \cdot \eta(t)\left(\Delta_{1} \tau(s, t)+\Delta_{2} \tau(s, t)\right)\right) d s \tag{107}
\end{equation*}
$$

where $\Delta_{1} \tau$ and $\Delta_{2} \tau$ are the same as (59). From the same arguments from the above,

$$
\begin{align*}
& \left\|\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\nabla \cdot \eta(t) \Delta_{1} \tau(s, t)\right) d s\right\|_{\alpha, p}  \tag{108}\\
& \leq \frac{C t}{\nu}\|\eta\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{\operatorname{Lip}\left(0, T ; C^{\alpha, p}\right)} M_{X}^{\alpha}
\end{align*}
$$

On the other hand,

$$
\begin{align*}
& \Delta g_{\nu(t-s)} *\left(\nabla \cdot \eta(t) \Delta_{2} \tau(s, t)\right)(x)=\int_{\mathbb{R}^{d}}(K(x, z, t, s)(\nabla \cdot \eta)(X(z, t), t)  \tag{109}\\
& \left.+\Delta g_{\nu(t-s)}(x-X(z, t))((\nabla \cdot \eta)(X(z, s), t)-(\nabla \cdot \eta)(X(z, t), t))\right) d z
\end{align*}
$$

where $K$ is as in (63). Then as in the proof of Lemma 3, by the generalized Young's inequality we have

$$
\begin{gather*}
\left\|\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\nabla \cdot \eta(t) \Delta_{2} \tau(s, t)\right) d s\right\|_{L^{p} \cap L^{\infty}} \leq C\|\tau(t)\|_{L^{p} \cap L^{\infty}}\|\eta\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} \\
\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\left(\frac{t^{\alpha}}{\nu \alpha}+\left(\frac{t}{\nu^{3}}\right)^{\frac{1}{2}}+\frac{t^{2}}{\nu^{3}}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{3}\right) \tag{110}
\end{gather*}
$$

For the Hölder seminorm, we repeat the same argument in the proof of Theorem 5, using the bound (81). Then we obtain

$$
\begin{gather*}
\frac{1}{|h|^{\alpha}}\left\|\delta_{h}\left(\int_{0}^{t} \Delta g_{\nu(t-s)} * \Delta_{2} \tau(s, t) d s\right)\right\|_{L^{\infty}} \\
\leq \frac{C(\alpha)}{\nu}\left(1+\left(\frac{t}{\nu}\right)^{\frac{1}{2}}+\left(\frac{t}{\nu}\right)^{2}\right)\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}\|\eta\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} . \tag{111}
\end{gather*}
$$

Therefore,

$$
\begin{gather*}
\left\|I_{5}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq \frac{C(\alpha)}{\nu}\left(1+t+\left(\frac{t}{\nu}\right)^{2}\right)\left(1+\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\right)^{3} M_{X}^{1+2 \alpha}  \tag{112}\\
\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L i p\left(0, T ; C^{\alpha, p}\right)}
\end{gather*}
$$

The term $I_{4}(t)$ is treated in the exactly same way, by noting that

$$
\begin{align*}
\nabla \cdot(\eta(s)-\eta(t)) & =\nabla_{x} X^{-1}(s):\left(\Delta_{1} \nabla_{a} X^{\prime}(s, t)\right)+\nabla_{x} X^{-1}(s):\left(\Delta_{2} \nabla_{a} X^{\prime}(s, t)\right)  \tag{113}\\
+ & \left(\nabla_{x} X^{-1}(s)-\nabla_{x} X^{-1}(t)\right):\left(\nabla_{a} X^{\prime} \circ X^{-1}\right)(t)
\end{align*}
$$

where as in (59)

$$
\begin{align*}
\Delta_{1} \nabla_{a} X^{\prime}(x, s, t) & =\nabla_{a} X^{\prime}\left(X^{-1}(x, s), s\right)-\nabla_{a} X^{\prime}\left(X^{-1}(x, s), t\right)  \tag{114}\\
\Delta_{2} \nabla_{a} X^{\prime}(x, s, t) & =\nabla_{a} X^{\prime}\left(X^{-1}(x, s), t\right)-\nabla_{a} X^{\prime}\left(X^{-1}(x, t), t\right)
\end{align*}
$$

and

$$
\begin{equation*}
\nabla_{x}\left(X^{-1}(x, s)-X^{-1}(x, t)\right)=\left(\nabla_{a} X \circ X^{-1}\right)(x, t)\left(\nabla_{a}(X-\mathrm{Id})\right)\left(X^{-1}(x, t), t-s\right) \tag{115}
\end{equation*}
$$

so that

$$
\begin{equation*}
\left\|\nabla_{x} X^{-1}(s)-\nabla_{x} X^{-1}(t)\right\|_{C^{\alpha}} \leq|t-s|\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)} M_{X}^{1+2 \alpha} \tag{116}
\end{equation*}
$$

Also note that

$$
\begin{equation*}
\left\|\Delta_{2} \nabla_{a} X^{\prime}(s, t)\right\|_{L^{\infty}} \leq\left\|\nabla_{a} X^{\prime}(t)\right\|_{C^{\alpha}}\|X-\operatorname{Id}\|_{L i p\left(0, T ; L^{\infty}\right)}^{\alpha}|t-s|^{\alpha} \tag{117}
\end{equation*}
$$

so that

$$
\begin{gather*}
\left\|\int_{0}^{t} \Delta g_{\nu(t-s)} *\left(\nabla_{x} X^{-1}(s):\left(\Delta_{2} \nabla_{a} X^{\prime}(s, t)\right) \tau \circ X^{-1}(s)\right) d s\right\|_{C^{\alpha, p}} \\
\leq \frac{C(\alpha)}{\nu}\left(1+t^{\alpha}+\left(\frac{t}{\nu}\right)^{2}\right) M_{X}^{1+2 \alpha}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}  \tag{118}\\
\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}
\end{gather*}
$$

The final result is

$$
\begin{gather*}
\left\|I_{4}(t)\right\|_{\alpha, p} \leq \frac{C(\alpha)}{\nu}\left(1+t+\left(\frac{t}{\nu}\right)^{2}\right) M_{X}^{2+4 \alpha}\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}  \tag{119}\\
+C \frac{t}{\nu} M_{X}^{1+3 \alpha}\left\|X^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} .
\end{gather*}
$$

Finally, $I_{2}$ can be bounded using the combination of the technique in Theorem 5 and Theorem 6. First, we have

$$
\begin{gather*}
I_{2}(x, t)= \\
\int_{0}^{t} \int_{\mathbb{R}^{d}} \nabla \Delta g_{\nu(t-s)}(y) \cdot y \cdot\left(\int_{0}^{1} \nabla \eta(x-(1-\lambda) y, t) d \lambda\left(\Delta_{1} \tau(x-y, s, t)\right)\right) d y d s  \tag{120}\\
+\int_{0}^{t} \int_{\mathbb{R}^{d}} \nabla \Delta g_{\nu(t-s)}(x-z) \cdot(x-z) \cdot\left(\int_{0}^{1} \nabla \eta(\lambda x+(1-\lambda) z, t) d \lambda\left(\Delta_{2} \tau(z, s, t)\right)\right) d z d s .
\end{gather*}
$$

Then applying the argument of the proof of Theorem 6 , the first term is bounded by

$$
\begin{equation*}
\frac{C}{\nu} t M_{X}^{\alpha}\|\eta\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L i p\left(0, T ; C^{\alpha, p}\right)} \tag{121}
\end{equation*}
$$

The second term is treated using the method used in Theorem 5. By changing variables to form a kernel similar to (63), and applying generalized Young's inequality, the $L^{p} \cap L^{\infty}$ norm of the second term is bounded by

$$
\begin{equation*}
\frac{C(\alpha)}{\nu}\left(t^{\alpha}+\left(\frac{t}{\nu}\right)^{\frac{1}{2}}+\left(\frac{t}{\nu}\right)^{2}\right)\left(1+\|X-\operatorname{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\right)^{4}\|\eta\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L^{\infty}\left(0, T ; L^{p} \cap L^{\infty}\right)} \tag{122}
\end{equation*}
$$

Finally, the Hölder seminorm of the second term is bounded by the same method as Theorem 5. The only additional point is the finite difference of $\nabla \eta$ term, but this term is bounded by a straightforward estimate. The bound for the Hölder seminorm of the second term is

$$
\begin{equation*}
\frac{C(\alpha)}{\nu}\left(1+t^{\alpha}+\left(\frac{t}{\nu}\right)^{\frac{1}{2}}+\left(\frac{t}{\nu}\right)^{2}\right)\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\|\eta\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{123}
\end{equation*}
$$

To sum up, we have

$$
\begin{gather*}
\left\|I_{2}(t)\right\|_{\alpha, p} \leq \frac{C(\alpha)}{\nu}\left(1+t+\left(\frac{t}{\nu}\right)^{2}\right)\left(1+\|X-\mathrm{Id}\|_{\operatorname{Lip}\left(0, T ; C^{1+\alpha}\right)}\right)^{4} M_{X}^{1+3 \alpha}  \tag{124}\\
\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}\|\tau\|_{L i p\left(0, T ; C^{\alpha, p}\right)}
\end{gather*}
$$

If we put this together,

$$
\begin{gather*}
\left\|\left[X^{\prime} \circ X^{-1} \cdot \nabla, \mathbb{G}\right]\left(\tau \circ X^{-1}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \\
\leq C\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} M_{X}^{1+2 \alpha}\left\|\Gamma\left(\tau \circ X^{-1}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}  \tag{125}\\
+\left(\left\|X^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)}+\left\|X^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha}\right)} T^{\frac{1}{2}}\right) F_{1}\left(\nu, \alpha, X,\|\tau\|_{L i p\left(0, T ; C^{\alpha, p}\right)}, T\right)
\end{gather*}
$$

where $F_{1}$ depends on the written variables and grows like polynomial in $T,\|\tau\|_{\operatorname{Lip}\left(0, T ; C^{\alpha, p}\right)}$, and $\|X-\mathrm{Id}\|_{\text {Lip }\left(0, T ; C^{1+\alpha}\right)}$. The bound on $\Gamma\left(\tau \circ X^{-1}\right)$ is given by Theorem 5.

## 4 Bounds on variations and variables

Using the results from the previous section we find bounds for variations and variables. For simplicity, we adopt the notation

$$
\begin{equation*}
M_{\epsilon}=1+\left\|X_{\epsilon}-\mathrm{Id}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} \tag{126}
\end{equation*}
$$

First, we bound $\frac{d}{d \epsilon} \mathcal{V}_{\epsilon}$. Note that $X_{\epsilon}(0)=\mathrm{Id}$, so $X_{\epsilon}^{\prime}(0)=0$ and by Theorem 2 and since $X_{\epsilon}^{\prime} \in$ $\operatorname{Lip}\left(0, T ; C^{1+\alpha, p}\right)$ we have

$$
\begin{gather*}
\left\|X_{\epsilon}^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} \leq T\left\|X_{\epsilon}^{\prime}\right\|_{\operatorname{Lip}\left(0, T ; C^{1+\alpha, p}\right)}  \tag{127}\\
\left\|\eta_{\epsilon}(t)\right\|_{C^{\alpha}} \leq t\left\|X^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha, p}\right)} M_{\epsilon}^{\alpha}
\end{gather*}
$$

Then by the Theorem 3, we have

$$
\begin{gather*}
\left\|\eta_{\epsilon} \cdot \mathbb{L}_{\nu}\left(\nabla_{x} u_{\epsilon, 0}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left(\frac{T}{\nu}\right)^{\frac{1}{2}} M_{\epsilon}^{\alpha}\left\|X_{\epsilon}^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha, p}\right)}\left\|u_{\epsilon, 0}\right\|_{1+\alpha, p}  \tag{128}\\
\left\|\mathbb{L}_{\nu}\left(u_{\epsilon, 0}^{\prime}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left\|u_{\epsilon, 0}^{\prime}\right\|_{\alpha, p}
\end{gather*}
$$

By Theorem 6, we have

$$
\begin{gather*}
\left\|\left[\eta_{\epsilon} \cdot \nabla_{x}, \mathbb{U}\right]\left(\sigma_{\epsilon}-u_{\epsilon} \otimes u_{\epsilon}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left(\left(\frac{T}{\nu}\right)^{\frac{1}{2}}+\left(\frac{T}{\nu}\right)\right) M_{\epsilon}^{2+4 \alpha}  \tag{129}\\
\left\|X_{\epsilon}^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\left\|\tau_{\epsilon}-v_{\epsilon} \otimes v_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}
\end{gather*}
$$

and by Theorem 4, we have

$$
\begin{gather*}
\left\|\mathbb{U}\left(\delta_{\epsilon}-\left(v_{\epsilon}^{\prime} \otimes v_{\epsilon}+v_{\epsilon} \otimes v_{\epsilon}^{\prime}\right) \circ X_{\epsilon}^{-1}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left(\frac{T}{\nu}\right)^{\frac{1}{2}} M_{\epsilon}^{\alpha}  \tag{130}\\
\left\|\tau_{\epsilon}^{\prime}-\left(v_{\epsilon}^{\prime} \otimes v_{\epsilon}+v_{\epsilon} \otimes v_{\epsilon}^{\prime}\right)\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}
\end{gather*}
$$

Therefore,

$$
\begin{gather*}
\left\|\frac{d}{d \epsilon} \mathcal{V}_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left\|u_{\epsilon, 0}^{\prime}\right\|_{\alpha, p}  \tag{131}\\
+S_{1}(T)\left(\left\|X_{\epsilon}^{\prime}\right\|_{\operatorname{Lip}\left(0, T ; C^{1+\alpha, p}\right)}+\left\|v_{\epsilon}^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}+\left\|\sigma_{\epsilon, 0}^{\prime}\right\|_{\alpha, p}+\left\|\tau_{\epsilon}^{\prime}\right\|_{L i p\left(0, T ; C^{\alpha, p}\right)}\right) Q_{1}
\end{gather*}
$$

where $S_{1}(T)$ vanishes as $T^{\frac{1}{2}}$ as $T \rightarrow 0$ and $Q_{1}$ is a polynomial in $\left\|u_{\epsilon, 0}\right\|_{1+\alpha, p},\left\|X_{\epsilon}-\operatorname{Id}\right\|_{\text {Lip }\left(0, T ; C^{1+\alpha, p}\right)}$, $\left\|\tau_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}$, and $\left\|v_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}$, whose coefficients depend on $\nu$. Similarly,

$$
\begin{equation*}
\left\|g_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq M_{X}^{\alpha}\left\|u_{0}\right\|_{1+\alpha, p}+C_{1}\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\left\|\sigma_{\epsilon, 0}\right\|_{\alpha, p}+S_{2}(T) Q_{2} \tag{132}
\end{equation*}
$$

where $S_{2}(T)$ vanishes as $T^{\frac{1}{2}}$ as $T \rightarrow 0$ and $Q_{2}$ is polynomial in $\|\tau\|_{L i p\left(0, T ; C^{\alpha, p}\right)}$ and $\|X-\operatorname{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}$, whose coefficients depend on $\alpha$ and $\nu$. Also

$$
\begin{gather*}
\left\|g_{\epsilon}^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq C\left(\left\|u_{\epsilon, 0}^{\prime}\right\|_{1+\alpha, p}+\|X-\mathrm{Id}\|_{L i p\left(0, T ; C^{1+\alpha}\right)}^{\alpha}\left\|\tau_{\epsilon, 0}^{\prime}\right\|_{\alpha, p}\right)  \tag{133}\\
+S_{3}(T)\left(\left\|X_{\epsilon}^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha, p}\right)}+\left\|\sigma_{\epsilon, 0}^{\prime}\right\|_{\alpha, p}+\left\|\tau_{\epsilon}^{\prime}\right\|_{L i p\left(0, T ; C^{\alpha, p}\right)}+\left\|v_{\epsilon}^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha, p}\right)}\right) Q_{3}
\end{gather*}
$$

where $S_{3}(T)$ vanishes as $T^{\frac{1}{2}}$ as $T \rightarrow 0$ and $Q_{3}$ is polynomial in $\left\|u_{\epsilon, 0}\right\|_{1+\alpha, p},\|X-\operatorname{Id}\|_{L i p\left(0, T ; C^{1+\alpha, p}\right.}$, $\|\tau\|_{\operatorname{Lip}\left(0, T ; C^{\alpha, p}\right)}$, and $\left\|v_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha, p}\right)}$, whose coefficients depend on $\nu$ and $\alpha$. Then we have

$$
\begin{equation*}
\left\|\nabla_{a} \frac{d}{d \epsilon} \mathcal{V}_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq T\left\|X_{\epsilon}^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha}\right)}\left\|g_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}+M_{\epsilon}\left\|g_{\epsilon}^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \tag{134}
\end{equation*}
$$

and

$$
\begin{gather*}
\left\|\frac{d}{d \epsilon} \mathcal{T}_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq 2\left\|g_{\epsilon}^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}\left(\left\|\tau_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}+2 \rho K\right)  \tag{135}\\
+\left\|\tau_{\epsilon}^{\prime}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}\left(\left\|g_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}+2 k\right)
\end{gather*}
$$

## 5 Local existence

We define the function space $\mathcal{P}_{1}$ and the set $\mathcal{I}$,

$$
\begin{gather*}
\mathcal{P}_{1}=\operatorname{Lip}\left(0, T ; C^{1+\alpha, p}\right) \times \operatorname{Lip}\left(0, T ; C^{\alpha, p}\right) \times L^{\infty}\left(0, T ; C^{1+\alpha, p}\right) \\
\mathcal{I}=\left\{(X, \tau, v):\|(X-\operatorname{Id}, \tau, v)\|_{\mathcal{P}_{1}} \leq \Gamma, v=\frac{d X}{d t}\right\} \tag{136}
\end{gather*}
$$

where $\Gamma>0$ and $T>0$ are to be determined. Now, for given $u_{0} \in C^{1+\alpha, p}$ divergence free and $\sigma_{0} \in C^{\alpha, p}$ we define the map

$$
\begin{equation*}
(X, \tau, v) \rightarrow \mathcal{S}(X, \tau, v)=\left(X^{\text {new }}, \tau^{n e w}, v^{n e w}\right) \tag{137}
\end{equation*}
$$

where

$$
\left\{\begin{align*}
X^{n e w}(t)= & \mathrm{Id}+\int_{0}^{t} \mathcal{V}(X(s), \tau(s), v(s)) d s  \tag{138}\\
\tau^{\text {new }}(t)= & \sigma_{0}+\int_{0}^{t} \mathcal{T}(X(s), \tau(s), v(s)) d s \\
& v^{\text {new }}(t)=\mathcal{V}(X, \tau, v)
\end{align*}\right.
$$

If $(X-\operatorname{Id}, \tau, v) \in \mathcal{P}_{1}$, then $\left(X^{n e w}-\mathrm{Id}, \tau^{\text {new }}, v^{\text {new }}\right) \in \mathcal{P}_{1}$ for any choice of $T>0$. Moreover, we have the following:
Theorem 8. For given $u_{0} \in C^{1+\alpha, p}$ divergence free and $\sigma_{0} \in C^{\alpha, p}$, there is a $\Gamma>0$ and $T>0$ such that the map $\mathcal{S}$ of (138) maps $\mathcal{I}$ to itself.

Proof. It is obvious that $\frac{d}{d t} X^{n e w}=v^{n e w}$. For the size of $\mathcal{S}(X, \tau, v)$, first note that if $(X-\operatorname{Id}, \tau, v)_{\mathcal{P}_{1}} \leq \Gamma$, then

$$
\begin{equation*}
M_{X}=1+\|X-\mathrm{Id}\|_{L^{\infty}\left(0, T ; C^{1+\alpha}\right)} \leq 1+T \Gamma \tag{139}
\end{equation*}
$$

Applying Theorem 3 and Theorem 4, we know that

$$
\begin{equation*}
\|\mathcal{V}\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq\left\|u_{0}\right\|_{\alpha, p}+A_{1}(T) B_{1}\left(\Gamma,\left\|u_{0}\right\|_{\alpha, p},\left\|\sigma_{0}\right\|_{\alpha, p}\right) \tag{140}
\end{equation*}
$$

where $A_{1}(T)$ vanishes like $T^{\frac{1}{2}}$ for small $T>0$ and $B_{1}$ is a polynomial in its arguments, and some coefficients depend on $\nu$. We estimate

$$
\begin{equation*}
\|g\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq\left\|u_{0}\right\|_{1+\alpha, p}+C_{1} \Gamma^{\alpha}\left\|\sigma_{0}\right\|_{\alpha, p}+A_{2}(T) B_{2}\left(\Gamma,\left\|u_{0}\right\|_{1+\alpha, p},\left\|\sigma_{0}\right\|_{\alpha, p}\right) \tag{141}
\end{equation*}
$$

where $C_{1}$ is as in Theorem 5, depending only on $\alpha$ and $\nu, A_{2}(T)$ vanishes in the same order as $A_{1}(T)$ as $T \rightarrow 0$, and $B_{2}$ is a polynomial in its arguments, and some coefficients depend on $\nu$ and $\alpha$. From (24) we conclude

$$
\begin{equation*}
\|\mathcal{V}\|_{L^{\infty}\left(0, T ; C^{1+\alpha, p}\right)} \leq K_{1}\left(\left\|u_{0}\right\|_{1+\alpha, p}+\Gamma^{\alpha}\left\|\sigma_{0}\right\|_{\alpha, p}\right)+A_{3}(T) B_{3}\left(\Gamma,\left\|u_{0}\right\|_{1+\alpha, p},\left\|\sigma_{0}\right\|_{\alpha, p}\right) \tag{142}
\end{equation*}
$$

where $K_{1}$ is a constant depending only on $\nu$ and $\alpha$, and $A_{3}$ and $B_{3}$ have the same properties as previous $A_{i} \mathrm{~s}$ and $B_{i} \mathrm{~s}$. Now we measure $\mathcal{T}$. From (84) and the previous estimate on $g$ we have

$$
\begin{align*}
\|\mathcal{T}\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)} \leq K_{2}\left(\left\|u_{0}\right\|_{1+\alpha, p}(\rho K+\right. & \left.\left.\left\|\sigma_{0}\right\|_{\alpha, p}\right)+\left\|\sigma_{0}\right\|_{\alpha, p}\left(\Gamma^{\alpha}\left\|\sigma_{0}\right\|_{\alpha, p}+\rho K \Gamma^{\alpha}+k\right)\right)  \tag{143}\\
+ & A_{4} B_{4}
\end{align*}
$$

where $K_{2}$ is a constant depending on $\nu$ and $\alpha$, and $A_{4}$ and $B_{4}$ are as before. Since $\alpha<1$, we can appropriately choose large $\Gamma>\left\|\sigma_{0}\right\|_{\alpha, p}+\left\|u_{0}\right\|_{1+\alpha, p}$ and correspondingly small $\frac{1}{6}>T>0$ so that the right side of (142) and (143) are bounded by $\frac{\Gamma}{6}$. Then $\left\|\left(X^{\text {new }}-\mathrm{Id}, \tau^{\text {new }}, v^{\text {new }}\right)\right\|_{\mathcal{P}_{1}} \leq \Gamma$.

We show now that $\mathcal{S}$ is a contraction mapping on $\mathcal{I}$ for a short time.
Theorem 9. For given $u_{0} \in C^{1+\alpha, p}$ divergence free and $\sigma_{0} \in C^{\alpha, p}$, there is a $\Gamma$ and $T>0$, depending only on $\left\|u_{0}\right\|_{1+\alpha, p}$ and $\left\|\sigma_{0}\right\|_{\alpha, p}$, such that the map $\mathcal{S}$ is a contraction mapping on $\mathcal{I}=\mathcal{I}(\Gamma, T)$, that is

$$
\begin{equation*}
\left\|\mathcal{S}\left(X_{2}, \tau_{2}, v_{2}\right)-\mathcal{S}\left(X_{1}, \tau_{1}, v_{1}\right)\right\|_{\mathcal{P}_{1}} \leq \frac{1}{2}\left\|\left(X_{2}-X_{1}, \tau_{2}-\tau_{1}, v_{2}-v_{1}\right)\right\|_{\mathcal{P}_{1}} \tag{144}
\end{equation*}
$$

Proof. First from Theorem 8 we can find a $\Gamma$ and $T_{0}>0$, depending only on the size of initial data, say

$$
\begin{equation*}
N=\max \left\{\left\|u_{0}\right\|_{1+\alpha, p},\left\|\sigma_{0}\right\|_{\alpha, p}\right\} \tag{145}
\end{equation*}
$$

which guarantees that $\mathcal{S}$ maps $\mathcal{I}$ to itself. This property still holds if we replace $T_{0}$ by any smaller $T>0$. In view of the fact that $\mathcal{I}$ is convex, we put

$$
\begin{gather*}
X_{\epsilon}=(2-\epsilon) X_{1}+(\epsilon-1) X_{2} \\
\tau_{\epsilon}=(2-\epsilon) \tau_{1}+(\epsilon-1) \tau_{2}, 1 \leq \epsilon \leq 2 \tag{146}
\end{gather*}
$$

Then $\left(X_{\epsilon}, \tau_{\epsilon}, v_{\epsilon}\right) \in \mathcal{I}, v_{\epsilon}=(2-\epsilon) v_{1}+(\epsilon-1) v_{2}, u_{\epsilon, 0}=u_{0}$, and $\sigma_{\epsilon, 0}=\sigma_{0}$. This means that

$$
\begin{equation*}
X_{\epsilon}^{\prime}=X_{2}-X_{1}, v_{\epsilon}^{\prime}=v_{2}-v_{1}, u_{\epsilon, 0}^{\prime}=0, \sigma_{\epsilon, 0}^{\prime}=0 \tag{147}
\end{equation*}
$$

Then from the results of Section 4, we see that

$$
\begin{align*}
&\left\|\frac{d}{d \epsilon} \mathcal{V}_{\epsilon}\right\|_{L^{\infty}\left(0, T ; C^{1+\alpha, p}\right)} \leq\left(\left\|X_{2}-X_{1}\right\|_{L i p\left(0, T ; C^{1+\alpha, p}\right)}+\left\|v_{2}-v_{1}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}\right. \\
&\left.+\left\|\tau_{2}-\tau_{1}\right\|_{L i p\left(0, T ; C^{\alpha, p}\right)}\right) S_{1}^{\prime}(T) Q_{1}^{\prime}(\Gamma) \\
&\left\|\mathcal{X}_{\epsilon}^{\prime}\right\|_{L i p\left(0, T ; C^{1+\alpha, p}\right)} \leq\left(\left\|X_{2}-X_{1}\right\|_{L i p\left(0, T ; C^{1+\alpha, p}\right)}+\left\|v_{2}-v_{1}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}\right.  \tag{148}\\
&\left.+\left\|\tau_{2}-\tau_{1}\right\|_{L i p\left(0, T ; C^{\alpha, p}\right)}\right) S_{2}^{\prime}(T) Q_{2}^{\prime}(\Gamma) \\
&\left\|\pi_{\epsilon}\right\|_{L i p\left(0, T ; C^{\alpha, p}\right)} \leq\left(\left\|X_{2}-X_{1}\right\|_{L i p\left(0, T ; C^{1+\alpha, p}\right)}+\left\|v_{2}-v_{1}\right\|_{L^{\infty}\left(0, T ; C^{\alpha, p}\right)}\right. \\
&\left.+\left\|\tau_{2}-\tau_{1}\right\|_{L i p\left(0, T ; C^{\alpha, p}\right)}\right) S_{3}^{\prime}(T) Q_{3}^{\prime}(\Gamma)
\end{align*}
$$

where $\mathcal{X}_{\epsilon}^{\prime}$ and $\pi_{\epsilon}$ are defined in $(20), S_{1}^{\prime}(T), S_{2}^{\prime}(T), S_{3}^{\prime}(T)$ vanish at the rate of $T^{\frac{1}{2}}$ as $T \rightarrow 0$, and $Q_{1}^{\prime}(\Gamma), Q_{2}^{\prime}(\Gamma), Q_{3}^{\prime}(\Gamma)$ are polynomials in $\Gamma$, whose coefficients depend only on $\nu$ and $\alpha$. By choosing $0<T<T_{0}$ small enough, depending on the size of $Q_{i}^{\prime}(\Gamma)$ s, we conclude the proof.

We have obtained a solution to the system (6) in the path space $\mathcal{P}_{1}$ for a short time, that is, we have $(X, \tau, v)$ satisfying $v=\frac{d X}{d t}$ and satisfying (14). We also have Lipschitz dependence on initial data, Theorem 1 .

Proof. We repeat the calculation of the Theorem 9, but this time $u_{\epsilon, 0}^{\prime}=u_{1}(0)-u_{2}(0)$ and $\sigma_{\epsilon, 0}^{\prime}=$ $\sigma_{1}(0)-\sigma_{2}(0)$. Then we choose $T_{0}$ small enough that $S_{i}^{\prime}\left(T_{0}\right) Q_{1}^{\prime}(\Gamma)<\frac{1}{2}$.

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## References

[1] P. Constantin. Lagrangian-Eulerian methods for uniqueness in hydrodynamic systems Adv. Math., 278:67-102, 2015.
[2] P. Constantin. Analysis of hydrodynamic models, volume 90 of CBMS-NSF Regional Conference Series in Applied Mathematics, Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 2017.


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